



NUMERICAL SOLUTION OF THE EFFECTS OF HEAT AND MASS TRANSFER ON UNSTEADY MHD FREE CONVECTION FLOW PAST AN INFINITE VERTICAL PLATE

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ABSTRACT

This study attempts to explore a qualitative analysis of the effects of Soret on an unsteady magnetohydrodynamics free convection flow of a chemically reacting incompressible fluid past an infinite vertical plate embedded in a porous medium taking the source of heat and thermal radiation into account as well as viscous dissipation. The central equations are scrupulously converted into sets of coupled nonlinear partial differential equations for providing logical solutions. The method of Galerkin finite element is used considering appropriate boundary conditions for diverse physical metrics and then numerically analyzed employing MATLAB. A significant change in velocity, temperature, concentration profiles is observed for various values of Prandtl number, Grashof number, Eckert number, Soret number. It is noticed that the velocity profile enhances by enhancing the values of porous medium as well as it decreases when 'M' and Prandtl number increases. The temperature profile decreases as for the increasing vales of heat source parameter and also the concentration profile increases as Soret number increases.

Keywords: Galerkin Method; Magnetohydrodynamics; Chemical Reaction; Viscous Dissipation; Heat Source; Soret Effect

1. INTRODUCTION

The natural physical process of convection flow has been gaining increasing research attention owing to its increased utilization in diverse physical, chemical and engineering applications. Natural convection is a physical process of transferring heat and mass among the fluids initiated by the variation in the concentration of both the temperature as well as the species resulting in the differences in density and encouraging the forces of buoyancy to operate on the fluids. The natural convection of the Nano fluids from a vertical accelerated plate in the presence of the radiation flux and magnetic field is observed in this study by Astuti et al. (2019). Nature contains enormous quantities of such flows. This process of convection is studied in depth due to its extensive free applications in the environments of geophysical sciences and engineering. Room heating inside the buildings utilizing radiators could well illustrate the application of convection for transferring heat. Buoyancy or decrease in gravity Middleman (1998), Rubin et al. (2001) is the prime factor causing these natural flows of convection due to the variation in temperature resulting in corresponding differences in fluid density, including liquids and gases. It has been identified that the flows induced by buoyancy in the fluid-saturated porous media have been extensively employed in a wide variety of applications of thermal engineering applications and domains such as geothermic system spring water pollution, exchangers of heat, thermal insulation, piling up of nuclear effluence, packed bed catalytic reactors, cooling of electronic devices, the security of energy systems, atmosphere and oceanic circulation.

Several researchers are drawn towards the field of Magnetohydrodynamics due to its vast and varied applications in the realms of geophysical science and astrophysics. MHD is largely utilized in the advanced studies on the astronomical matter, solar and stellar structures, propagation of radio through the medium of ionosphere, etc. Several applications associated with oil extraction, recovery of heat, flow-through devices of filtration, and storage of thermal energy largely make use of the applications based on the progression of convection through permeable media, besides MHD bearings and MHD pumps. The modern times witness an enhanced focus on this heat transfer process through convection due to its ever-enhancing prominence in diverse technological innovations and applications such as fiber and granular insulation, electronic system cooling, cool combustors, and porous material regenerative heat exchangers, geo-thermal reservoirs, and drying the porous solids. This feature of concurrent heat and mass transfer from entirely diverse geometrics embedded in porous media encompasses a range of engineering and geophysical science applications as have been mentioned.

Soundalgekar and Wavre (1977) have explored an unstable free convection flow of convection past an infinite vertical plate that contains continuous suction and shifting of mass. Agrawal H et al. (1984) furnish the influence that Hall currents could have on hydro-magnetic free convection with the transfer of mass in an extremely rotating fluid. Jang and Ni (1989) have analyzed transient-free convection with a transfer of mass from an equal vertical plate embedded in a medium that has an extremely permeable nature. Jha (1991) studied the characteristics of the free convection of MHD and the subsequent flow of transfer of mass through a porous medium. Prabhakar Reddy (2020) has studied how the transfer of mass impacts the free convective flow on an unsteady MHD

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containing incompressible viscous dissipative fluid past an infinite vertical porous plate. MHD flow of carreau nano fluid explores using CNT over a nonlinear elongated sheet is studied by Nagalakshmi et al. (2020).

The contribution of radiation becomes important for air in instances where the wall temperature ranges between the var6000 – 10,000K. Such instances occur when vehicles in space attempt reentry into the atmosphere. Korycki (2006) illustrates this radioactive heat transfer as a vital and fundamental phenomenon that exists in all sensible engineering processes as can be found in radiation within the buildings, foundry engineering, and set processes, die formation, chemical engineering, and diverse composite structures applied in manufacturing.

Thermal radiation has become an integral and potential component in the field of fluid dynamics emerging as one of the premium branches in engineering sciences with its prominent aspect of fusing diverse scenarios in mechanical, chemical, aerospace, environmental, stellar energy and hazards engineering. Bhaskara Reddy and Bathaiah (1981) & (1982) have analyzed the MHD free streamline flow of convection of an incompressible elastic fluid. Further, they have studied the combined and strained flow of convection through two parallel porous walls. Elabashbeshy (1997) has studied the transfer of mass and heat on a vertical plate along with magnetic flux. Samad, Karim, and Muhammad (2010) have made numerical computations to estimate the influence of the thermal radiation on steady MHD free convection flow considering the diffusion approximation of Rosseland. Loganathan and Arasu (2010) have analyzed the results of the deposition of thermophoresis on non-Darcy MHD combined convective transfer of mass and heat past a porous wedge while suction or injection is present.

This paper attempts to examine, the effects of transfer of mass and heat of free convection flow- considering the viscous dissipation of a fluid past an infinite vertical porous plate embedded in a permeable medium when a homogeneous magnetic field is present. Partial differential equations of coupled non-linear systems govern the principles of the problem with inbuilt complexities to obtain precise solutions. Hence, the method of Galerkin finite element, with its computational cost efficiency, has been employed for obtaining the solutions. The consequent changes and compoment of diverse aspects such as the concentration velocity, temperature, Sherwood and Nusselt numbers have been comprehensively focused for observing the possible changes in the governing framework and metrics.

Crepeau and Clarksean (1997) have considered the classical drawback in instances of natural convection flow from an equal vertical plate along with the term of supplementary heat generation within the equation of energy. They found out that there must be exponential decay in internal heat generation with the classical variable of similarity for finding a real similar solution. Reddy and Rao (2011) examined the Soret effect on unstable MHD free convection transfer flow of mass and heat past an infinite vertical permeable plate in conjunction with the Hall current, heat source, and consequent thermal radiation. Pal and Talukdar (2012) have detailed the impact that the thermal radiation and the chemical reaction of first-order could have on unstable MHD convective transfer flow of mass and heat along with the Soret effect of a viscous fluid past a semi-infinite vertical flat plate when heat source and oscillatory suction are present.

Chemical reaction processes are both uniform and heterogeneous depending on the location of their occurrences such as an interface or single section volume. Gehart et al. (1971) have studied how foreign mass can impact the free flow of convection beyond a semi-infinite vertical plate. The far-off mass existing either in water or in the air can cause significant chemical process. In general, the reaction rate of chemical processes corresponds to the concentration levels of the species. For example, convection associated with the internal generation of heat has a pivotal role to play in the framework of the method of complete heat transfer i.e., the acquisition of metal waste from the utilized fuel and the processes of thermal combustion activity.

Alim et al. (2008) have investigated the impact of viscous dissipation over natural flow of convection along a sphere having heat loss of radiation, while Salina et al. (2010) have examined the effect of viscous dissipation over natural convection flow of convection along a sphere with the generation of heat. Alam et al. (2007) have examined the impact of viscous dissipation on MHD natural convection flow on a sphere besides heat generation.

This paper attempts to examine, the effects of transfer of mass and heat of free convection flow- considering the viscous dissipation of a fluid past an infinite vertical porous plate embedded in a permeable medium when a homogeneous magnetic field is present. Partial differential equations of coupled non-linear systems govern the principles of the problem with inbuilt complexities to obtain precise solutions. Hence, the method of Galerkin finite element, with its computational cost efficiency, has been employed for obtaining the solutions. The consequent changes and compoment of diverse aspects such as the concentration velocity, temperature, Sherwood and Nusselt numbers have been comprehensively focused for observing the possible changes in the governing framework and metrics.

2. MATHEMATICAL ANALYSIS

Our purpose in this paper is to study the unstable spontaneous flow of convection considering the viscous dissipation of a fluid which has the properties of electric conduction and hydro-magnetic radiation, past on infinite heated vertical plate which has been embedded in a porous medium. Fig. 1 represents the physical model including the coordinate system. The x' - axis is considered here for alongside the vertical plate, while y' - axis for normal to the plate. At the point of

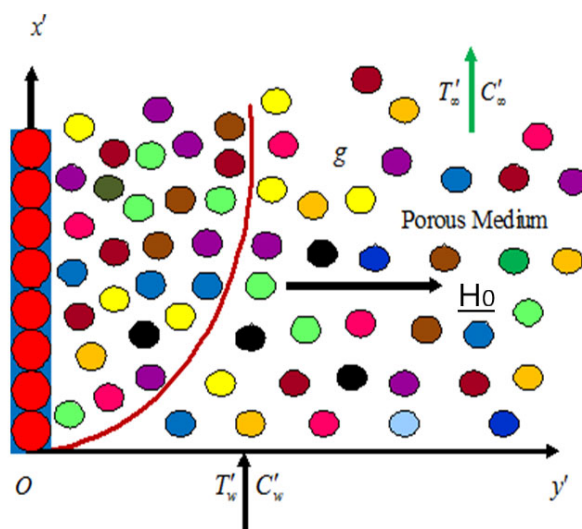


Fig.1 Schematic Diagram

time $t'=0$, the temperatures of T' and T_∞' have been steadily maintained for the plate as well as the fluid respectively and presumed to be identical. Similarly, it has been assumed that concentration levels of pertaining to the species, C_w' at the plate and C_∞' all through the fluid, to be identical. The plate's temperature at time $t'<0$, which has been hitherto maintained as constant, now is altered resulting in the flow of convection currents very near to the plate. Consequently, the supply of mass to the plate is at a constant rate, while the functions of time t' and y' in this scenario, become the flow variables.

We now consider flow of an incompressible fluid making the following assumptions.

- i) All the fluid properties except the density in the buoyancy force term are constant.

- ii) A homogenous magnetic field with an intensity of H_0 is executed diagonally to the plate.
- iii) The magnetic field which has been thus induced is ignored since the flow's magnetic Reynolds number is taken to be very inconsequential.
- iv) $\bar{E} = 0$ the electric field is zero.
- v) The Boussinesq approximation is applied.
- vi) The Hall effect of magnetohydrodynamics and magnetic dissipation (Joule heating of the fluid) are neglected.
- vii) $T_w > T_\infty$
- viii) The heat source Q^* is of the type $Q^* = Q_0(T'_\infty - T')$.

- ix) The concentration of the diffusing species in the binary mixture is assumed to be very small in comparison with the other chemical species Dufour effect is negligible.
- x) The chemical reaction is considered to be homogeneous and of first order.
- xi) Also, no applied or polarized voltages are assumed to exist, so that the effect of polarization of fluid is negligible.

With the foregoing assumptions and under the usual boundary layer approximations, unsteady flow is governed by the following partial differential equations

Continuity Equation:

$$\frac{\partial v'}{\partial y'} = 0 \tag{1}$$

Momentum Equation:

$$\frac{\partial u'}{\partial t'} = \nu \frac{\partial^2 u'}{\partial y'^2} + g\beta(T'_\infty - T') + g\beta^*(C' - C'_\infty) - \frac{\sigma\mu_0^2 H_0^2 u'}{\rho} - \frac{\nu u'}{K'} \tag{2}$$

Energy Equation:

$$\rho C_p \frac{\partial T'}{\partial t'} = k \frac{\partial^2 T'}{\partial y'^2} - \frac{\partial q_r}{\partial y'} + \mu \left(\frac{\partial u'}{\partial y'} \right)^2 - Q_0(T' - T'_\infty) \tag{3}$$

Concentration Equation:

$$\frac{\partial C'}{\partial t'} = D_M \frac{\partial^2 C'}{\partial y'^2} - K_r(C' - C'_\infty) + \frac{D_M K_T}{T_m} \frac{\partial^2 T'}{\partial y'^2} \tag{4}$$

The boundary conditions for the fields of velocity, temperature, and concentration respectively are as follows:

$$\left. \begin{aligned} t' \leq 0, T' = T'_\infty, C' = C'_\infty \quad \forall y' \\ u' = 0, T' = T'_w, C' = C'_w \quad \text{at } y' = 0 \\ u' = 0, T' \rightarrow T'_\infty, C' \rightarrow C'_\infty \quad \text{as } y' \rightarrow \infty \end{aligned} \right\} \tag{5}$$

Where u' is the velocity along the x' - axis, ν is the kinematic coefficient of viscosity, g is the acceleration due to gravity, β is the coefficient of volume expansion for the heat transfer, β^* is the volume coefficient of expansion with species concentration, T' is the fluid temperature, T'_∞ is the fluid temperature at infinity, T'_w and C'_w are the wall dimensional temperature and concentration respectively. C' is the species concentration, C'_∞ is the species concentration at infinity, D_M is the chemical molecular diffusivity, k_1 is the mean absorption coefficient, K' is the constant permeability of the medium, μ is the coefficient of viscosity, C_p is the specific heat at constant pressure, ρ is the density of fluid, K_r - the chemical reaction parameter and t is the time. Equation (1) gives $v' = -v_0(v_0 > 0)$

Where v_0 is the constant suction velocity with the negative sign implying it in the direction of the plate.

Employing the method of Roseland approximation, the term for radiative heat flux is depicted by $q_r = -\frac{4\sigma}{3k_1} \frac{\partial T'^4}{\partial y'}$ (6)

where $\sigma = 5.67 \times 10^{-8} W/m^2 K^4$ is the Stefan- Boltzmann constant and k_1 is the Roselan mean absorption coefficient. Here, it has been assumed that the variations in the temperature in the flow are insignificant enough that T'^4 could be expressed as the temperature's linear function. This is achieved through the expansion of T'^4 a Taylor series and ignoring the terms of a greater order. Thus

$$T'^4 \cong 4T'^3_\infty T' - 3T'^4_\infty \tag{7}$$

Then, by the application of (5) and (6) in equation (2), can be reduced as given below:

$$\rho C_p \frac{\partial T'}{\partial t'} = k \frac{\partial^2 T'}{\partial y'^2} + \frac{16\sigma T'^3_\infty}{3k_1} \frac{\partial^2 T'}{\partial y'^2} + \mu \left(\frac{\partial u'}{\partial y'} \right)^2 \tag{8}$$

The non-dimensional quantities mentioned below are introduced, so that the governing equations and boundary condition can be rendered in dimensionless form:

$$\begin{aligned} M &= \frac{\sigma\mu_0^2 H_0^2 \nu}{\rho\nu_0^2}, u = \frac{u'}{u_0}, y = \frac{y' \nu_0}{V}, t = \frac{t' \nu_0^2}{V}, \theta = \frac{T' - T'_\infty}{T'_w - T'_\infty}, \\ C &= \frac{C' - C'_\infty}{C'_w - C'_\infty}, P_r = \frac{\mu C_p}{k}, G_m = \frac{g\beta^*(C'_w - C'_\infty)V}{u_0\nu_0^2}, K_r = \frac{K_r' \nu}{\nu_0^2}, \\ S_c &= \frac{\nu}{D_M}, E_c = \frac{u_0^2}{C_p(T'_w - T'_\infty)}, R = \frac{4\sigma T'^3_\infty}{kk_1}, S_r = \frac{D_M K_T}{\nu T_m(C'_w - C'_\infty)} \\ G_r &= \frac{\nu g\beta(T'_w - T'_\infty)}{u_0\nu_0^2}, C_f = \frac{\tau \omega}{\rho U_0 \nu} = \left(\frac{\partial u}{\partial y} \right)_{y=0}, S = \frac{Q_0}{\rho C_p \nu} \end{aligned} \tag{9}$$

In the context of (8), equations (1), (3), and (7) have been reduced to the non-dimension form as given hereunder:

$$\frac{\partial u}{\partial t} = G_r \theta + G_m C + \frac{\partial^2 u}{\partial y^2} - \left(M + \frac{1}{K} \right) u \tag{10}$$

$$\frac{\partial \theta}{\partial t} = \frac{1}{P_r} \left(1 + \frac{4R}{3} \right) \frac{\partial^2 \theta}{\partial y^2} + E_c \left(\frac{\partial u}{\partial y} \right)^2 - S \theta \tag{11}$$

$$\frac{\partial C}{\partial t} = \frac{1}{S_c} \frac{\partial^2 C}{\partial y^2} - K_r C + S_r \frac{\partial^2 \theta}{\partial y^2} \tag{12}$$

And these corresponding boundary conditions are

$$\left. \begin{aligned} t \leq 0 : u = 0, \theta = 0, C = 0 \forall y \\ t > 0 : u = 0, \theta = 1, C = 1 \text{ at } y = 0 \\ u \rightarrow 0, \theta \rightarrow 0, C \rightarrow 0 \text{ as } y \rightarrow \infty \end{aligned} \right\} \tag{13}$$

3. NUMERICAL TECHNIQUE

The technique of finite element has been employed for solving the issues of non-dimensional momentum as well as the equations of energy (10), (11) and (12) in addition to the imposed conditions of the boundary (13).

In the context of the differential equations, the Galerkin expression (10) becomes

$$\int_{y_j}^{y_k} \left\{ N^T \left[\frac{\partial^2 u^{(e)}}{\partial y^2} - \frac{\partial u^{(e)}}{\partial t} - Ru^{(e)} + P \right] \right\} dy = 0$$

Where $N^T = [N_j \quad N_k]^T = \begin{bmatrix} N_j \\ N_k \end{bmatrix}$

$$R = M + \frac{1}{K}, \quad P = (Gr)\theta + (G_m)C;$$

Let us assume the linear piece-wise approximation solution to be

$$u^{(e)} = N_j(y)u_j(t) + N_k(y)u_k(t) = N_j u_j + N_k u_k$$

Then, the equation of element is represented as

$$\int_{y_j}^{y_k} \left\{ \begin{bmatrix} N'_j N'_j & N'_j N'_k \\ N'_j N'_k & N'_k N'_k \end{bmatrix} \begin{bmatrix} u_j \\ u_k \end{bmatrix} \right\} dy + \int_{y_j}^{y_k} \left\{ \begin{bmatrix} N_j N_j & N_j N_k \\ N_j N_k & N_k N_k \end{bmatrix} \begin{bmatrix} \dot{u}_j \\ \dot{u}_k \end{bmatrix} \right\} dy + R \int_{y_j}^{y_k} \left\{ \begin{bmatrix} N_j N_j & N_j N_k \\ N_j N_k & N_k N_k \end{bmatrix} \begin{bmatrix} u_j \\ u_k \end{bmatrix} \right\} dy = P \int_{y_j}^{y_k} \begin{bmatrix} N_j \\ N_k \end{bmatrix} dy$$

Where prime and dot represent differentiation w.r.to 'y' and 't' respectively. On simplification, we derive

$$\frac{1}{l^{(e)^2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} u_j \\ u_k \end{bmatrix} + \frac{1}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{bmatrix} \dot{u}_j \\ \dot{u}_k \end{bmatrix} + \frac{R}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{bmatrix} u_j \\ u_k \end{bmatrix} = \frac{P}{2} \begin{bmatrix} 1 \\ 1 \end{bmatrix} \text{ Where}$$

$$l^{(e)} = y_k - y_j = h$$

For obtaining the differential equation at the knot x_i , we have written the element equations for the two elements $y_{i-1} \leq y \leq y_i$ and $y_i \leq y \leq y_{i+1}$.

Assembling these two, we get

$$\frac{1}{l^{(e)^2} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} u_{i-1} \\ u_i \\ u_{i+1} \end{bmatrix} + \frac{1}{6} \begin{bmatrix} 2 & 1 & 0 \\ 1 & 4 & 1 \\ 0 & 1 & 2 \end{bmatrix} \begin{bmatrix} \dot{u}_{i-1} \\ \dot{u}_i \\ \dot{u}_{i+1} \end{bmatrix} + \frac{R}{6} \begin{bmatrix} 2 & 1 & 0 \\ 1 & 4 & 1 \\ 0 & 1 & 2 \end{bmatrix} \begin{bmatrix} u_{i-1} \\ u_i \\ u_{i+1} \end{bmatrix} = \frac{P}{2} \begin{bmatrix} 1 \\ 2 \\ 1 \end{bmatrix}$$

The row equation of the knot 'i', is represented as

$$\frac{1}{l^{(e)^2} [-u_{i-1} + 2u_i - u_{i+1}] + \frac{1}{6} [\dot{u}_{i-1} + 4\dot{u}_i + \dot{u}_{i+1}] + \frac{R}{6} [u_{i-1} + 4u_i + u_{i+1}] = P$$

On the application of the method of Crank-Nicholson to the above equation, we derive

$$A_1 u_{i-1}^{n+1} + A_2 u_i^{n+1} + A_3 u_{i+1}^{n+1} = A_4 u_{i-1}^n + A_5 u_i^n + A_6 u_{i+1}^n + 12Pk$$

Where

$$\left. \begin{aligned} A_1 &= -6r + 2 + Rk, \\ A_2 &= 12r + 8 + 4Rk, \\ A_3 &= -6r + 2 + Rk, \\ A_4 &= 6r + 2 - Rk, \\ A_5 &= -12r + 8 - 4Rk, \\ A_6 &= 6r + 2 - Rk \end{aligned} \right\} \quad (14)$$

When a similar procedure is applied to equation (10), we obtain

$$B_1 \theta_{i-1}^{n+1} + B_2 \theta_i^{n+1} + B_3 \theta_{i+1}^{n+1} = B_4 \theta_{i-1}^n + B_5 \theta_i^n + B_6 \theta_{i+1}^n + 12kQP_r$$

Where

$$\left. \begin{aligned} B_1 &= 2P_r - 6r - 8Rr + SkP_r, \\ B_2 &= 8P_r + 12r + 16Rr + 4SkP_r, \\ B_3 &= 2P_r - 6r - 8Rr + SkP_r, \\ B_4 &= 2P_r + 6r + 8Rr - SkP_r, \\ B_5 &= 8P_r - 12r - 16Rr - 4SkP_r, \\ B_6 &= 2P_r + 6r + 8Rr - SkP_r \end{aligned} \right\} \quad (15)$$

When a similar procedure is applied to equation (11), we obtain

$$E_1 C_{i-1}^{n+1} + E_2 C_i^{n+1} + E_3 C_{i+1}^{n+1} = E_4 C_{i-1}^n + E_5 C_i^n + E_6 C_{i+1}^n + Sr \frac{\partial^2 \theta}{\partial y^2} 12kSc$$

Where

$$\left. \begin{aligned} E_1 &= -6r + 2S_c + K_r S_c k, \\ E_2 &= 12r + 8S_c + 4K_r S_c k, \\ E_3 &= -6r + 2S_c + K_r S_c k, \\ E_4 &= 6r + 2S_c - K_r S_c k, \\ E_5 &= -12r + 8S_c - 4K_r S_c k, \\ E_6 &= 6r + 2S_c - K_r S_c k \end{aligned} \right\} \quad (16)$$

Here, $r = \frac{k}{h^2}$, while k, h are the sizes of mesh in the direction along y

and time-direction, with indices I and J referring to space and time respectively. The mesh system consists of $h=0.1$ and $k=0.001$. In the equations (11), (12) and (13) considering $i=1$ in (1) and employing the boundary conditions (10), we obtain the subsequent system equations:

$A_i X_i = B_i$ for $i=1(1)n$ Where A_i 's are matrices of order n and X_i, B_i 's are column matrices with n components. Thomas algorithm employed to acquire the relevant solutions with regard to the system of equations considered above for the aspects of velocity, temperature, and concentration. Besides the numerical solutions can be acquired for the considered system of equations by running the same on MATLAB. MATLAB has been employed with reduced values of 'h' and 'k' to establish the convergence and stability pertaining to the the method of Galerkin finite element, where considerable variations in the values of u, T and C have not been perceived. Consequently, it can be established that the method of Galerkin finite element is stable and convergent.

Study of grid independence:

To analyze grid independency/dependency it is better to test the solution for several mesh sizes and get a range at which there is no variation in the solution. The mathematical estimation of velocity for specific inputs of mesh(grid) size at time $t= 0.2$, are shown in the above Table 1. From this table, we observed that, variation of velocity is closer for various mesh size at time $t=0.2$. Hence, we conclude that, the computational results are stable and converge.

Table 1:

The numerical values of u for variation of mesh (Grid) sizes (when $Pr=0.71, K=0.5, M=2, Ec=0.1, Sc=0.22, Gm=3, Gr=3, Kr=1.0, R=2, S=2, Sr=2$).

y	u		
	Size of Mesh k=0.001, h=0.1	Size of Mesh k = 0.0005, h=0.05	Size of Mesh k = 0.00025, h=0.025
0	0	0	0
0.5	0.364470	0.364025	0.363979
1	0.280719	0.280738	0.280831
1.5	0.155177	0.155504	0.155662
2	0.074048	0.074378	0.074514
2.5	0.031897	0.032127	0.032218
3	0.012464	0.012601	0.012653
3.5	0.004401	0.004473	0.004499
4	0.001398	0.001431	0.001443
4.5	0.000398	0.000411	0.000416
5	0.000101	0.000105	0.000107

4. RESULT AND DISCUSSION

This current study addresses the issues pertaining to an unsteady free convection flow of mass transfer considering the viscous dissipation of a fluid past an infinite vertical permeable plate. Innumerable numerical computations have been applied for the velocity and the non-dimensional, temperature, and concentration 'C', for various parameter indicated below:

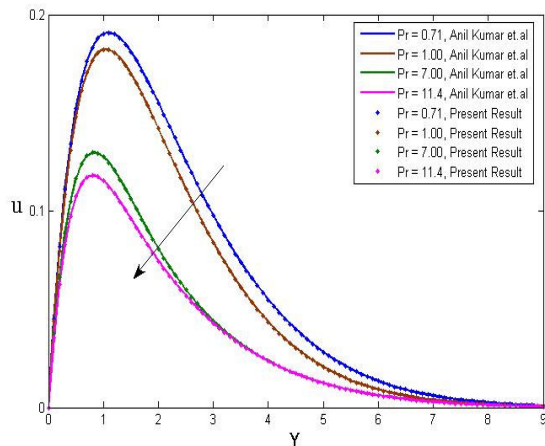


Fig. 2 Profile of Velocity for diverse values of Pr when $K=0.5$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $Gr=3$, $Kr=1.0$, $R=2$.

It can be perceived from Fig. 2, that the profile of velocity decreases corresponding to the increase in Prandtl number Pr . This study makes a comparative analysis with the available previous studies without the source of heat and the Soret effect for establishing the accuracy and validity of the output numerical results. The profiles of velocity for diverse values of Pr have been studied in comparison with the available solutions proposed by Anil Kumar et al. (2018) and found to largely concur. It is further perceived that the profile of velocity decreases corresponding to the increase in Pr . This could be consequent to the fact that fluids having greater Pr contain higher levels of viscosities causing a reduction in velocities. Fig. 3 illustrates that the profile of velocity decreases corresponding to the increase in Prandtl number Pr . It is further perceived that the profile of velocity decreases corresponding to the increase in Pr . This could be consequent to the fact that fluids having greater Pr contain higher levels of viscosities causing a reduction in velocities.

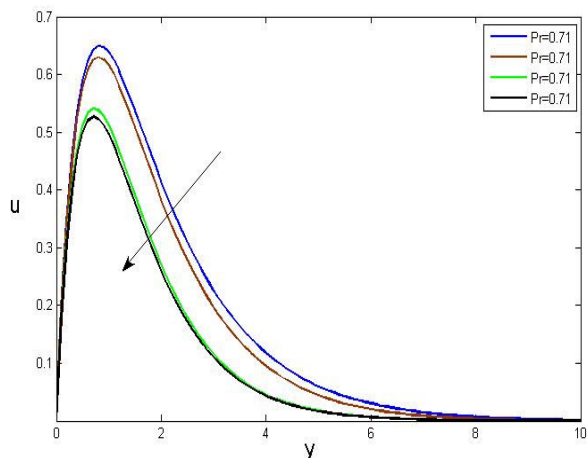


Fig. 3 Profile of Velocity for diverse values of Pr when $K=0.5$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $Gr=3$, $Kr=1.0$, $R=2$, $S=2$, $Sr=2$.

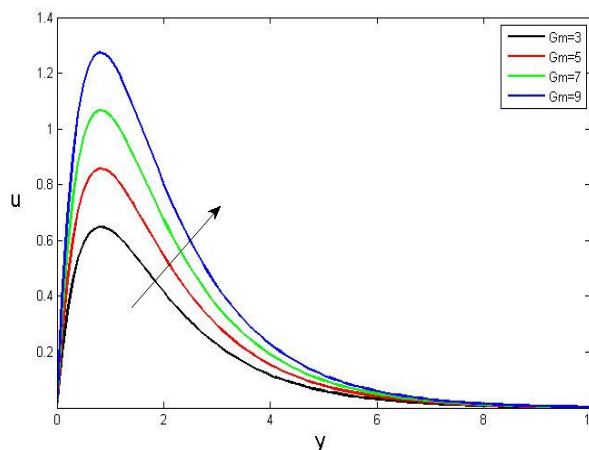


Fig. 4 Profile of Velocity for diverse values of Gm when $Pr=0.71$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gr=3$, $R=2$, $K=0.5$, $Kr=1.0$, $S=2$, $Sr=2$.

It is evident from Fig. 4 that the velocity increases corresponding to the increase in Grashof number Gm . The ratio between the viscous and buoyancy forces operating on a fluid is approximated through Modified Grashof number, where any increase of Gm results in a corresponding increase in the forces of buoyancy besides a corresponding decrease in the viscous forces. In instances of decrease of viscosity, the fluid's internal resistance decreases leading to an automatic increase in the fluid velocity. The ratio between the forces of viscosity and buoyancy operating on a fluid is approximated through Grashof number Gr . Here, any increase of Gr results in a corresponding increase in the forces of buoyancy besides a corresponding decrease in the viscous forces as depicted in Fig. 5. In instances of decrease of viscosity, the fluid's internal resistance decreases leading to an automatic increase in the fluid velocity. Fig. 6 illustrates the profile of velocity for diverse values of the parameter for porosity K . In the characterization of the properties of transport of mass and heat in a porous medium, the very factor of permeability stands as a prominent parameter. It is evident that there is an increase in K in proportion to the increase in the velocity, as the increased permeability of the medium indicates decreased resistance owing to the presence of the porous matrix within the medium

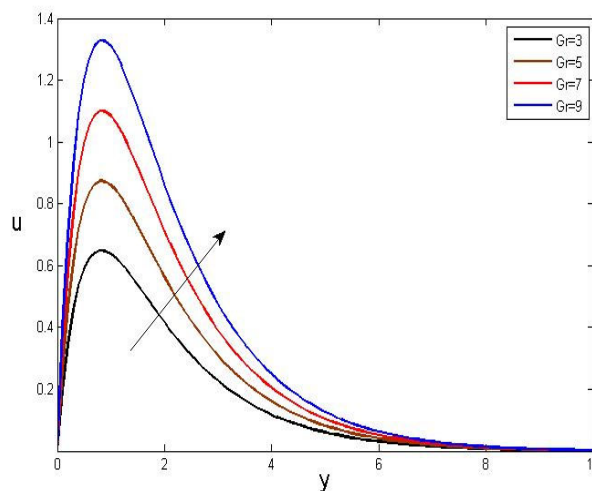


Fig. 5 Profile of Velocity for diverse values of Gr when $Pr=0.71$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $R=2$, $K=0.5$, $Kr=1.0$, $S=2$, $Sr=2$.

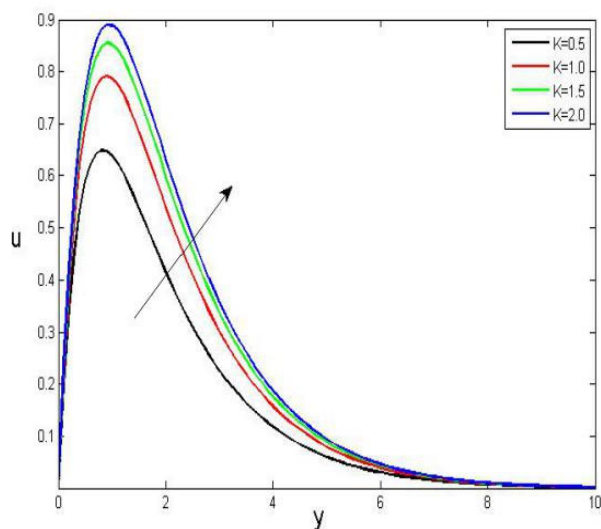


Fig. 6 Profile of Velocity for various values of K when $Pr=0.71$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $R=2$, $Gr=3$, $Kr=1.0$, $S=2$, $Sr=2$.

Fig. 7 illustrates the differences in the profile of velocity in conjunction with the Magnetic field parameter M . It is evident that increasing the Magnetic field results in a corresponding increase in the magnetic strength and the mass of the fluid particle causing the corresponding decrease in velocity. Consequent to the impact of the transverse magnetic field on a fluid that conducts electrically, there is a considerable increase in a resistive sort of force known as Lorentz force, which almost acts as a drag force. Any increase in the value of M leads to a corresponding force of resistance. Fig. 7. has been plotted in such a way as to illustrate the differences in the profile of velocity corresponding to M . It is evident that any increase in M results in a proportionate increase in the magnetic strength while slowing down the fluid motion. Fig. 8 illustrates the profile of temperature for diverse values of Eckert number Ec . The Ec represents the interrelationship of the flow's kinetic energy and enthalpy. This represents the energy transformation from kinetic to internal in terms of work done against the stress of the viscous fluid. This energy manifests in the form of heat during dissipation. Consequently, the dissipative heat triggers an increase in temperature.

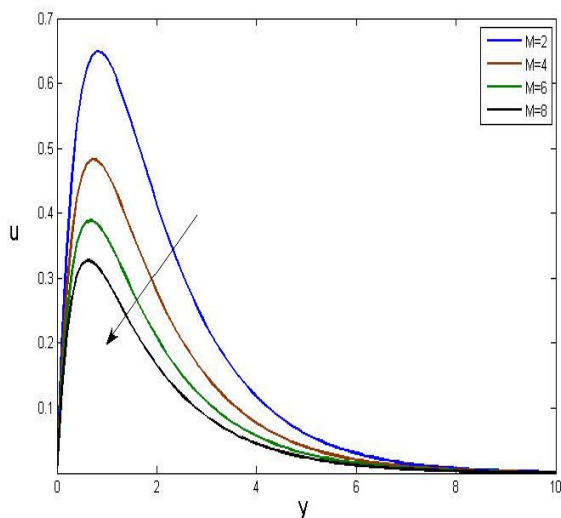


Fig. 7 Profile of Velocity for diverse values of M when $K=0.5$, $Pr=0.71$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $R=2$, $Gr=3$, $Kr=1.0$, $S=2$, $Sr=2$.

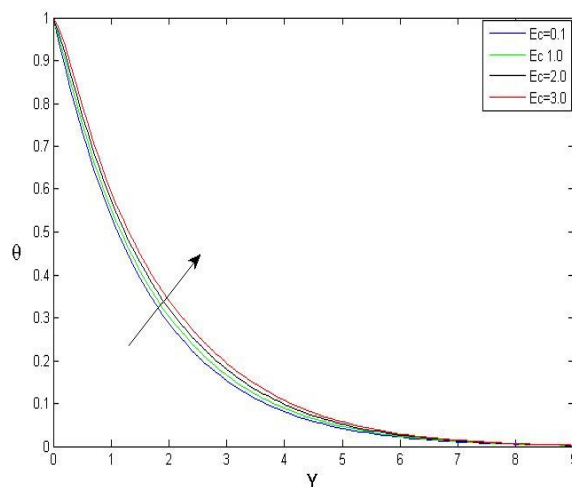


Fig. 8 Profile of Temperature for various values of Ec when $K=0.5$, $M=2$, $Pr=0.71$, $Sc=0.22$, $Gm=3$, $R=2$, $Gr=3$, $Kr=1.0$, $S=2$, $Sr=2$.

Fig. 9 depicts the profile of temperature for diverse values of the parameter of Radiation R . Any increase in R increases the viscosity of the fluid, finally resulting in the decrease both in the temperature and the thickness of the thermal boundary layer. The results as depicted in Fig. 10 illustrate an increase in the value of the Prandtl number Pr causing a corresponding decrease in the field of temperature. Besides, any rise in Pr causes a corresponding decrease both in the temperature and the thickness of the thermal boundary layer. Fig. 11 illustrates the rise in the value of the Heat Source Parameter S correspondingly decreasing the boundary layer as can be expected. This is because the absorption of heat automatically reduces the force of buoyancy which in turn slows down the rate of the flow. Fig. 12 reflects the profile of concentration for diverse values of Schmidt number Sc . The Sc approximates the ratio between the mass diffusivity and kinematic viscosity. Any increase in Sc results in a corresponding increase in the viscous forces. The Sc represents the interrelationship and ratio between the momentum and mass diffusivity. The rise in Sc reduces the concentration correspondingly causing the effect of concentration buoyancy, thereby reducing the fluid velocity considerable.

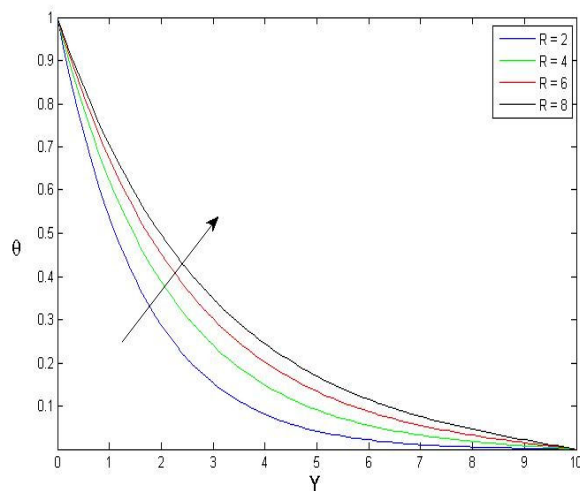


Fig. 9 Profile of Temperature for diverse values of R when $K=0.5$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $Pr=0.71$, $Gr=3$, $Kr=1.0$, $S=2$, $Sr=2$.

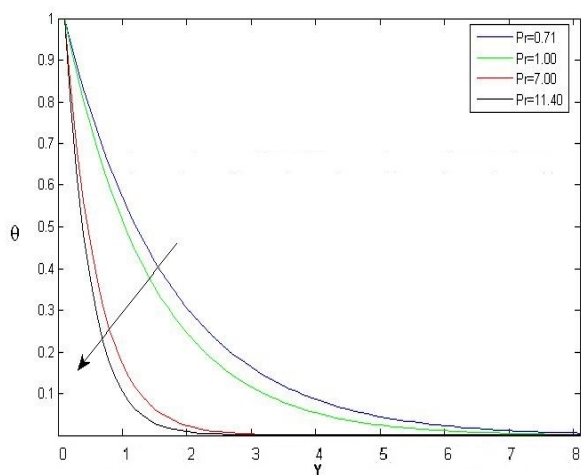


Fig. 10 Profile of Temperature for diverse values of Pr when $K=0.5$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $Gr=3$, $Kr=1.0$, $R=2$, $S=2$, $Sr=2$.

Fig. 13 illustrates the effect of Soret Number Sr on diverse profiles of concentration. This graph considers Sr values as $Sr = 2, 4, 6, 8$. It is evident from the graph that any increase in the values of Sr causes a corresponding increase in the concentration profiles uniformly at every point in the field of flow. This phenomenon occurs because the diffusive species having greater Sr values tend to enhance the profile of concentration. Hence, conclusions can be drawn that the Sr values considerably influence the concentration distributions. Fig. 14 depicts the influence of the parameter of chemical reaction Kr on the profiles of concentration. It is obvious that an increase in Kr results in the reduction of concentration. Fig. 15 illustrates the Skin-friction τ in the context of time t for diverse parameter values. Any increase in M causes a corresponding increase in the τ . The corresponding increase in the τ along with M can be attributed to the increased Lorentz force, which brings in supplementary momentum within the boundary layer. In contrast, an increase in K , Gm and Gr leads to a corresponding reduction in the τ . The extent of τ for $Pr = 0.71$ is low in comparison to that of $Pr = 7$.

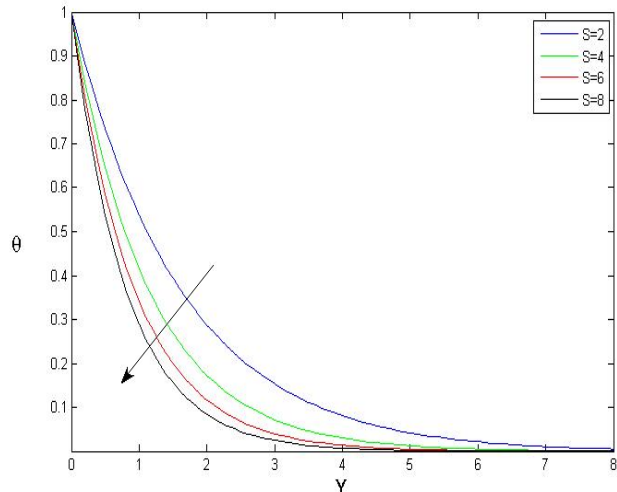


Fig. 11 Profile of Temperature for diverse values of S when $K=0.5$, $M=2$, $Ec=0.1$, $Sc=0.22$, $Gm=3$, $Pr=0.71$, $Gr=3$, $Kr=1.0$, $Sr=2$

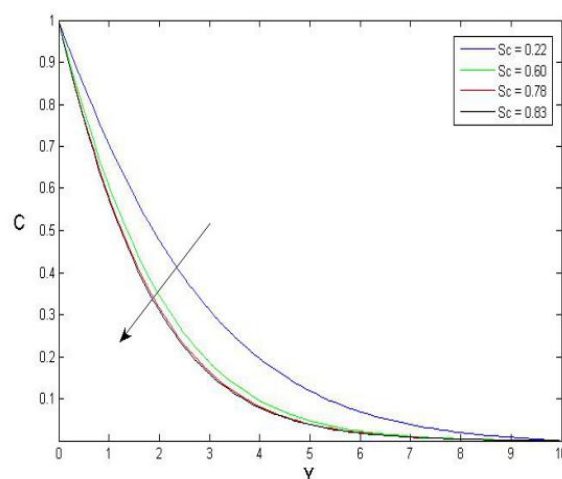


Fig. 12 Profile of Concentration for diverse values of Sc when $K=0.5$, $M=2$, $Ec=0.1$, $Pr=0.71$, $Gm=3$, $R=2$, $Gr=3$, $Kr=1.0$, $S=2$, $Sr=2$.

Fig. 16 illustrates a perceivable increase in the concentration transfer rate corresponding to the rising values of Sc , and Kr , whereas the same is found to decrease when the Sr value is increased. Fig. 17 illustrates the Nusselt number Nu along with the time t for different values of metrics such as Pr , R , Ec , and S . Nu for $Pr=7$ (in this case for water) is found to be greater than the number for $Pr=0.71$. This is because the lower Pr values are equivalent to correspondingly enhancing the thermal conductivities. Consequently, the heat can more swiftly diffuse further away from the plate than with higher Pr values. The same applies to S also. Hence, an increased heat transfer rate can be perceived. It can also be perceived that there is a considerable fall in the heat transfer rate corresponding to the rise in R and Ec .

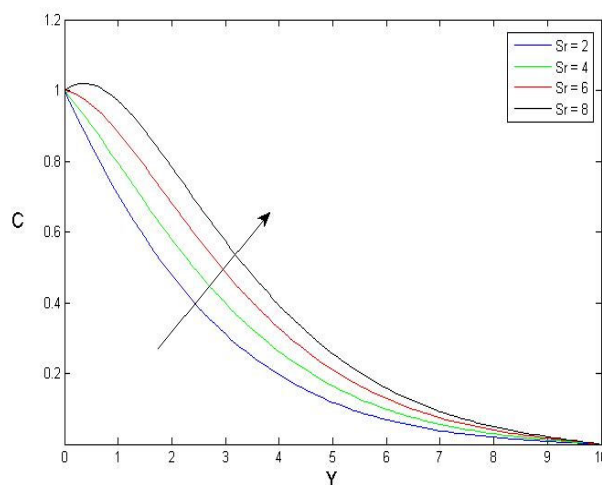


Fig. 13 Profile of Concentration for diverse values of Sr , when $K=0.5$, $M=2$, $Ec=0.1$, $Pr=0.71$, $Gm=3$, $R=2$, $Gr=3$, $Sc=0.22$, $S=2$.

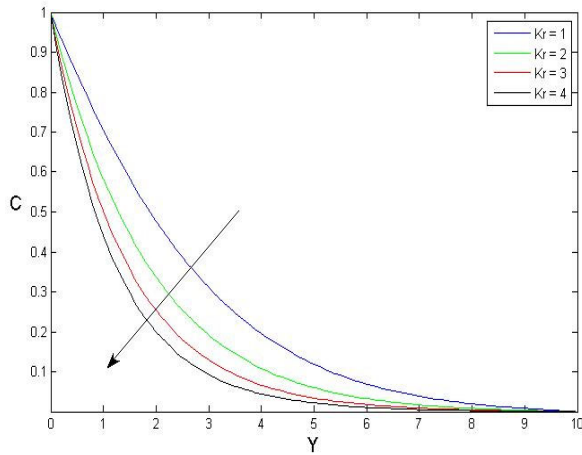


Fig. 14 Profile of Concentration for diverse values of Kr when $K=0.5$, $M=2$, $Ec=0.1$, $Pr=0.71$, $Gm=3$, $R=2$, $Gr=3$, $Sc=0.22$, $S=2$, $Sr=2$.

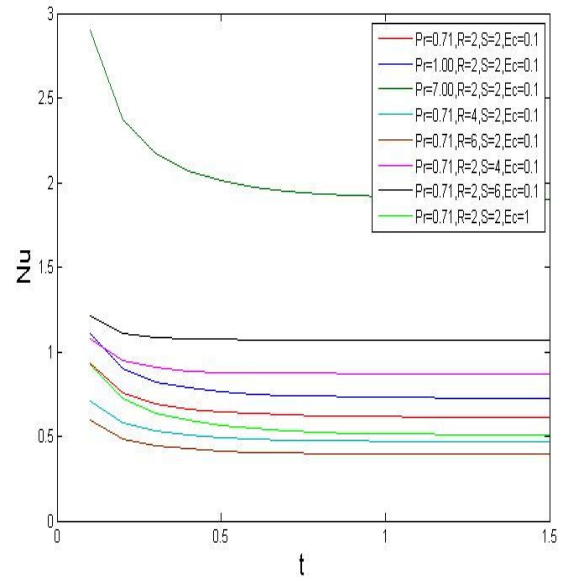


Fig. 17 Profile of the Nusselt number for diverse values of Pr , R , S , and Ec .

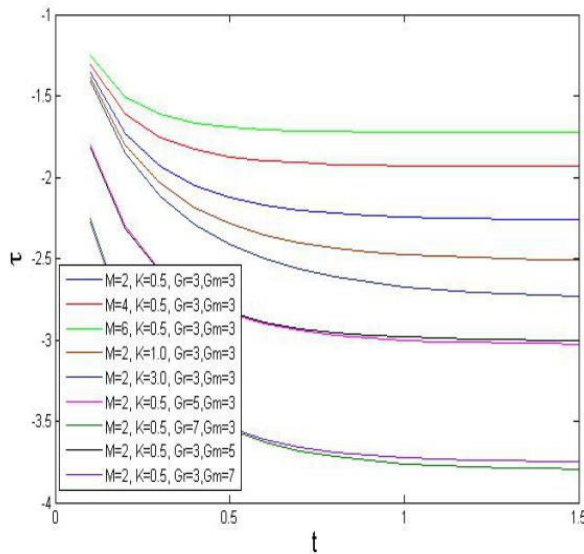


Fig. 15 Profile of Skin Friction for diverse values of M , K , Gr , and Gm .

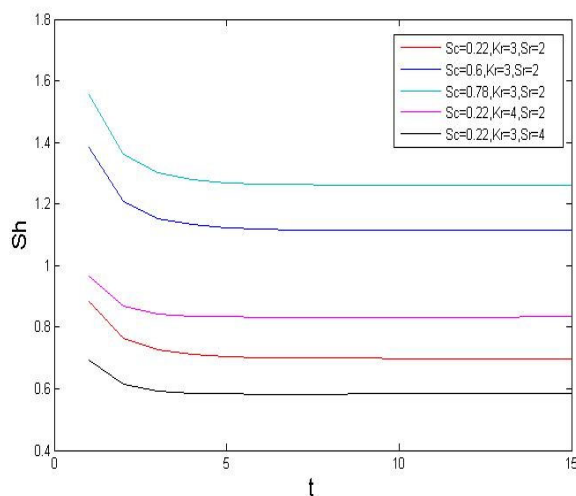


Fig. 16 Profile of the Sherwood number for diverse values of Sc , Kr , and Sr .

5. CONCLUSIONS

In this study, an attempt has been made to analyze numerically the impact of viscous dissipation of thermal radiation on the diffusion flow of mass and heat of MHD past an infinite vertical plate which has been embedded in a permeable medium containing variable conditions of the surface. The method of Galerkin finite element is employed for solving the equations which governs the flow. The following are some of the interesting outcome factors of the study about the physical aspects of concentration, temperature, and the velocity of the flow:

- Velocity increases corresponding to the enhancement in the Gr and Gm numbers.
- Velocity decreases corresponding to the enhancement in the magnetic parameter M and the permeable nature of the porous medium K .
- An increase in Sc and time t causes a decrease in the velocity.
- The temperature rises corresponding to the rise in Ec as well as time t , but reduces with the rise in the parameter of radiation R .
- An increase in Sc causes a corresponding decrease in the level of concentration.
- Greater Sr values tend to enhance the level of concentration.
- A rise in M causes an increase in the Skin-friction τ .
- Increased K , Gm , and Gr cause a corresponding decrease in the τ .
- A rise in the values of Sc correspondingly increases the Sherwood number Sh .
- Increasing M causes a corresponding increase in the τ .
- Increasing the values of K , Gm and Gr causes a corresponding decrease in the τ .
- Enhanced values of Sc cause a corresponding rise in the Sh .
- Enhanced values of Sr cause a corresponding decrease in the Sh .
- Rise in the values of Pr and S result in a corresponding increase in Nu .

NOMENCLATURE:

C_{∞}'	Concentration of the fluid far away from the plate ($Kg\ m^{-3}$)
C_w'	Concentration of the plate ($Kg\ m^{-3}$)
y	Dimensionless displacement (m)
T_{∞}'	Fluid temperature away from the plate (K)
Gm	Grashof number for mass transfer
Gr	Grashof number for heat transfer
u	Non-dimensional fluid velocity ($m\ s^{-1}$)
Sh	The local Sherwood number
u'	Velocity component in x' - direction ($m\ s^{-1}$)
g	Acceleration of gravity, $9.81\ (m\ s^{-2})$
x'	Coordinate axis along the plate (m)
Re	Reynolds number
D	Solute mass diffusivity ($m\ s^{-2}$)
Nu	The local Nusselt number
τ	Skin Friction
B_0	Uniform magnetic field (Tesla)
H_0	Magnetic Induction
ρ	The constant density ($Kg\ m^{-3}$)
β	Volumetric coefficient of thermal expansion (K^{-1})
y'	Co-ordinate axis normal to the plate (m)
Ec	Eckert number
C'	Fluid Concentration ($Kg\ m^{-3}$)
T_w'	Fluid temperature at the wall (K)
T'	Fluid temperature (K)
M	Magnetic parameter
Pr	Prandtl number
q_r	Radiative heat transfer coefficient Sc Schmidt number
C_p	Specific heat at constant Pressure ($J\ Kg^{-1}K$)
ν	Kinematic viscosity ($m^2\ s^{-1}$)
μ_0	Magnetic Permeability ($N.A^{-2}$)
C	Species concentration ($Kg\ m^{-3}$)

Greek Symbols:

K	Thermal conductivity of the fluid ($Wm^{-1}K^{-1}$)
θ	Non dimensional fluid temperature (K)
β^*	Volumetric Coefficient of thermal expansion with concentration ($m^{-3}\ Kg$)
σ	Electric conductivity of the fluid ($s\ m^{-1}$)

Superscripts:

' Dimensionless Properties

Subscripts:

∞	Free stream conditions
p	Plate

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