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Discussion of the Fluid Acceleration Quality of a Ducted Propulsion System on the Propulsive Performance

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ABSTRACT

This paper focuses on the ducted propulsion with the accelerating nozzle, and discusses the influence of its fluid acceleration quality on its propulsive performances, including the hull efficiency, the relative rotative efficiency, the effective wake, and the thrust deduction factor. An actual ducted propulsion system is used as an example for computational analysis. The computational conditions are divided into four combinations, which are provided with different propeller pitches, cambers, and duct lengths. The method applied in this study is the Computational Fluid Dynamics (CFD) technology, and the contents of the calculation include the hull's viscous resistance, the wave-making resistance, the propeller performance curve, and the self-propulsion simulation in order to obtain the ship's effective wake, thrust deduction factor, hull efficiency, and relative rotative efficiency. The performance curve of the propeller and resistance estimation results are compared with the experimental values for determining the correctness of the self-propulsion simulation. According to the computational analysis, it is known that increasing the propeller pitch cannot effectively increase the hull efficiency. The duct acceleration quality can be reduced by shortening the duct length; hence, when the effective wake fraction and thrust deduction factor decrease, the hull efficiency is increased. In addition, the pressure inside the duct is relatively low if the acceleration quality of the duct is too high, which is unfavorable for controlling the propeller cavitation. Moreover, if the hull bottom in front of the propeller is tapered up from the front to the back at an overly steep angle, the thrust deduction factor will be too large and lead to a relatively low hull efficiency.

KEYWORDS

Ducted propellers; accelerating nozzle; hull efficiency; relative rotative efficiency; effective wake; thrust deduction factor

Nomenclature

ρ	Seawater density
Q_0	Propeller open-water torque
η_H	Hull efficiency
R_S	Suction-effect resistance
η_R	Relative rotative efficiency
R_{SM}	Sea-margin resistance



$-C_p$	Pressure coefficient
R_T	Total bare hull resistance
D	Propeller diameter
R_V	Viscous resistance
h	Shaft immersion
R_W	Waving-making resistance
J_A	Advanced ratio
T	Single propeller thrust
K_{Qo}	Torque coefficient
t	Thrust deduction fraction
K_T	Thrust coefficient
V_A	Propeller inflow velocity
n	Propeller speed
V_S	Ship speed
P_0	Atmospheric pressure
w	Wake fraction
P_L	Local pressure on the foil

1 Introduction

The ducted propeller, given its multiple excellent characteristics, has been used increasingly on a variety of vessels in recent years. For a shallow-draft ship, the duct can protect the propeller and prevent stranding. For a tugboat, more thrust is generated by using ducted propellers with accelerating nozzles, so that it can pull large-sized vessels. For some high-speed propellers, the ducted propellers with decelerating nozzles can increase the static pressure inside the duct and reduce propeller cavitation. Therefore, this study aims to adopt the CFD to conduct the self-propulsion simulation, and discuss the fluid acceleration quality on the propulsive performance.

Sung et al. [1] applied the CFD to simulate the performance of underwater vehicles. Kim et al. [2] recommended three turbulence models, namely, the standard $k-\epsilon$ model (SKE), the RNG-based $k-\epsilon$ model (RNG), and the realizable $k-\epsilon$ model (RKE), for calculating the nominal wake of surface vessels. Among the three models, the RKE model could identify the strong vortex position most accurately. Huuva [3] discussed the two-phase flow problem that resulted from the Reynolds-averaged-Navier-Stokes (RANS), Detached-Eddy Simulation (DES), and Large-Eddy simulations (LES) equations. Yang et al. [4] employed the commercial software Fluent with the $k-\omega$ turbulent model to conduct the self-propulsion simulation of large oil tankers. They found that the axial wakes calculated at the circumferential positions, except for the 180° position (six o'clock position), are close to the measured values. Toxopeus et al. [5] analyzed the wake field of the DARPA SUBOFF submarine by the DES turbulent. They used RANS for calculating the boundary layer and the DES for predicting the flow field. Although the difference between the computed results and the experimental values was low, they suggested that the mesh should be optimized, different inflow angles should be considered, and the results of different turbulent models should be compared. Vaz et al. [6] applied the MARIN ReFRESKO program proposed in Vaz et al. [7] and AcuSolve to predict the flow field of the DARPA SUBOFF submarine. They analyzed the maneuvering quality, and found that the difference between the numerical simulations and the measured values was within the acceptable range. Seo et al. [8] adopted flexible meshing techniques in the CFD for the self-propulsion simulation of ships, and the obtained stern wake was relatively small. Korkmaz [9] employed the RANS with explicit algebraic stress model and the potential theory for the self-propulsion simulation of a bulk carrier. The predicted results

were compared with the experimental values, and the comparisons indicated that the bare hull resistances were quite close. Kao et al. [10] used the CFD for the self-propulsion simulation of the DARPA SUBOFF submarine. They compared the simulation results with the measured values of Vaz et al. [6], and found that the maximum error was only 3.5%. This study refers to the calculation settings and methods mentioned by Kao et al. [10], and compares the calculation results on the propeller performance and resistance estimation with the experimental values to determine the correctness of the self-propulsion simulation.

Yu et al. [11] combined the panel method panMARE with the RANS code ANSYS-CFX to calculate the open-water performance of a Ka series propeller with a 19A duct. Their analysis on the flow regime of the fluid at the tip clearance suggested that the SST turbulence model was more suitable for the calculation of a ducted propeller system than the k-epsilon turbulence model. Focused on a ducted propeller, Yongle et al. [12] discussed the influence of its tip clearance on its tip clearance loss at different advance ratios. Gao et al. [13] integrated the CFD technique with the finite element and boundary element methods to estimate the vibration and noise generated by a ducted propeller. Gaggero et al. [14] adopted the OpenFOAM blockMesh library and the RANS solver simpleFOAM for the geometric design of accelerating and decelerating nozzles, in order to increase the thrust and delay the cavitation phenomenon. Gong et al. [15] applied the DES model to calculate the wake vortex evolution of a ducted propeller, and confirmed that the duct can delay the wake contraction and result in different trajectory behaviors of the tip vortices. Andersson et al. [16] utilized the commercial CFD software STAR-CCM + v10.06 to conduct the self-propulsion simulation on the ducted propeller. They also discussed the main engine horsepower loss that resulted from the viscous losses. Gong et al. [17] employed the CFD to calculate the flow field of a ducted propeller with an inclined shaft angle, and measured the tip vortex and blade surface pressure. Kinnas et al. [18] introduced a perturbation potential based panel method with full wake alignment to calculate the ducted propellers. It was shown that the fully aligned blade wake gets more accurate results than a simplified wake alignment model. In Bosschers et al. [19], an iterative procedure to analyze ducted propellers was carried out by a RANS solver and a boundary element method (BEM). The RANS solver analyzed the viscous flow by body forces obtained from the BEM. However, the absence of the tip leakage vortex which leads to a different flow behavior in the tip region should be further investigated. Bontempo et al. [20] applied the axial momentum theory and the nonlinear semi-analytical actuator disk model in analyzing the ducted propeller with a decelerating nozzle. It was indicated that some duct geometrical parameters such as the profile camber, the chord, and the maximum thickness affect the appearance of the cavitation. Bontempo et al. [21] investigated the accelerating duct by the axial momentum theory and the nonlinear semi-analytical actuator disk model. According to Bontempo et al. [21], the propulsive efficiency can be improved by increasing the thickness, chord, camber, or decreasing the incidence angle. Du et al. [22] provided the 3D VII method (the viscous/inviscid interaction method) to solve the flow around three-dimensional wings and propellers. It was shown that the 3D VII method can improve the results significantly as compared with the full-blown RANS simulations. Villa et al. [23] used the viscous-based CFD code to analyze the flow features around the ducted propellers, and got a good agreement with LDV measurements in terms of predicted wake shape and vortex positions. According to Villa et al. [23], the effects of the shaft configuration were non-negligible, while the shape of the boss cap had practically no effects on the computed forces. Amiri Delouei et al. [24], Karimnejad et al. [25], and Afra et al. [26] applied the Immersed boundary method (IBM) for dealing with moving objects in fluids. The IBM has the behavior of good stability. There is no arbitrary parameter used in IBM. The IBM along with diffuse interface scheme is often employed to take into account the forces

arising from the particle and fluid interactions. By using of the fixed, non-body-conformal grids in the IBM, time-consuming re-meshing process for simulation of moving bodies in fluid flows is unnecessary.

This paper is mainly to observe the effect of the fluid acceleration due to the ducted propulsion system. The fluid acceleration is related to the propulsive efficiency and cavitation control. A real workboat case is treated as the computing sample. According to the propulsive theory, the propeller pitch ratio, propeller camber ratio, and the duct length are modified to avoid the problems due to the fluids acceleration. The most effective way to solve the problems caused by the flow acceleration is recommended in this paper. The pressure distribution induced by the fluid acceleration on the hull is also discussed in order to investigate the influence on the hull resistance. This study provides four computing samples with different propellers and ducts. The calculated resistances include the viscous resistance and wave-making resistance. In the self-propulsion simulation, the effective wake, thrust deduction factor, hull efficiency, and relative rotative efficiency of an actual ship are obtained. Besides, in order to correctly simulate the self-propulsion test, the calculated propeller performance curve (K-J chart), and the resistance are compared with the measured values.

2 Numerical Schemes

In this study, the CFD is employed to compute the viscous resistance of the hull, the wave-making resistance, the propeller performance curve (i.e., K–J chart), and simulate the self-propulsion test. The turbulent model adopted in this paper is the SST $k-\omega$ model. The viscous flow is solved in the Reynolds-averaged Navier-Stokes (RANS) equations, and the SIMPLE algorithm is applied. The spatial discretization scheme is second-order accurate. The SIMPLE method, as deduced by Patankar et al. [27], is commonly used to calculate the incompressible flow. The turbulent model, computing settings and mesh generation used in this study are referred to Kao et al. [10].

The self-propulsion simulation is conducted in this study to discuss the performance of a ducted propulsion system with the accelerating nozzle. When the total propeller thrust is equal to the total hull resistance (including the resistance generated by the propeller suction effect (R_S), the viscous resistance (R_V), the waving-making resistance (R_W), and the sea-margin resistance (R_{SM})), the balance point of the self-propulsion is reached. The balance point is as indicated in Eq. (1).

$$\sum^T = R_S + R_V + R_W + R_{SM} \quad (1)$$

Here T is the single propeller thrust.

Then, the thrust coefficient, K_T , is obtained by Eq. (2).

$$K_T = \frac{T}{\rho n^2 D^4} \quad (2)$$

Here ρ is the seawater density, n is the propeller speed, and D is the propeller diameter. According to the thrust identity, the advanced ratio ($J_A = \frac{V_A}{nD}$) and torque coefficient ($K_{Qo} = \frac{Q_0}{\rho n^2 D^5}$) can be interpolated by the propeller K–J chart calculated by CFD. Here V_A is the propeller inflow velocity, Q_0 is the propeller torque in the open-water condition. Thus, at the balance point,

the ship effective wake ($1 - w$), thrust deduction ($1 - t$), hull efficiency (η_H), and relative rotative efficiency (η_R) can be obtained according to following equations:

$$(1 - w) = \frac{V_A}{V_S} = \frac{J_A \times n \times D}{V_S} \tag{3}$$

$$(1 - t) = \frac{R_T}{T} \tag{4}$$

$$\eta_R = \frac{Q_0}{Q} \tag{5}$$

$$\eta_H = \frac{1 - t}{1 - w} \tag{6}$$

Here V_S is the ship speed, R_T is the total bare hull resistance ($R_T = R_V + R_W + R_{SM}$), and Q is the propeller torque in the behind condition. The above computing process is also shown in Fig. 1. Fig. 1 indicates the procedure of the self-propulsion simulation. According to Fig. 1, the bare hull resistances (R_V and R_W) at the specified ship speed, and the propeller performance (i.e., K-J chart) are calculated in advance by the CFD. The propeller speed is unknown, and set to an initial value. Then the total propeller thrust (ΣT), and the self-propulsion resistance (R) without the free surface are computed by the self-propulsion simulation. Determine whether the results of the self-propulsion simulation meet Eq. (7), if not, the propeller speed will be modified, and the process will be executed once again. As the results of the self-propulsion simulation meet Eq. (7), the propulsive factors can be obtained by interpolating K_T value in the K-J chart.

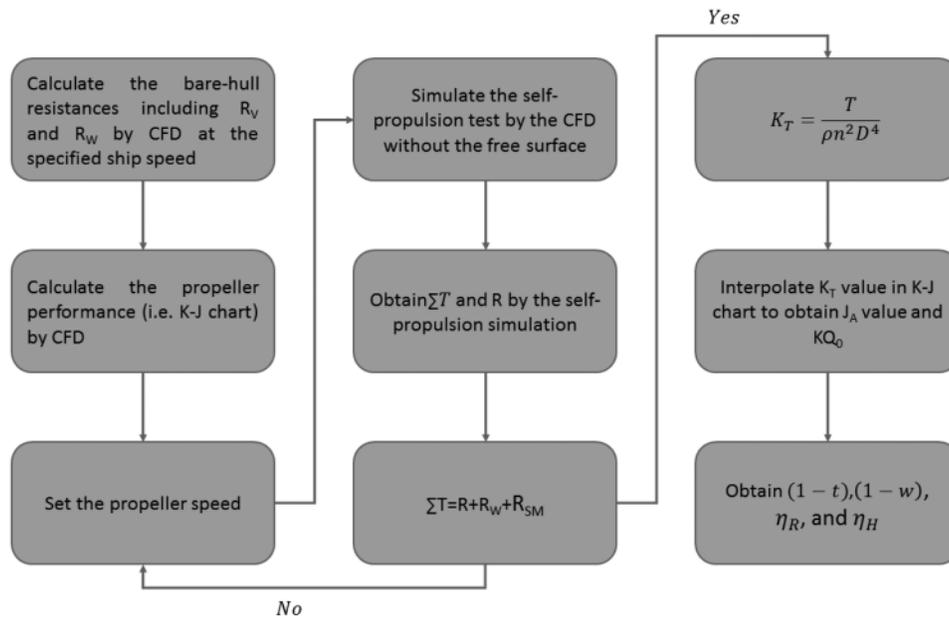


Figure 1: The flowchart for the self-propulsion simulation

The operating propeller and ship hull are both considered in the self-propulsion simulation. It is noteworthy that, in order to save computational time, the free surface is disregarded in the self-propulsion simulation (i.e., disregarding the wave-making resistance (R_W)). The value of wave-making resistance (R_W) is computed by using the bare hull. In other words, the ship resistance (R) obtained by the self-propulsion simulation only includes the resistance caused by the propeller suction effect (R_S) and viscous resistance (R_V). Therefore, when the self-propulsion reaches the balance point, the total propeller thrust is equal to the summation of the ship resistance obtained by the self-propulsion simulation (R), the wave-making resistance (R_W) and the sea-margin resistance (R_{SM}), as expressed in Eq. (7).

$$\begin{aligned}\sum^T &= R + R_W + R_{SM} \\ &= R_S + R_V + R_W + R_{SM}\end{aligned}\quad (7)$$

The dimensional formulation and governing equations in this section are referenced in HYKAT HSVA [28].

The turbulent model, computing settings and mesh generation used in this study are referred to Kao et al. [10]. Table 1 shows that the predicated results by Kao et al. [10] are in good agreement with the measured values. In other words, the turbulent model, computing settings and mesh generation used in this study are verified.

Table 1: A comparison of propeller performance

	J_A	$1 - w$	KT	KQ
Exp.	0.766	0.6851	0.2362	0.4343
CDF	0.751	0.6701	0.2444	0.4448
Error (%)	1.96%	2.19%	3.49%	2.42%

3 Numerical and Experimental Results

In this paper, an actual two-screw workboat is used as a computing sample. Table 2 shows the related data of the hull geometry.

Table 2: Related data of the workboat

Ship length (LWL) in meter	41
Breadth in meter	8
Draft in meter	1.61
Block coefficient (C_B)	0.66
Midship section coefficient (C_M)	0.88
Waterplane area coefficient (C_{WP})	0.85
Duct length in meter	0.675
No. of screw	2

3.1 Resistance Calculation

3.1.1 Hull Viscous Resistance (R_V)

This section presents the computation of a single bare hull. The viscous resistance induced by the fluid viscosity, when the hull is advancing, is calculated by the CFD. The computational settings for the bare hull without a free surface are shown in Table 3. As indicated in Table 3, the inflow velocity is 14 (Kn) and the turbulent model is SST $k - \omega$. Fig. 2 presents the flow field of a bare hull without the free surface effect. As shown in Fig. 2, the distance from the inlet to the bow is twice as large as the ship's length (LWL), the distance from the stern to the outlet is three times the ship's length, the distance from the hull to the slip wall is twice if the ship's length, and the ship's length is 41 m. The boundary condition of the inlet is the velocity inlet, that of the outlet is the pressure outlet, and that of the water plane is set as symmetry. The grid independence analysis is shown in Fig. 3. As can be seen, the resistance difference between 7 million meshes and 8.19 million meshes is 0.3%. Thus, 7 million meshes are used in this study to save the computing time, and the obtained viscous resistance is 44917 N.

Table 3: Computational settings for bare-hull resistance without a free surface

Ship speed (Kn)	14
Turbulent model	SST $k - \omega$
No. of meshes	6948834
Computing status	Steady

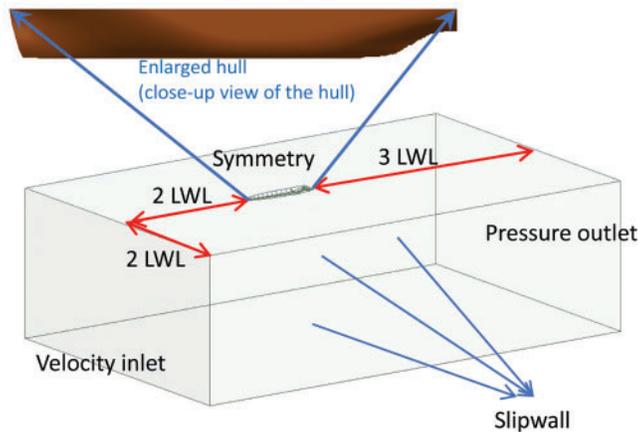


Figure 2: The flow field of a bare hull without the free surface effect

3.1.2 Wave-making Resistance (R_W)

This section calculates the bare-hull resistance with the free surface effect. The result is then compared with the result in Section 3.1.1 to obtain the wave-making resistance (R_W).

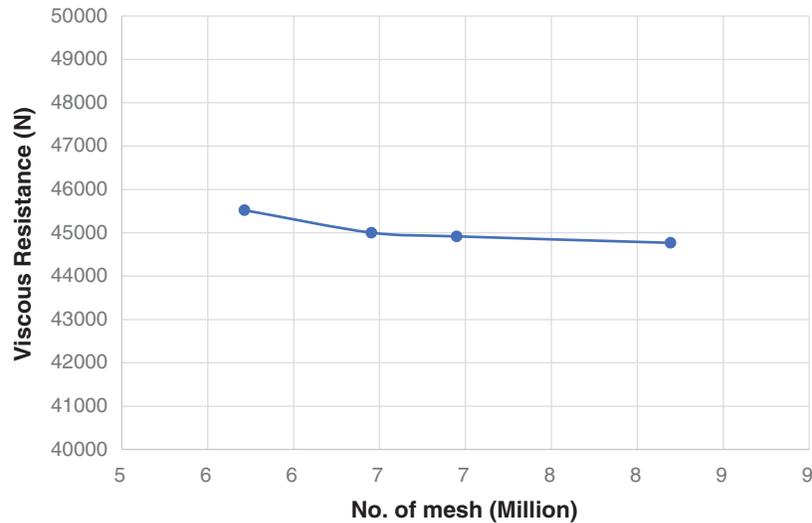


Figure 3: Grid independence analysis of the bare-hull resistance without a free surface

The computational settings for a bare hull with the free surface effect are shown in [Table 4](#). According to the table, the inflow velocity is 14 (Kn) and the turbulent model is SST $k - \omega$. The boundary conditions are listed in [Table 5](#). [Fig. 4](#) presents the flow field of a bare hull with the free surface effect. The trapezoid on the free surface is the region where the grids are refined. In order to shorten the computational time, only half of the hull is calculated and the symmetric boundary condition is applied. The resistance value of the other half of the hull can be obtained by the symmetric boundary condition. The Volume of Fluid model is used to simulate the behavior of waves. The height of the free surface is shown in [Fig. 5](#), and the water surface elevation is set as 1.61 m. [Fig. 6](#) presents the grid independence analysis of the bare-hull resistance with a free surface effect. The resistance difference of the half hull between 4 million meshes and 6 million meshes is less than 1%; it is within the acceptable range. Therefore, 4 million meshes are used in this study for the half hull to obtain satisfactory results and shorten the computational time. In order to describe the wave phenomenon clearly, the Cut-cell meshes are used. The grid arrangements are shown in [Fig. 7](#). The numerical beach is activated at the boundary to avoid the wave reflection at the boundary. The computation in steady state is conducted first to obtain appropriate initial values, and then the transient state is applied. The time step for the transient state is 0.1 sec, and after 4500-step iterations, the resistance value is nearly stable, and there is no obvious jitter. As a result, the average resistance of the half hull from the 4000th step to the 5000th steps is 37.45 KN; thus, the resistance of the entire bare ship with the free surface is 74.9 KN (see [Fig. 8](#)). [Table 6](#) compares the experimental result of the bare-hull resistance with the CFD result. As can be seen, the total resistance difference is 4%; it is within the acceptable range.

[Fig. 9](#) is the distribution of the wave height. In [Fig. 9a](#), the maximum pressure occurs at the bow stagnation point. Consequently, the waves begin to form near the bow, and a wave crest is generated. The top view of the wave height on the free surface is shown in [Fig. 9b](#). The distribution in [Fig. 9b](#) is reasonably similar to the Kelvin wake distribution.

Table 4: Computational settings for the bare-hull resistance with the free surface effect

Ship speed (Kn)	14
Turbulent model	SST $k-\omega$
No. of meshes	4316752
Time step (s)	0.1
Computing status	Transient

Table 5: Boundary conditions for the bare-hull resistance with the free surface effect

	Boundary condition
Inlet	Pressure inlet
Outlet	Pressure outlet
Side plane	Symmetry

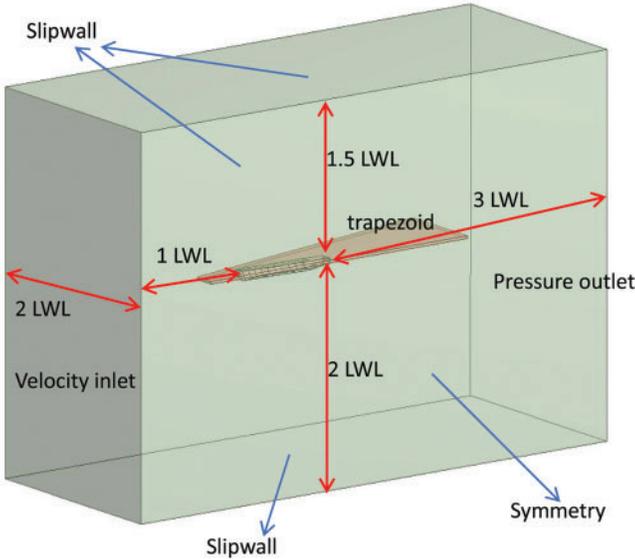


Figure 4: The computing field of the bare hull with the free surface effect

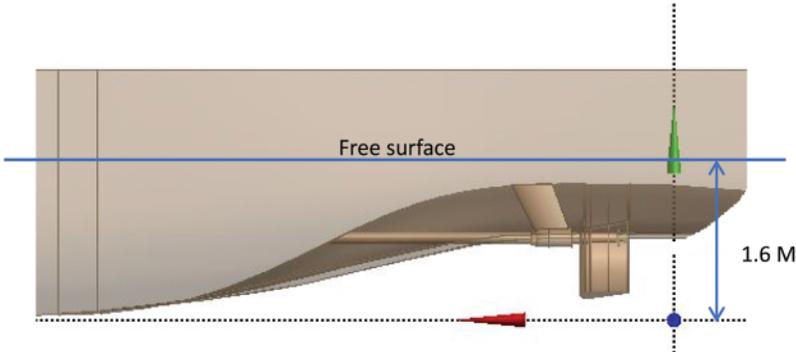


Figure 5: Schematic diagram of the height for the free surface

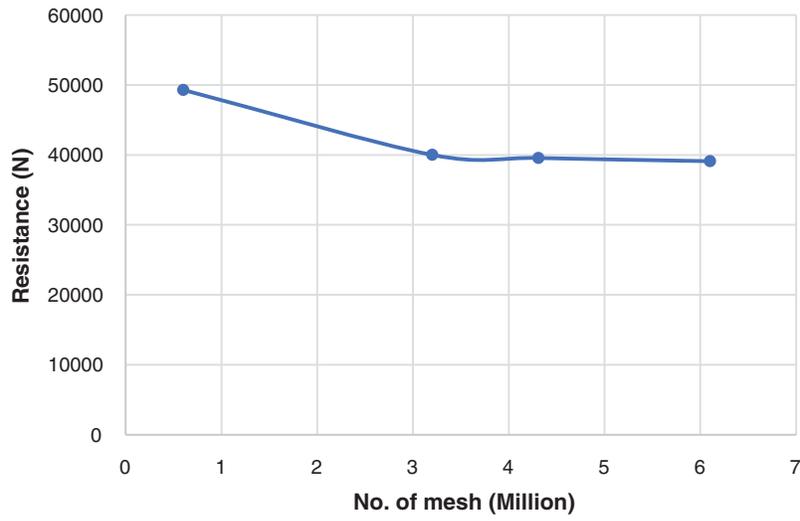


Figure 6: Grid independence analysis of the half bare-hull resistance with the free surface effect

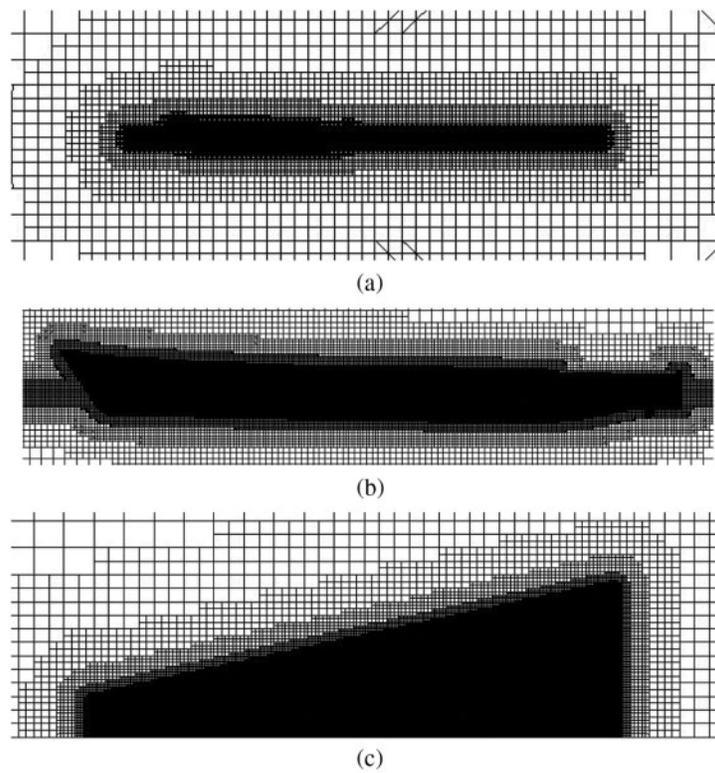


Figure 7: Mesh layout for the flow field of a bare hull with the free surface (a) lateral view of the meshes (b) lateral enlarged view of the meshes (c) top enlarged view of the meshes on the free surface

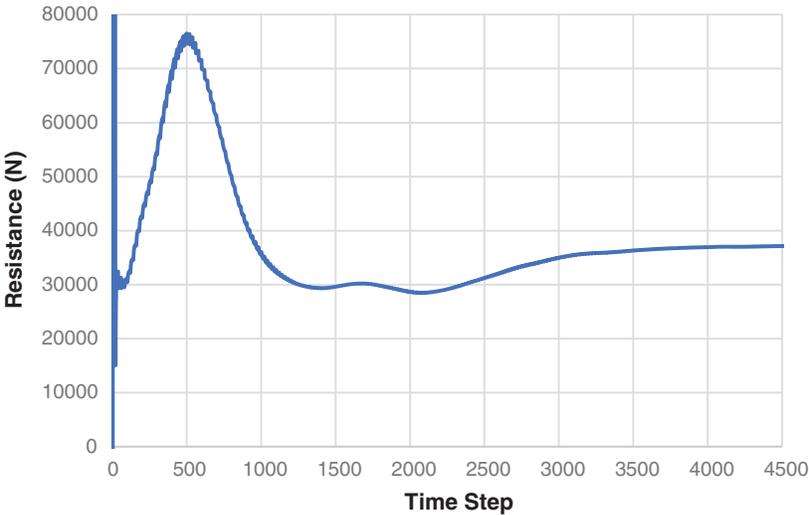


Figure 8: The computing process for the bare-hull resistance with the free surface effect

Table 6: The comparison between the experimental and CFD results of the bare-hull resistance with the free surface

Vs (Kn)	Total bare-hull resistance (KN)	
	By experiment	By CFD
14.0	72	74.9
Error (%)	4%	

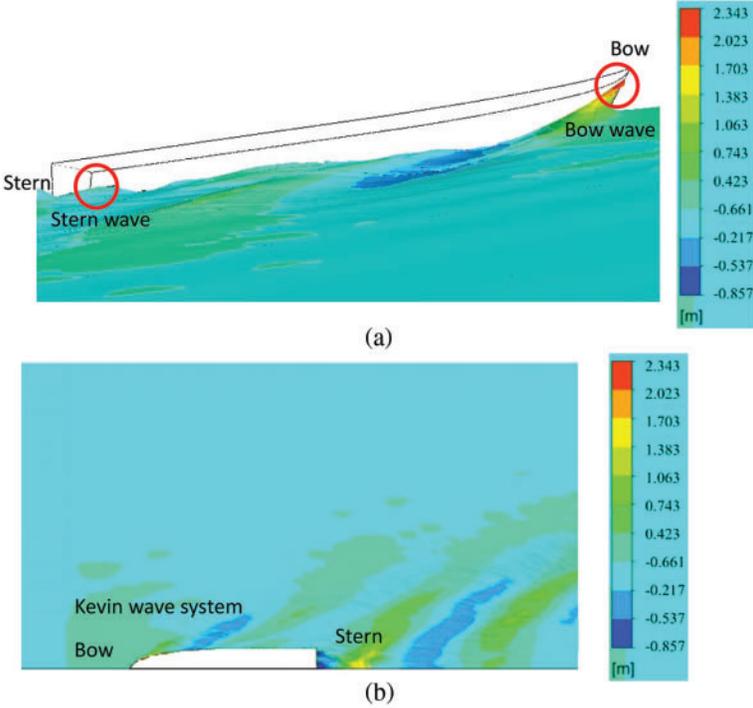


Figure 9: Schematic diagram of the wave height for the free surface (a) lateral view of the wave height of the free surface (b) top view of the wave height of the free surface

3.2 Propeller Performance Curve (K–J Chart)

In addition to the bare-hull resistances described in Section 3.1, the performance curve of the propeller (K–J chart) should be also calculated before the self-propulsion simulation. Table 7 lists the dimensions of the prototype duct and prototype propeller that are used in this paper. Fig. 10 presents the geometries of the prototype propeller and duct. Four combinations of the self-propulsion simulation are calculated in this paper, as shown in Table 8. When the hull is equipped with different versions of the duct and propeller, the thrust and resistance change, and the related propulsion factors have different behaviors.

Table 7: Dimensions of the prototype duct and prototype propeller

Duct length (m)	0.675
Duct inner diameter (m)	1.37
Propeller diameter in meter	1.36
Propeller pitch ratio (P/D) at 0.7R	0.82
Propeller area ratio (E.A.R)	1.05
No. of propeller blades	5

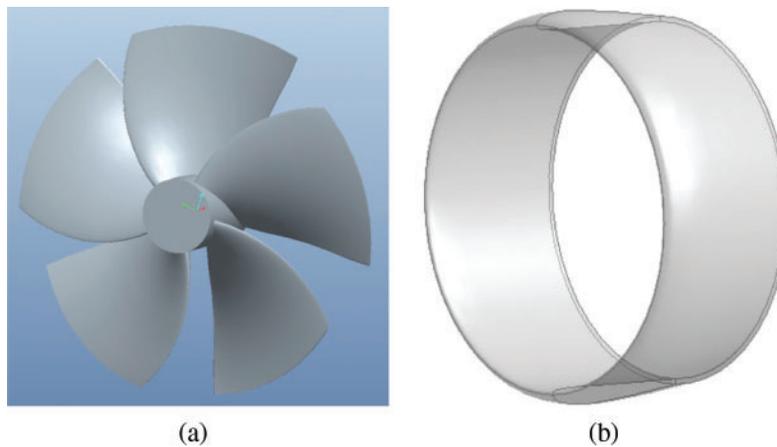


Figure 10: Geometries of the prototype propeller and duct (a) prototype propeller (b) prototype duct

Table 8: Four combinations of the self-propulsion simulations

Computational combination	Duct	Propeller
Combination 1	Prototype duct	Prototype propeller
Combination 2	Prototype duct	Pitch ratio of the prototype propeller $\times 1.05$
Combination 3	Shortened duct (prototype duct is shortened by 53 mm)	Pitch ratio of the prototype propeller $\times 1.04$
Combination 4	Shortened duct (prototype duct is shortened by 53 mm)	1. Pitch ratio of the prototype propeller $\times 1.08$, 2. Camber ratio of the prototype propeller $\times 0.9$

Figs. 11–14 show the K–J charts of the propellers for the four combinations listed in Table 8. For the prototype propeller, the comparison between the predicted values by the CFD and the experimental results is shown in Fig. 11. As can be seen, the calculated results are highly consistent with the experimental values. For the computing sample in this paper, the J_A value of the self-propulsion balance point is between 0.6 and 0.7. The difference between the computed and the experimental results for $J_A = 0.6\sim 0.7$ is very small.

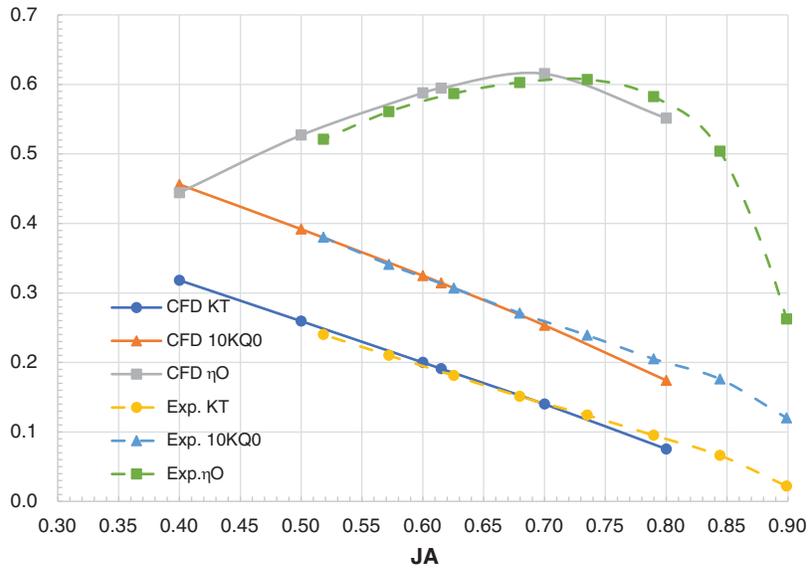


Figure 11: The K–J chart of Combination 1

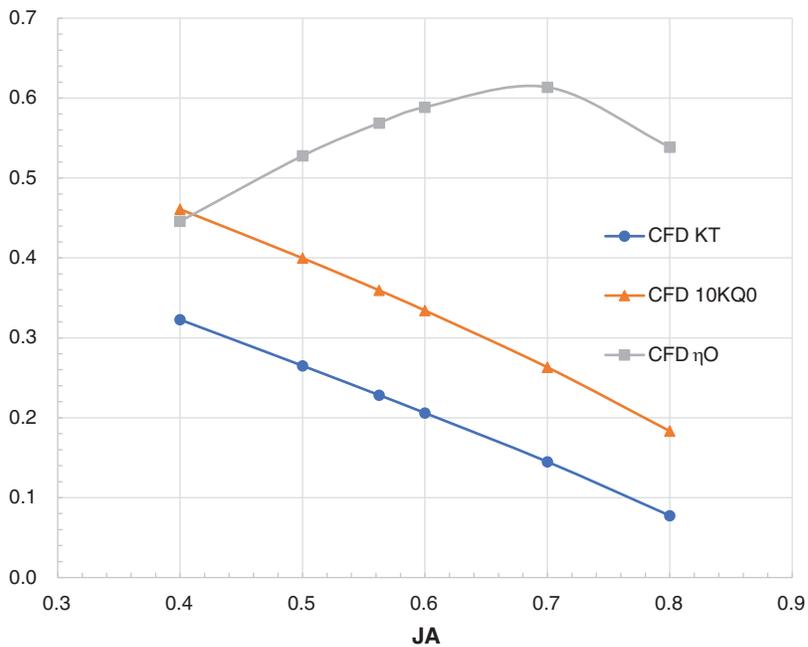


Figure 12: The K–J chart of Combination 2

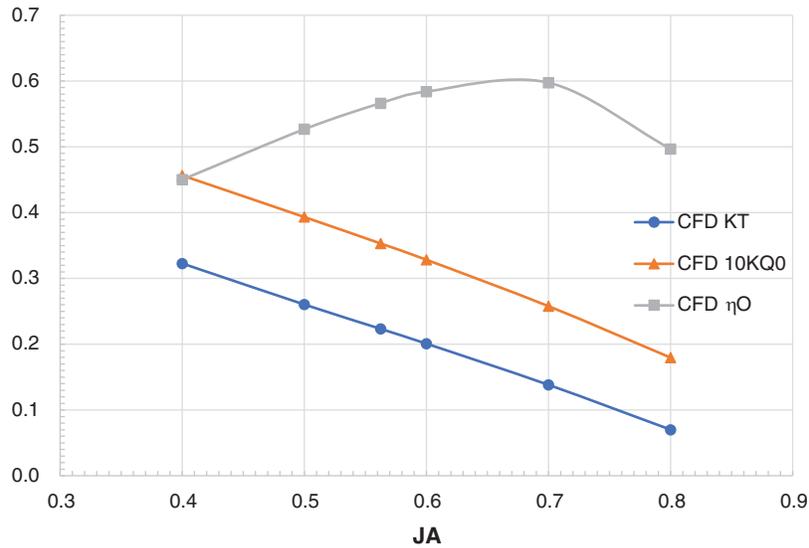


Figure 13: The K–J chart of Combination 3

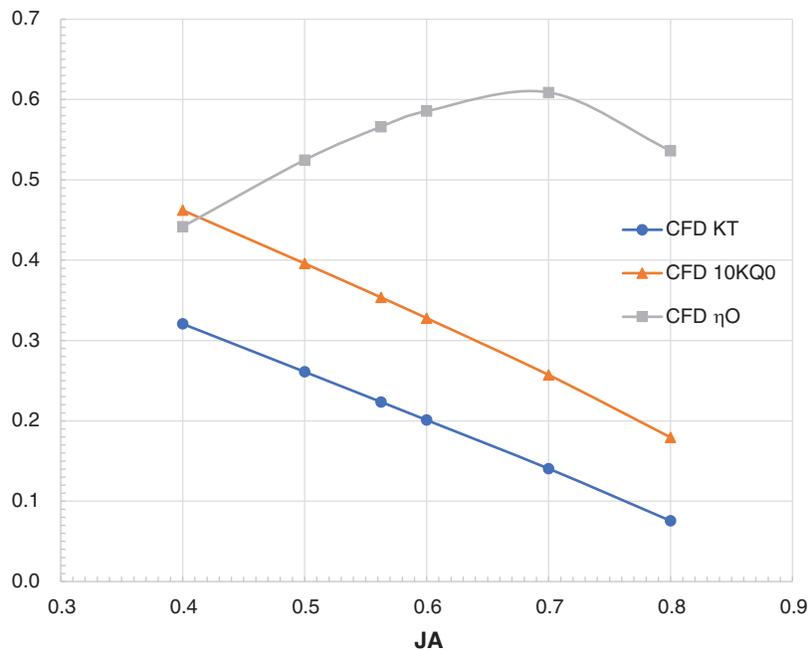


Figure 14: The K–J chart of Combination 4

3.3 Self-Propulsion Simulation

The setting of the computing domains and the mesh generation for the self-propulsion simulations in this paper are based on Kao et al. [10]. Kao et al. [10] applied the CFD for the self-propulsion simulation of the DARPA SUBOFF submarine, and compared the results with the measured results shown in Vaz et al. [6]. The maximum different compared by Vaz et al. [6] was only 3.5%. In addition, the numerical results in Sections 3.1 and 3.2 are compared with

experimental values, and the differences are within the acceptable range. Therefore, the correctness of the self-propulsion simulation results in this section can be indirectly confirmed.

Table 9 lists the CFD settings of the self-propulsion simulation. Fig. 15 presents the computing domains of the self-propulsion simulation. As shown in Fig. 15, the number of the meshes in the internal domain is about 4.8 million, and that in the external domain is about 7.2 million. On the external-domain inlet, the steady uniform velocity, V_s , is imposed, the external-domain outlet is set as the pressure outlet, and the free surface is imposed by the symmetry boundary condition. The propeller speed is 10.195 RPS, and the inflow velocity is 14 (Kn). The inclined shaft angle of the propeller is 0.94° .

In Fig. 16, the meshes of the hull, propeller surface and internal domain are refined. Fig. 17 shows the grid independence analysis. The relative difference between the hull resistances (R) calculated by 12 million meshes and 14 million meshes is 1.18%; the relative difference between the propeller thrusts calculated by 12 million meshes and 14 million meshes is 0.1%. Since both differences are within the acceptable range, 12 million meshes are used for the self-propulsion simulations in this paper.

Table 9: The computing settings for the self-propulsion simulation without the free surface

Ship speed (Kn)	14
Propeller rotating speed (RPS)	10.195
Turbulent model	SST $k - \omega$
No. of meshes	12000000
Time step (s)	0.0016
Computing status	Transient

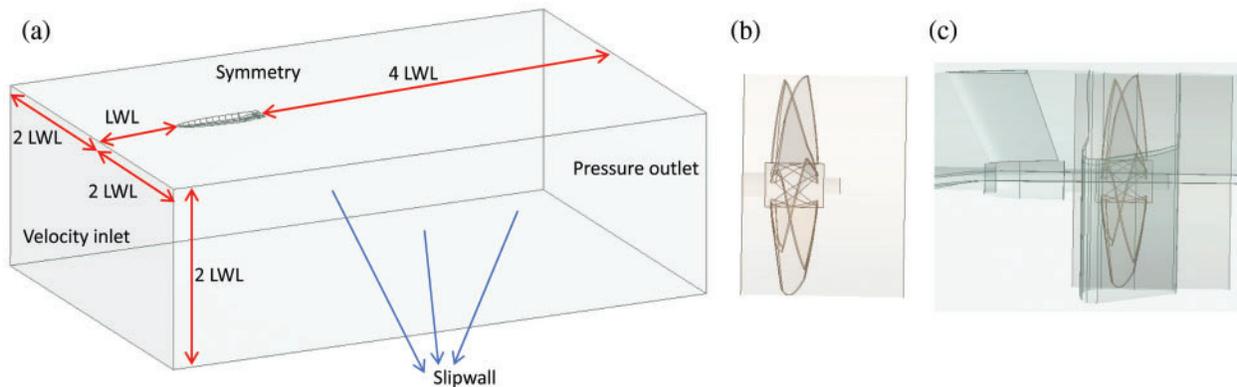


Figure 15: The computing domains of the self-propulsion simulation (a) computing domain for the self-propulsion simulation-external field (b) computing domain for the self-propulsion simulation-internal field (c) enlarged view of stern after combining the internal and external fields

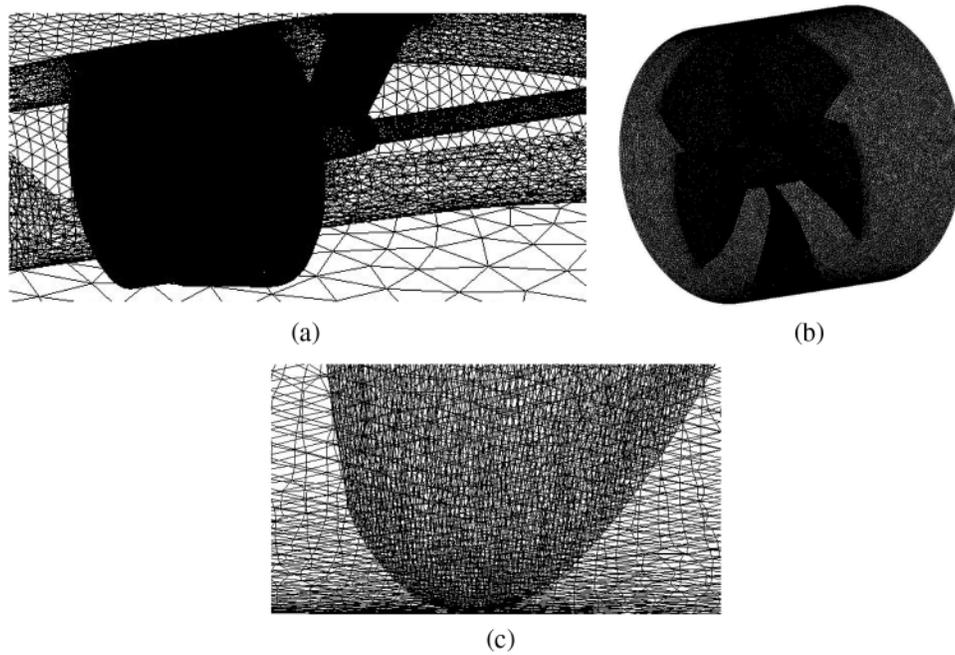


Figure 16: Mesh layout for the self-propulsion simulation (a) enlarged view of the mesh layout near the stern for the self-propulsion simulation (b) enlarged view of the mesh layout in the internal domain for the self-propulsion simulation (c) enlarged view of the mesh layout near the blade tip for the self-propulsion simulation

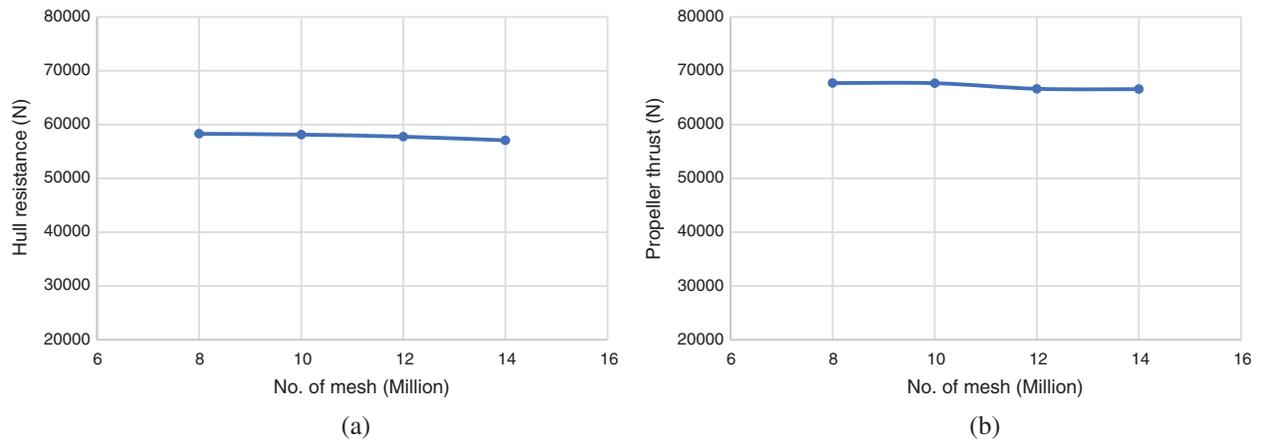


Figure 17: Grid independence analysis of the self-propulsion simulation (a) hull resistance (R) (b) propeller thrust (T)

3.3.1 Results of Combination 1

The results of Combination 1 are listed in Table 10. Table 10 indicates that the calculated hull efficiency is relatively low because $(1 - w)$ is too large (i.e., the flow acceleration due to the duct is too significant). Another reason is that the suction effect of the operating propeller on the hull is too large (i.e., the value of $(1 - t)$ is too small). Therefore, the pitch of the prototype propeller in Combination 2 is increased by 5% to increase the thrust coefficient (K_T). When the thrust coefficient (K_T) is increased, the advanced ratio (J_A) and $(1 - w)$ should be decreased (i.e., the propeller loading is increased). Fig. 18 shows the pressures near the stern. When the operating propeller is not included (Fig. 18a), the pressure on the zone marked 1 is 6000–8000 Pa. However, as the operating propeller is considered, the suction effect causes low pressure at the zone marked 3, and the pressure drops to 4000–6000 Pa (i.e., the resistance due to the propeller suction effect is increased). At the zone marked 2, the diffusion effect decreases the flow velocity, thus the pressure is increased. At the zone marked 4, the suction side and the pressure side of the operating propeller result in the apparent pressure junction.

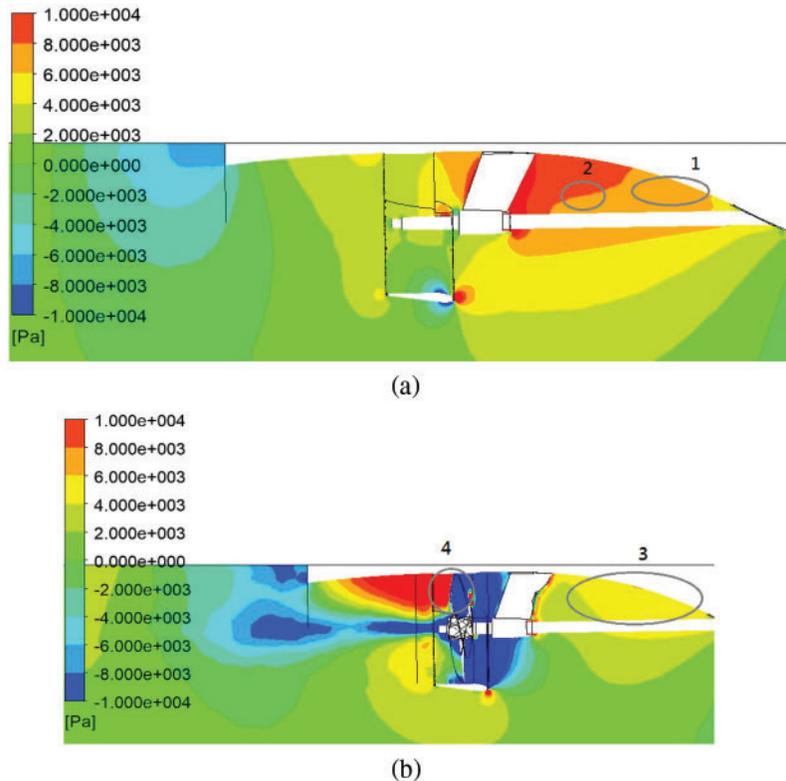


Figure 18: Pressure distribution near the stern for Combination 1. (a) pressure distribution near the stern without an operating propeller (b) pressure distribution near the stern with an operating propeller

Table 10: Results of the self-propulsion simulation for Combination 1

Input	Vs (kn)	14
	Vs (m/s)	7.2016
	D (m)	1.36
	Propeller rotating speed (rps)	10.195
CFD results	$R_V + R_W$ (N)	74917
	T(N)	59276
	R_{SM} (N)	23750
	R_S	19882
	K_T	0.162639
	Q (N.m)	13216
	Self-propulsion results	J_A (KT interpolated K–J chart)
KQ_0 (KT interpolated K–J chart)		0.0265
$1 - w$		1.213
$1 - t$		0.832
Q_0 (N.m)		12789.26
η_R		0.968
η_H		0.686

3.3.2 Results of Combination 2

Table 11 indicates the results for Combination 2. According to Table 11, although the thrust coefficient (K_T) increases from 0.16 to 0.18 by increasing the pitch ratio, the effective wake of Combination 2 is not expectedly reduced. The reason is that the K_T curve in the K–J chart of Combination 2 (Fig. 12) is slightly shifted to the right side as compared with that of Combination 1 (Fig. 12) due to the increment of the propeller pitch. In other words, the interpolated advanced ratio (J_A) is increased in Combination 2; therefore, the effective wake is not decreased as expected and hull efficiency is smaller than that of Combination 1.

Fig. 19 presents the pressure distribution near the stern with the operating propeller. In Fig. 19, the zone circled by the red line reveals that the pressure is obviously raised at the six-o'clock position of the propeller pressure side. The pressure difference between the propeller suction and pressure sides becomes larger as compared with Combination 1; consequently, the thrust coefficient is increased. For effectively deducing the acceleration caused by the duct, the duct length is shortened in Combination 3.

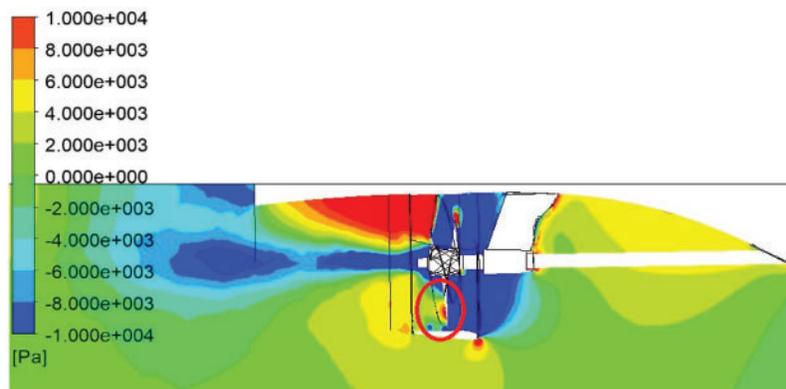
**Figure 19:** Pressure distribution near the stern with an operating propeller for Combination 2

Table 11: Results of the self-propulsion simulation for Combination 2

Input	Vs (kn)	14
	Vs (m/s)	7.2016
	D (m)	1.36
CFD results	Propeller rotating speed (rps)	10.195
	$R_V + R_W$ (N)	74917
	T(N)	67135
	R_{SM} (N)	32500
	R_S	26559
	K_T	0.184202
	Q (N.m)	15101
Self-propulsion results	J_A (K_T interpolated K–J chart)	0.637
	KQ_0 (K_T interpolated K–J chart)	0.0308
	$1 - w$	1.226
	$1 - t$	0.800
	Qo	14864.5
	η_R	0.984
	η_H	0.652

3.3.3 Results of Combination 3

The results of Combination 3 are shown in [Table 12](#). By comparing the results between Combinations 2 ([Table 11](#)) and 3 ([Table 12](#)), it is found that $(1 - w)$ is decreased from 1.226 to 1.186 by shortening the duct length (i.e., the propeller inflow V_A is decreased). The decrease in $(1 - w)$ leads to a rise in the hull efficiency, η_H . Besides, a decline in the inflow velocity V_A results in the less cavitation.

Table 12: Results of the self-propulsion for Combination 3

Input	Vs (kn)	14
	Vs (m/s)	7.2016
	D (m)	1.36
CFD results	Propeller rotating speed (rps)	10.195
	$R_V + R_W$ (N)	74917
	T(N)	69311
	R_{SM} (N)	34700
	R_S	28985
	K_T	0.190173
	Q (N.m)	15015
Self-propulsion results	J_A (K_T interpolated K–J chart)	0.616
	KQ_0 (K_T interpolated K–J chart)	0.0318
	$1 - w$	1.186
	$1 - t$	0.791
	Qo	15347.11
	η_R	1.022
	η_H	0.667

Fig. 20 illustrates the pressure distribution on the leading edge of the foil at $0.8R_p$ (R_p the propeller radius) when the propeller blade is at the six-o'clock position (i.e., at the 180-degree position). In Fig. 20, the high pressure is detected on the suction side of the propeller foil. In other words, there is a negative attacked angle for the blade at the six-o'clock position; thus, the face cavitation appears easily. In Combination 4, the propeller pitch is increased to avoid the negative attacked angle.

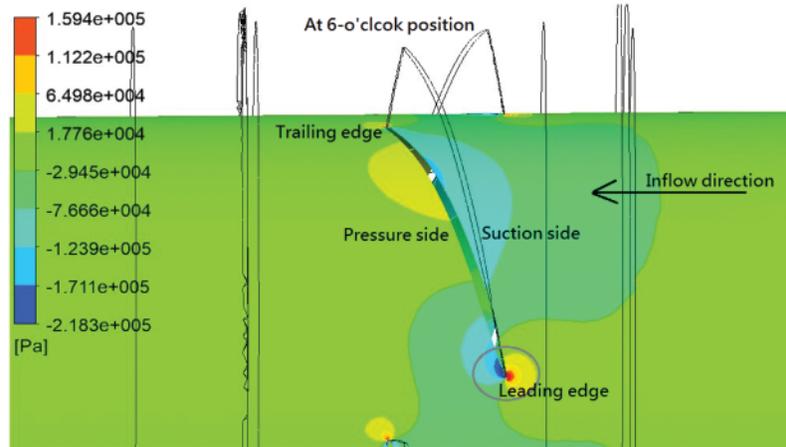


Figure 20: Pressure distribution on the leading edge of the propeller blade for Combination 3 at the six o'clock position

3.3.4 Results of Combination 4

Table 13 presents the results of the self-propulsion simulation for Combination 4. There is only slight difference between the results of Combination 3 and Combination 4. Fig. 21 presents the pressure distribution on the foil at $0.8R_p$ when the propeller blade is at the six-o'clock position (180° position). As indicated in the circled zone, the negative attacked angle still exists. Fig. 22 illustrates the pressure coefficient ($-C_p$) distribution of the foil at $0.8R_p$ as the blade is at different angles (0° and 180°). $-C_p$ is defined by Eq. (8).

$$-C_p = \frac{P_0 + \rho gh - P_L}{\frac{1}{2}\rho V_S^2} \quad (8)$$

Here P_0 is the atmospheric pressure, V_S is the ship speed, P_L is the local pressure on the foil, g is the gravity acceleration, and h is the shaft immersion. For the blade, the inflow velocity at the 180° position is higher than that in the 0° position; thus, the negative attacked angle appears easily at the 180° position. Consequently, the $-C_p$ value on the suction side of the leading edge at 180° position is decreased. For the same reason, the $-C_p$ on the suction side of the mid chord at 180° position is expected to decrease together; however, the computational results do not meet the expectation. This is because the acceleration quality of the duct is still too strong, so that the suction side of the mid-chord at the 180° position becomes a low-pressure area, and the $-C_p$ on this area is still larger than that at the 0° position.

Table 13: Results of the self-propulsion simulation for Combination 4

Input	V_s (kn)	14
	V_s (m/s)	7.2016
	D (m)	1.36
	Propeller rotating speed (rps)	10.195
CFD results	$R_V + R_W$ (N)	74917
	T(N)	68907
	R_{SM} (N)	37200
	R_S	25669
	K_T	0.189064
	Q (N.m)	15352
	J_A (K_T interpolated K-J chart)	0.621
Self-propulsion results	KQ_0 (K_T interpolated K-J chart)	0.0315
	$1 - w$	1.196
	$1 - t$	0.814
	Q_0	15202.33
	η_R	0.990
	η_H	0.680

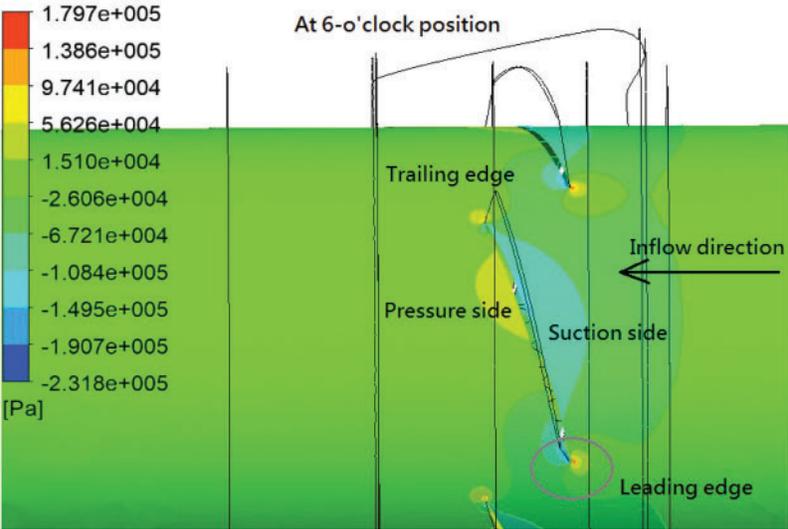


Figure 21: Pressure distribution on the leading edge for the propeller blade for Combination 4 at the six o'clock position

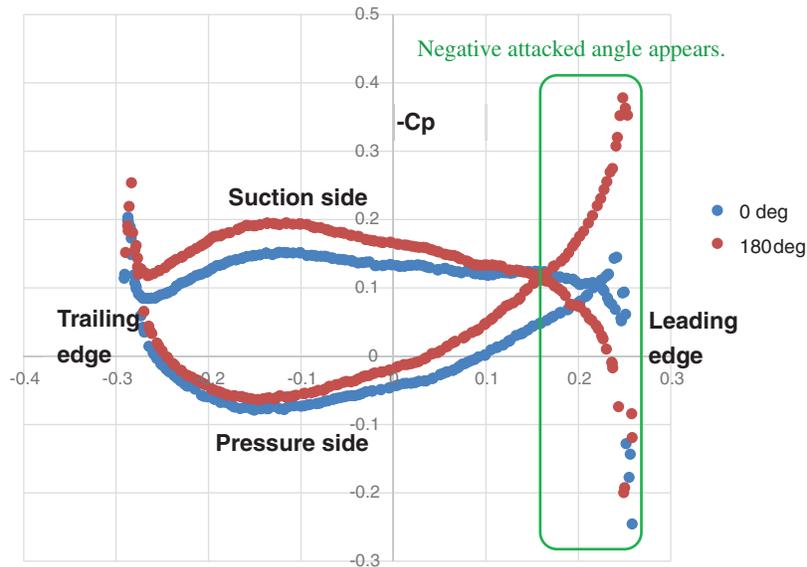


Figure 22: Pressure coefficient ($-C_p$) distribution on the foil of the $0.8R_p$ radial position at different blade angles

4 Conclusion

This study employed CFD to conduct the self-propulsion simulation of a ducted propeller with four combinations of different propeller geometries and duct lengths. The conclusions are as follows:

- (1) In order to compute the wave-making resistance correctly, the mesh needs to be refined on the free surface. The result is consistent with the Kelvin wave system. The maximum pressure occurs at the stagnation point of the bow. Consequently, the waves begin to form, and a wave crest is generated. The results are consistent with the physical phenomenon.
- (2) The predicted propeller performance curve (K–J chart) and total bare-hull resistance are compared with the experimental results. The comparison comparisons are in good agreement. In other words, the correctness of the self-propulsion simulation is confirmed indirectly.
- (3) Four combinations with different parameters including the propeller pitch ratio, propeller camber ratio, and the duct length are computed in this paper. These parameters are changed according to the detected results of the fluid acceleration which will affect the propulsive efficiency. As the acceleration is too significant (i.e., the hull efficiency is relatively low, and the face cavitation appears), three strategies are provided to solve this problem according to the propulsive theory. The first strategy is to increase the propeller pitch ratio, the second strategy is to shorten the duct length, and the last strategy is to increase the propeller camber ratio. The computed results show that the most effective way to improve the problem caused by the fluid acceleration is to shorten the duct length for the ducted propulsion system.
- (4) As the effective wake ($1 - w$) is too large, the hull efficiency is relatively low. In Combination 2, the pitch of the prototype propeller is increased in order to increase the propeller thrust coefficient and consequently reduce the effective wake. However, the increasing pitch ratio does not work. The K_T curve in the K–J chart is slightly shifted to the right side

as the propeller pitch is increased. The interpolated advanced ratio (J_A) and the effective wake ($1 - w$) are both therefore increased, and hull efficiency is brought down.

- (5) The duct is shortened to reduce the duct acceleration effect. As a result, ($1 - w$) decreases from 1.226 to 1.186 in Combination 3. The hull efficiency is also improved significantly.
- (6) The propeller inflow is remarkably accelerated by the duct. The strong acceleration causes a negative attacked angle for the propeller; thus, the propeller face cavitation can be detected easily. In the future, it is recommended to redesign the duct geometry for reducing the flow acceleration.

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References

1. Sung, C. H., Griffin, M., Tsai, J., Huang, T. (1993). Incompressible flow computation of forces and moments on bodies of revolution at incidence. *31st Aerospace Sciences Meeting*, pp. 787. Reno, USA.
2. Kim, W. J., Kim, D. H., Van, S. H. (2002). Computational study on turbulent flows around modern tanker hull forms. *International Journal for Numerical Methods in Fluids*, 38(4), 377–406. DOI 10.1002/(ISSN)1097-0363.
3. Huuva, T. (2008). *Large eddy simulation of cavitating and non-cavitating flow*. Gothenburg, Sweden: Chalmers University of Technology.
4. Yang, J., Rhee, S. H., Kim, H. (2009). Propulsive performance of a tanker hull form in damaged conditions. *Ocean Engineering*, 36(2), 133–144. DOI 10.1016/j.oceaneng.2008.09.012.
5. Toxopeus, S., Vaz, G. (2009). Calculation of current or manoeuvring forces using a viscous-flow solver. *International Conference on Offshore Mechanics and Arctic Engineering*, 43451, 717–728. DOI 10.1115/OMAE2009-79782.
6. Vaz, G., Toxopeus, S., Holmes, S. (2010). Calculation of manoeuvring forces on submarines using two viscous-flow solvers. *International Conference on Offshore Mechanics and Arctic Engineering*, vol. 49149, pp. 621–633. DOI 10.1115/OMAE2010-20373.
7. Vaz, G., Jaouen, F., Hoekstra, M. (2009). Free-surface viscous flow computations: Validation of URANS code FRESCO. *International Conference on Offshore Mechanics and Arctic Engineering*, vol. 43451, pp. 425–437. DOI 10.1115/OMAE2009-79398.
8. Seo, J. H., Seol, D. M., Lee, J. H., Rhee, S. H. (2010). Flexible CFD meshing strategy for prediction of ship resistance and propulsion performance. *International Journal of Naval Architecture and Ocean Engineering*, 2(3), 139–145. DOI 10.2478/IJNAOE-2013-0030.
9. Korkmaz, K. B. (2015). *CFD predictions of resistance and propulsion for the JAPAN Bulk Carrier (JBC) with and without an energy saving device (Master's Thesis)*. Chalmers University of Technology.
10. Kao, J. H., Lin, Y. J., Chin, S. S., Chang, F. N., Tsai, Y. H. et al. (2019). Computational analysis of underwater noise of DARPA SUBOFF submarine. *Annual Meeting of Society of Naval and Marine Engineering*. Kaohsiung, Taiwan.
11. Yu, L., Greve, M., Druckenbrod, M., Abdel-Maksoud, M. (2013). Numerical analysis of ducted propeller performance under open water test condition. *Journal of Marine Science and Technology*, 18(3), 381–394. DOI 10.1007/s00773-013-0215-4.
12. Yongle, D., Baowei, S., Peng, W. (2015). Numerical investigation of tip clearance effects on the performance of ducted propeller. *International Journal of Naval Architecture and Ocean Engineering*, 7(5), 795–804. DOI 10.1515/ijnaoe-2015-0056.

13. Gao, D., Su, J., Li, M., Hua, H. (2017). Flow-induced vibro-acoustic analysis of a ducted propeller propulsion system. *INTER-NOISE and NOISE-CON Congress and Conference Proceedings*, vol. 255, no. 6, pp. 1884–1892. Institute of Noise Control Engineering.
14. Gaggero, S., Villa, D., Tani, G., Viviani, M., Bertetta, D. (2017). Design of ducted propeller nozzles through a RANSE-based optimization approach. *Ocean Engineering*, 145(1), 444–463. DOI 10.1016/j.oceaneng.2017.09.037.
15. Gong, J., Guo, C. Y., Zhao, D. G., Wu, T. C., Song, K. W. (2018). A comparative DES study of wake vortex evolution for ducted and non-ducted propellers. *Ocean Engineering*, 160(1), 78–93. DOI 10.1016/j.oceaneng.2018.04.054.
16. Andersson, J., Eslamdoost, A., Vikström, M., Bensow, R. E. (2018). Energy balance analysis of model-scale vessel with open and ducted propeller configuration. *Ocean Engineering*, 167(1), 369–379. DOI 10.1016/j.oceaneng.2018.08.047.
17. Gong, J., Guo, C. Y., Phan-Thien, N., Khoo, B. C. (2020). Hydrodynamic loads and wake dynamics of ducted propeller in oblique flow conditions. *Ships and Offshore Structures*, 15(6), 645–660. DOI 10.1080/17445302.2019.1663664.
18. Kinnas, S. A., Fan, H., Tian, Y. (2015). A panel method with a full wake alignment model for the prediction of the performance of ducted propellers. *Journal of Ship Research*, 59(3), 246–257. DOI 10.5957/jsr.2015.59.3.246.
19. Bosschers, J., Willemsen, C., Peddle, A., Rijpkema, D. (2015). Analysis of ducted propellers by combining potential flow and RANS methods. *Fourth International Symposium on Marine Propulsors*, 39–648.
20. Bontempo, R., Cardone, M., Manna, M. (2016). Performance analysis of ducted marine propellers. Part I–Decelerating duct. *Applied Ocean Research*, 58, 322–330. DOI 10.1016/j.apor.2015.10.005.
21. Bontempo, R., Manna, M. (2018). Performance analysis of ducted marine propellers. Part II–Accelerating duct. *Applied Ocean Research*, 75, 153–164. DOI 10.1016/j.apor.2018.03.005.
22. Du, W., Kinnas, S. A. (2019). Applying 3D viscous/Inviscid interaction method to predict flows around wings and propellers. *SNAME 24th Offshore Symposium*, OnePetro. Houston, Texas, USA.
23. Villa, D., Gaggero, S., Tani, G., Viviani, M. (2020). Numerical and experimental comparison of ducted and non-ducted propellers. *Journal of Marine Science and Engineering*, 8(4), 257. DOI 10.3390/jmse8040257.
24. Delouei, A. A., Nazari, M., Kayhani, M. H., Ahmadi, G. (2017). Direct-forcing immersed boundary–non-newtonian lattice boltzmann method for transient non-isothermal sedimentation. *Journal of Aerosol Science*, 104, 106–122. DOI 10.1016/j.jaerosci.2016.09.002.
25. Karimnejad, S., Delouei, A. A., Nazari, M., Shahmardan, M. M., Rashidi, M. M. et al. (2019). Immersed boundary–thermal lattice boltzmann method for the moving simulation of non-isothermal elliptical particles. *Journal of Thermal Analysis and Calorimetry*, 138(6), 4003–4017. DOI 10.1007/s10973-019-08329-y.
26. Afra, B., Delouei, A. A., Mostafavi, M., Tarokh, A. (2021). Fluid-structure interaction for the flexible filament’s propulsion hanging in the free stream. *Journal of Molecular Liquids*, 323(3), 114941. DOI 10.1016/j.molliq.2020.114941.
27. Patankar, S. V., Spalding, D. B. (1983). A calculation procedure for heat, mass and momentum transfer in three-dimensional parabolic flows. *Numerical prediction of flow, heat transfer, turbulence and combustion*. 15(10), 54–73. Pergamon. DOI 10.1016/B978-0-08-030937-8.50013-1.
28. HYKAT HSVA (2016). *Calm water model tests for a container vessel*, Hamburg: Hamburgische Schiffbau Versuchsanstalt.