



**ARTICLE**

# Bioconvective Hybrid Flow with Microorganisms Migration and Buongiorno's Model under Convective Condition

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Received: 21 July 2023 Accepted: 25 October 2023 Published: 20 May 2024

## ABSTRACT

Heat transfer improves significantly when the working fluid has high thermal conductivity. Heat transfer can be found in fields such as food processing, solar through collectors, and drug delivery. Considering this notable fact, this work is focused on investigating the bio-convection-enhanced heat transfer in the existence of convective boundary conditions in the flow of hybrid nanofluid across a stretching surface. Buongiorno fluid model with hybrid nanoparticles has been employed along the swimming microorganisms to investigate the mixture base working fluid. The developed nonlinear flow governing equations have been tackled numerically with the help of the bvp4c. The effects of relevant parameters on the flow dynamic have been portrayed in a graphical representation. The velocity profile decreases by raising the levels of buoyancy ratio and mixed convection in the range of  $0.1 < \lambda \leq 0.3$ . It has been discovered that when bioconvection levels rise, motile microbe migration abruptly slows, which results in a decrease in fluid acceleration. The concentration of fluid flow declined for the Lewis number, but the opposite trend has been observed for the elastic parameter, thermophoresis parameter, and buoyancy ratio. With rising values of Brownian motion and thermophoretic diffusion, the surface drag and Nusselt number decrease significantly. Whereas, the opposite trend has been observed when the values of the thermal Biot number, Prandtl number and buoyancy ratio are enhanced. Additionally, data from this study have been validated by comparison with those that have previously been published, and an appropriate rate of agreement has been observed.

## KEYWORDS

Prandtl hybrid nanofluid; mixed convection; stretched sheet; bioconvection; motile microorganisms

## Nomenclature

$C$	Concentration profile
$N$	Motile microorganism's density
$T$	Temperature profile
$N_w$	Density at the wall of Motile microorganism
$T_w$	Surface temperature



$T_\infty$	Ambient temperature
$Nr$	Parameter of Buoyancy ratio
$D_2$	Mass diffusion coefficient
$C_{fx}$	Skin friction along x-direction
$N_\infty$	Ambient density of microorganisms
$D_1$	Temperature diffusion coefficient
$C_\infty$	Ambient concentration
$Re_x$	Reynold's number
$D_B$	Coefficient of Brownian diffusion
$C_w$	Surface concentration
$Nc$	Rayleigh number Bioconvection
$Pr$	Prandtl number
$Lb$	Parameter of Bioconvection Lewis number
$D_m$	Microorganism coefficient
$Bi$	The thermal Biot number
$Pe$	Peclet number
$Le$	Lewis number
$(u,v)$ ( $\text{ms}^{-1}$ )	Velocity components

### **Greek Symbols**

$u_w$ ( $\text{Nsm}^{-2}$ )	Stretching sheet's velocity
$\alpha_1$	Parameter of Prandtl fluid
$\mu_f$ ( $\text{Nsm}^{-2}$ )	Base fluid viscosity
$\omega$	Bio-convection constant
$\alpha$	Angle of inclination
$\alpha_{hmf}$ ( $\text{m}^{-2} \text{s}^{-1}$ )	Thermal diffusivity of HNF
$\lambda$	Mixed convection parameter
$\tau$ ( $\text{Nm}^{-2}$ )	Specific heat capacity ratio
$\alpha_2$	Elastic parameter
$\rho_p$ ( $\text{kgm}^{-3}$ )	Nanofluid density
$\omega_c$ ( $\text{ms}^{-1}$ )	Maximum cell velocity
$\rho_m$ ( $\text{kgm}^{-3}$ )	Density motile microorganism
$\eta$	Similarity variable
$\mu_{hmf}(\text{Nsm})^{-2}$	Kinematic viscosity of HNF
$\rho_f$ ( $\text{kgm}^{-3}$ )	Density of nanofluid
$g$ ( $\text{m}^{-2}\text{s}^{-1}$ )	Gravity
$Nu_x$	Nusselt number
HNF	Hybrid nanofluids

## **1 Introduction**

Recently, the new idea of “hybrid nanofluids” has surfaced in which more than one nano-meter nanoparticle is dispersed into the working fluid, and the formed solution is termed a hybrid nanofluid. The importance of improving heat transfer is emphasized by the industry's production sectors, where it finds notable uses. Mainstream fluids, such as water, ethylene glycol (EG), oils, biological fluids, etc., may transmit heat and are employed in many engineering projects, such as heating systems and electronic devices. These base liquids do not transmit heat well. A substance

known as nanofluid, a combination of base liquid and nanoparticles, was created to get a boost of thermal conductivity. By incorporating nano-scale solid particles into convection fluids, Choi et al. [1] first described the concept of nanofluids. Buongiorno [2] discussed the convective heat transfer in the nanofluids using thermophoretic particle diffusion and Brownian motion. With the use of the slip factor, Turkyilmazoglu [3] showed the comparative analysis of single and several phasing streams of nanofluids in radial annuli. Sarada et al. [4] studied the behaviour of a non-Newtonian fluid movement under an extended sheet under global thermal non-equilibrium conditions.

Babu et al. [5] explored non-Newtonian fluid flow across a thin stretched sheet under the influence of cross-diffusion phenomena. Sharma et al. [6] investigated the impact of heat radiation on a two-dimensional non-Newtonian fluid flow over a stretched surface. Ibrahim et al. [7] studied the impact of chemical reactions and radiation on the creation of heat and non-Newtonian fluid flow along a stretched sheet with an uneven depth. The flow of mixed convection in a non-Newtonian fluid flow with entropy and energy activation production was predicted by Ijaz Khan et al. [8]. The effects of activation energy and chemical reactions in viscoelastic liquids traveling over a stretched surface were examined by Ramesh [9]. In order to understand how activation energy affects the nanofluid's radiative stream under a convective boundary situation, Ijaz Khan et al. [10] looked into the issue. Saleem et al. [11] conducted theoretical research on the thermodynamic properties of hybrid nanofluids propelled by cilia. Beg et al. [12] discussed the Von Karman hydro-magnetic flow with heat transfer, joule heating, and viscous dissipation effect. Dinarvand et al. [13] studied the dual solution of the hybrid nanofluids over a thin moving needle. Mansourian et al. [14] examined hybrid nanofluids over non-linearly stretched porous surfaces. Naz et al. [15] investigated the cross fluid with entropy generation and motile microorganisms. Balla et al. [16] elaborated on the bioconvection in microorganisms in a square cavity with a thermal radiation effect. Hussain et al. [17] examined the magnetic hybrid nanofluid flow of carbon nanotube with thermal radiation. Arshad et al. [18] explored hybrid nanofluids between permeable rotating systems numerically. Rao et al. [19] used graphene nanoparticles to optimize linseed biodiesel using response surface methodology.

Bioconvection is the term used to describe the process by which microorganisms randomly travel in single-cell or colony-like shapes. To show the bioconvection phenomenon, Vincent et al. [20] used a floated algae solution. Li et al. [21] discussed bioconvection heat transfer in third-grade fluid using activation energy. Hussain et al. [22] explored the bioconvection effect on hybrid nanofluid with microorganisms Cattaneo-Christov heat flux. Elsebaee et al. [23] studied microorganism flow with tri-hybrid particles using a magnetic field over a slandering sheet. Khashi'ie et al. [24] studied the dual solutions of hybrid nanofluids using microorganisms towards a vertical plate. Sreedevi et al. [25] investigated the Williamson hybrid nanofluid with microorganisms and Cattaneo-Christov heat flux. Xu et al. [26] examined Maxwell fluid with microorganisms within parallel plates.

The heated surface has the potential to improve the heat transfer further over the infinite stretch surfaces. This provides researchers with options, along with higher thermal conductivity of the hybrid nanofluids. We can also employ the convective boundary to boost the heat transfer process in many engineering applications such as MEMS and solar through collectors. In the Newtonian and non-Newtonian flow examination over the stretched surface with different body force effects, many researchers have considered the heated surface while modeling the boundary conditions of the problems. Jusoh et al. [27] discussed heat transfer with viscous dissipation and convective boundary conditions using hybrid nanofluids. Rashid et al. [28] examined hyperbolic tangent fluid over a stretching surface with convective conditions. Waini et al. [29] explored heat transfer over stretching/shrinking surfaces with convective boundaries. Aly et al. [30] magneto-hydrodynamic flow with hybrid nanofluid over stretching/shrinking using heated boundary conditions. Hassan et al. [31] investigated the hybrid

ferrofluid over a stretching surface using the thermal slip condition with a nanoparticle shape effect. Srinivasulu et al. [32] studied inclined magnetization with convective boundary conditions for the non-Newtonian fluid over a stretched surface. Yamda et al. [33] discussed the effective thermal conductivity whereas, Xue [34] examined the thermal conductivity of the carbon nanotube base materials.

In the above literature review on the hybrid nanofluids, it was discovered that there is a lack of information on the dynamics of a mixture base as working fluid over a stretched surface. It is worth mentioning here the convective boundary condition plays a significant role in heat transfer. Additionally, the modified Buongiorno nanofluid model has not been employed with the Prandtl fluid model and microorganism profiles at the same time. Moreover, the magnetite and silver nanoparticles with a mixture base working fluid were not examined with the assumptions mentioned above. Therefore, in this article, we focus on investigating the Bioconvective hybrid nanofluid using a mixture base working fluid over the heated stretched surface, considering the modified Buongiorno nanofluid model with swimming microorganisms. We utilize a suitable similarity transformation to convert the partial differential equations into ordinary differential equations. In order to obtain the necessary numerical solution of modified partial differential equations, Matlab software is used with the bvp4c package. It will be possible to comprehend how emerging parameters affect fluid flows. Dimensionless parameters will be utilized to examine the behaviour of several physical and heat transfer characteristics of the hybrid nanofluid flow. This study focuses on the following research questions:

- How does mixed convection affect the flow of hybrid Prandtl nanofluid through a stretched heated sheet?
- How do the Prandtl number, thermophoresis, and Brownian motion affect the mixture-base hybrid nanofluid in the presence of microorganisms?
- Interpret the manner in which bio-convection affects the motile microbe dispersion.
- How do thermophoretic diffusion, Brownian motion, mixed convection, thermal Biot number, and Prandtl number affect the skin friction and Nusselt number?

## 2 Mathematical Modeling of the Problem

In this section, the flow governing mathematical model has been developed using the following assumptions:

The surface is stretched horizontally with some stretching velocity.

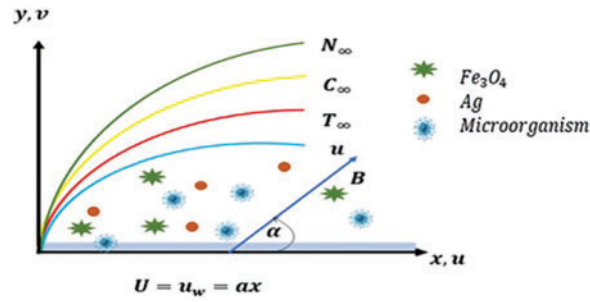
The convective boundary condition is applied at the surface.

Modified Buongiorno nanofluid has been taken into consideration.

Magnetite  $\text{Fe}_2\text{O}_4$  and silver Ag are used as nanoparticles. Mixture base fluid comprising (50%–50%) ethylene glycol  $\text{C}_6\text{H}_6\text{O}_2$  and water  $\text{H}_2\text{O}$  are used to constitute the desired hybrid nanofluid  $\text{Fe}_2\text{O}_4$ – $\text{Ag}/\text{C}_6\text{H}_6\text{O}_2$ – $\text{H}_2\text{O}$ .

### 2.1 Statement of Problem

Let us consider two dimensions: a steady, incompressible flow of hybrid nanofluid over the stretched surface, stretched with some stretching velocity in the x-direction. The flow is restricted to the upper half of the plane, i.e.,  $y > 0$ . Cartesian coordinated system has been taken into account. The surface is heated with applied heating. The surface contains microorganisms. Fig. 1 shows the configuration of the problem.



**Figure 1:** Geometry of problem

Buongiorno [2] presented a nanofluid model using thermophoretic particle diffusion and Brownian motion. We employed its modified form, as discussed by Hussain et al. [22]. Following are the flow governing equations conservation of mass (continuity) (1), conservation of momentum (2), energy Eq. (3), concentration (4), and motile microorganisms Eq. (5), respectively, developed using the above-mentioned flow assumptions.

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0. \tag{1}$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = V_{hmf} \frac{A}{C} \left( \frac{\partial^2 u}{\partial y^2} \right) + V_{hmf} \frac{A}{2C^3} \left( \frac{\partial^2 u}{\partial y^2} \right) \left( \frac{\partial u}{\partial y} \right)^2 + \frac{1}{\rho_{hmf}} [g(T - T_\infty) \rho_f \beta (1 - C_\infty) + g(C_\infty - C) (\rho_p - \rho_f) - (\rho_m - \rho_p) (N_\infty - N) g\gamma'] . \tag{2}$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha_{hmf} \left( \frac{\partial^2 T}{\partial y^2} \right) + \tau \left( \left( \frac{\partial^2 T}{\partial y^2} \right) \frac{D_T}{T_\infty} + \frac{\partial T}{\partial y} D_B \frac{\partial C}{\partial y} \right) . \tag{3}$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} + = D_B \frac{\partial^2 C}{\partial y^2} + \left( \frac{\partial^2 T}{\partial y^2} \right) \frac{D_T}{T_\infty} . \tag{4}$$

$$u \frac{\partial N}{\partial x} + v \frac{\partial N}{\partial y} = D_m \frac{\partial^2 N}{\partial y^2} + \frac{bw_c}{(C_\infty - C_w)} \frac{\partial}{\partial y} \left( N \frac{\partial C}{\partial y} \right) . \tag{5}$$

All of the variables used in these equations are specified using nomenclature. The next are the appropriate boundary conditions for the present literature [22]:

$$\text{When } y \rightarrow 0, N_w = N, C_w = C, v = 0, u = u_w = ax, \frac{\partial T}{\partial y} = -(T_w - T) \frac{h}{k_{hmf}} . \tag{6}$$

$$\text{When } y \rightarrow \infty, u \rightarrow 0, T \rightarrow T_\infty, C \rightarrow C_\infty, N \rightarrow N_\infty . \tag{7}$$

The flowing suitable similarity transform is used to transform dimensional equations into dimensionless flow governing equations [22]:

$$\begin{aligned} \eta &= \sqrt{\frac{a}{v_f}} y, \Psi(\eta) = f(\eta) \sqrt{av_f} x, \\ u &= \frac{\partial \Psi}{\partial y}, v = -\frac{\partial \Psi}{\partial x}, \theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty}, \\ \varphi(\eta) &= \frac{C - C_\infty}{C_w - C_\infty}, X(\eta) = \frac{N - N_\infty}{N_w - N_\infty}. \end{aligned} \tag{8}$$

In accordance with the above-mentioned similarity relations, Eq. (1) is demonstrated as being identical, and Eqs. (2) to (5) are reconstructed as:

$$f''' (\alpha_1 + \alpha_2 f'^2) + D_1 (ff'' - f'^2) + D_2 (\theta - N_r \varphi - N_c X) \lambda = 0. \quad (9)$$

$$\frac{1}{Pr D_3 D_4} \theta'' + f \theta' + N_T \theta'^2 + N_b \theta' \varphi' = 0. \quad (10)$$

$$\varphi'' + Le f \varphi' + \frac{N_T}{N_b} \theta'' = 0. \quad (11)$$

$$X'' - (X' \varphi' + (\omega + X) \varphi'') Pe + f L b X' = 0. \quad (12)$$

Then the Eqs. (6) and (7) are reconstructed as:

$$\text{When } \eta \rightarrow 0, f' = 1, f = 0, X = 1, \theta' = \frac{k_f}{k_{mf}} Bi (\theta - 1), \varphi = 1. \quad (13)$$

$$\text{When } \eta \rightarrow \infty, \varphi \rightarrow 0, \theta \rightarrow 0, f' \rightarrow 0, X \rightarrow 0. \quad (14)$$

where  $\theta, \varphi, X$  and  $f$  are functions of  $\eta$ . Furthermore, prime denotes the differentiation w.r.t  $\eta$ ,  $Pr = \frac{\nu_f (\rho C_p)_f}{k_f}$  is Prandtl number,  $Lb = \frac{\nu_f}{D_m}$  is Bioconvection Lewis number,  $Le = \frac{\nu_f}{D_B}$  indicates

Lewis number,  $\lambda = \frac{(1 - C_\infty) \beta g (T_w - T_\infty)}{ax^2}$  and  $Nt = \frac{\tau D_T (T_w - T_\infty)}{\nu_f T_\infty}$  denotes mixed convection and thermophoresis particle diffusion, respectively,  $Bi = \frac{h \sqrt{\nu_f}}{k_f \sqrt{\alpha}}$  is thermal Biot number, Peclet

number is  $Pe = \frac{b w_c}{D_m}$ , the bioconvection constant is denoted as  $\omega = \frac{N_\infty}{N_w - N_\infty}$ , the Rayleigh number

bioconvection is  $N_c = \frac{\gamma' (\rho_m - \rho_f) (N_\infty - N_w)}{(C_\infty - 1) \rho_f \beta (T_w - T_\infty)}$ ,  $Nb = \frac{\tau D_B (C_w - C_\infty)}{\nu_f}$  indicate the Brownian motion

parameter,  $Nr = \frac{(\rho_p - \rho_f) (C_w - C_\infty)}{(1 - C_\infty) \rho_f \beta (T_w - T_\infty)}$  denotes the buoyancy ratio parameter, the parameter of

Prandtl fluid is shown as  $\alpha_1 = \frac{A}{C}$ , the elastic parameter is defined as  $\alpha_2 = \frac{a^3 x^2 A}{2 \nu_f C^3}$ .

Now, the useful thermo-physical properties are the same as those explored by Hassan et al. [17,31]:

$$D_1 = (1 - \varphi_1) + \varphi_1 \frac{\rho_{s1}}{\rho_s} (1 - \varphi_2) + \varphi_2 \frac{\rho_{s2}}{\rho_s} (1 - \varphi_2)^{2.5} (1 - \varphi_1)^{2.5}, \quad (15)$$

$$D_2 = (1 - \varphi_2)^{2.5} (1 - \varphi_1)^{2.5}, \quad (16)$$

$$D_3 = \left( (1 - \varphi_1) + \varphi_1 \frac{(\rho C_P)_{s1}}{(\rho C_P)_f} \right) (1 - \varphi_2) + \frac{(\rho C_P)_{s1}}{(\rho C_P)_f} \varphi_1, \quad (17)$$

$$D_4 = \frac{\left( k_{b_f} - k_{s_2} \right) \varphi_2 + k_{s_2} + (m - 1) k_{b_f}}{(m - 1) k_f + k_{s_2} - (m - 1) (k_f - k_{s_2}) \varphi_2} \cdot \frac{(k_f - k_{s_1}) \varphi_1 + k_{s_1} + (m - 1) k_f}{(m - 1) k_f + k_{s_1} - (m - 1) (k_f - k_{s_1}) \varphi_1}. \quad (18)$$

Physical quantities are defined by Hussain et al. [22]:

$$C_{fx} = \frac{\tau_w}{\rho_f u_w^2}, \text{ And } Nu_x = \frac{xq_w^*}{k_f(T_w - T_\infty)}. \tag{19}$$

The heat flux  $q_w^* = -k_{hmf} \left( \frac{\partial T}{\partial y} \right)_{y=0}$ . The shear stress at the surface is  $\tau_w = \mu_{hmf} \left( \frac{A}{C} \frac{\partial u}{\partial y} + \frac{A}{2C^3} \left( \frac{\partial u}{\partial y} \right)^3 \right)$ .

Expressions (19) are transformed using the established similarity relations.

$$Cf_x = \frac{1}{B_2} R_{ex}^{-0.5} (\alpha_1 + \alpha_2 f''(0)^2) f''(0), \tag{20}$$

$$Nu_x = -\frac{k_{hmf}}{k_f} R_{ex}^{0.5} \theta'(0). \tag{21}$$

So, the relevant Reynolds number is  $R_{ex} = \frac{u_w x}{\nu_f}$ .

### 3 Numerical Solution

The non-dimensional Eqs. (9)–(12), which also have limits (13) and (14), have a nonlinear nature, making it challenging to compute them analytically. Due to the nonlinear nature of these differential equations, analytical solutions are not achievable. In the problem that is explained here, ordinary differential equations involving the boundary conditions are established to be capable of being resolved numerically. These solvable mathematical equations in the non-dimensional form have been numerically resolved as a result of adopting the `bvp4c` package, together with boundary conditions (13), (14). First-order PDEs are created by converting the associated nonlinear ordinary differential Eqs. (9)–(12) in Matlab using the `bvp4c` solver. The Matlab software is used to perform the above-mentioned mathematical programming in order to find a solution for the hybrid nanofluid over an extended sheet. The computing outcomes of the proposed method are compared with those acquired by other researchers, Wang [35] and Srinivasulu et al. [32], under restrictive conditions. Table 3 demonstrate that the present outcomes are consistent.

$$\begin{aligned} f' &= s_2, f = s_1, f'' = s_3, f''' = ss_1, \\ \theta &= s_4, \theta' = s_5, \theta'' = ss_2, \\ \varphi &= s_6, \varphi' = s_7, \varphi'' = ss_3, \\ \chi &= s_8, \chi' = s_9, \chi'' = ss_4. \end{aligned} \tag{22}$$

$$ss_1 = \frac{1}{(\alpha_1 + \alpha_2 s_3)} (-D_1 (s_1 s_3 - s_2^2) - D_2 \lambda (s_4 - Nr s_5 - Nc s_8)), \tag{23}$$

$$ss_2 = -Pr D_3 D_4 [(s_1 s_5) + N_b s_5 s_7 + N_T s_5^2], \tag{24}$$

$$ss_3 = -Les_1 s_7 - \frac{N_T}{N_b} ss_5, \tag{25}$$

$$ss_4 = (s_8 s_9 + (w + s_8) ss_3) Pe - Lbs_1 s_9. \tag{26}$$

The boundary conditions are also:

$$\text{When } \eta \rightarrow 0, s_2 = 1, s_1 = 0, s_8 = 1, s_5 = \frac{k_f}{k_{hmf}} Bi (\theta - 1), s_6 = 1. \quad (27)$$

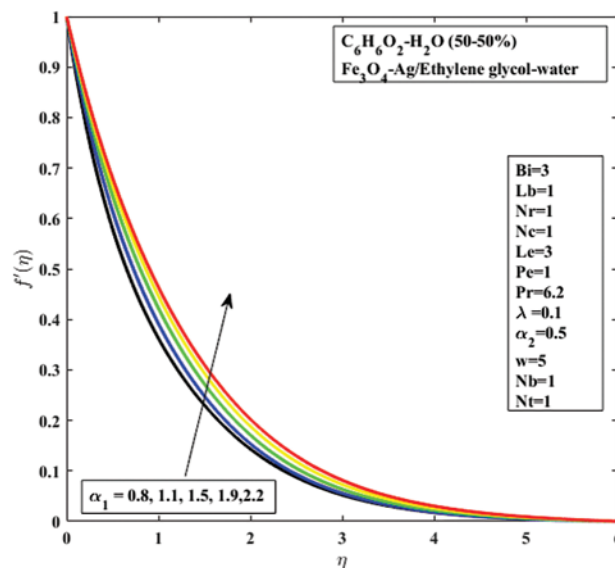
$$\text{When } \eta \rightarrow \infty, s_6 \rightarrow 0, s_4 \rightarrow 0, s_2 \rightarrow 0, s_8 \rightarrow 0. \quad (28)$$

### 3.1 Code Validation

In this part, the results of the current study are verified using previously released findings by Srinivasulu et al. [32], using modified parametric values. By altering the Prandtl number between  $0.7 < Pr \leq 70.2$  and while maintaining the elasticity ratio fixed at unity and setting all other parametric values to zero, we were able to compare the Nusselt number. The comparative findings are shown in Tables 3 and 4.

## 4 Results and Discussion

In this section the influence of different study parameters has been presented in a graphical manner on the distinct study profiles such as velocity, temperature, concentration, and microorganism profiles have been discussed. The effective range of the parametric values are as follows: elastic parameter ( $1 \leq \alpha_2 \leq 9$ ), Prandtl fluid parameter ( $0.8 \leq \alpha_1 \leq 2.2$ ), mixed convection ( $0.1 \leq \lambda \leq 3$ ), Prandtl number ( $1 \leq Pr \leq 1.8$ ), Lewis number ( $3 \leq Le \leq 3.8$ ), Bioconvection Lewis number ( $0.5 \leq Lb \leq 0.9$ ), and ( $1 \leq Pe \leq 1.16$ ). The influences of the elastic parameter ( $1 \leq \alpha_2 \leq 9$ ) and the Prandtl fluid parameter ( $0.8 \leq \alpha_1 \leq 2.2$ ) on the velocity profile are demonstrated in Figs. 2 and 3, respectively. With the increment in the elasticity and Prandtl fluid values, the velocity profile also enhances. Physically, when Prandtl fluid levels are increased the fluid viscosity decreases eventually raising the motion profile.



**Figure 2:** Prandtl fluid  $\alpha_1$  effect on velocity  $f'$

The overall outcome of mixed convection ( $0.1 \leq \lambda \leq 3$ ) on the distribution of velocity in the hybrid nanofluid is demonstrated in Fig. 4. This tendency is brought on by the beneficial mixed convection, which acts as a helpful pressure gradient and accelerates the fluid in the boundary layer.



The velocity profile for hybrid nanofluid flow grows as these mixed convection levels rise. Figs. 5 and 6 show the outcomes of the Rayleigh bioconvection number and buoyancy ratio, respectively, on the motion profile of hybrid nanofluid. The hybrid nanofluid flow velocity profile is grown by increasing the standard values of the  $N_r$  parameter.  $N_r$  has a sustaining force-like behaviour that accelerates the fluid nanoparticles close to the sheet. Similarly, the velocity profile for hybrid nanofluid flow rises as  $N_c$  parameter values increase.

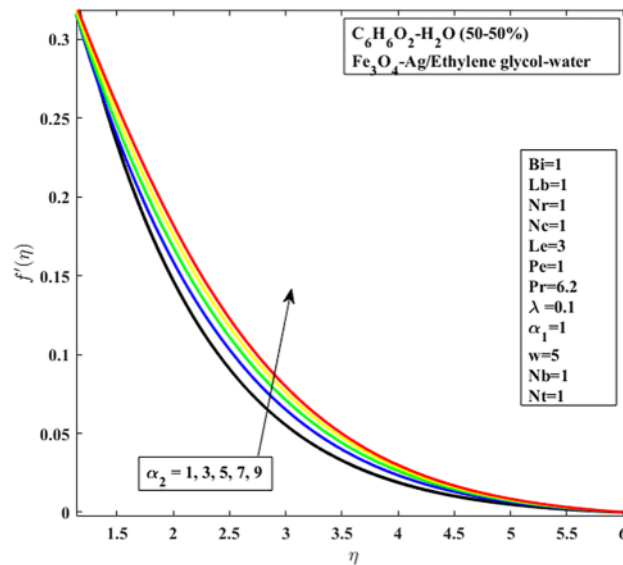


Figure 3: Elasticity  $\alpha_2$  impact on velocity profile  $f'$

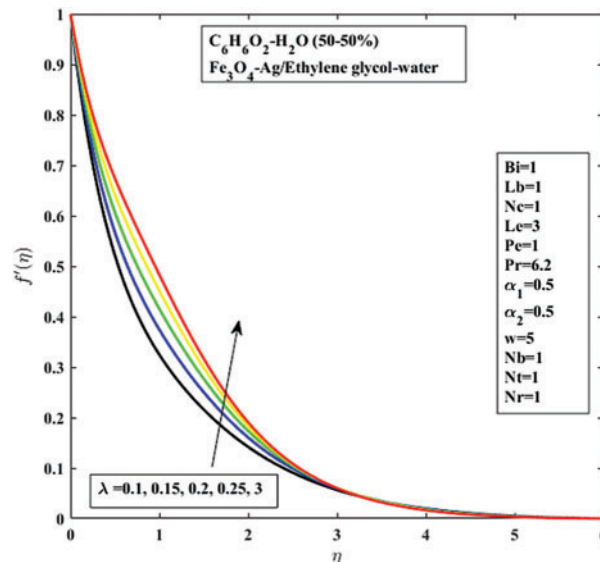
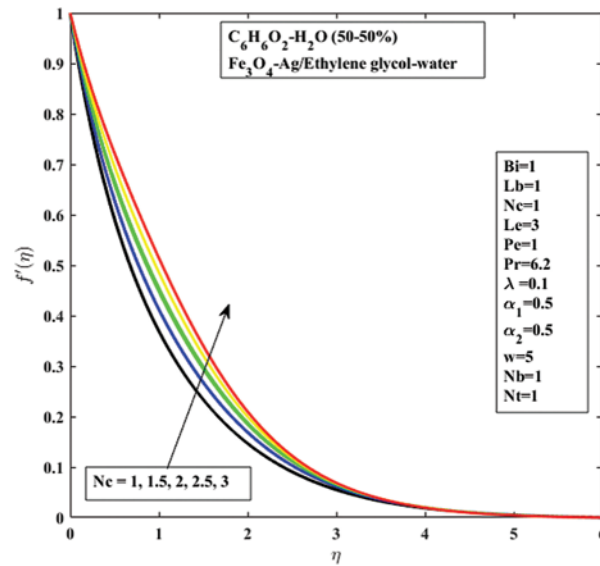
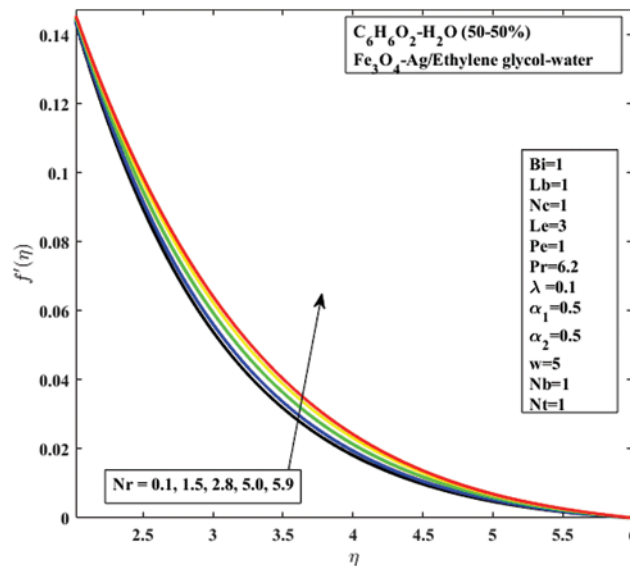


Figure 4: Mixed convection  $\lambda$  effect on  $f'$

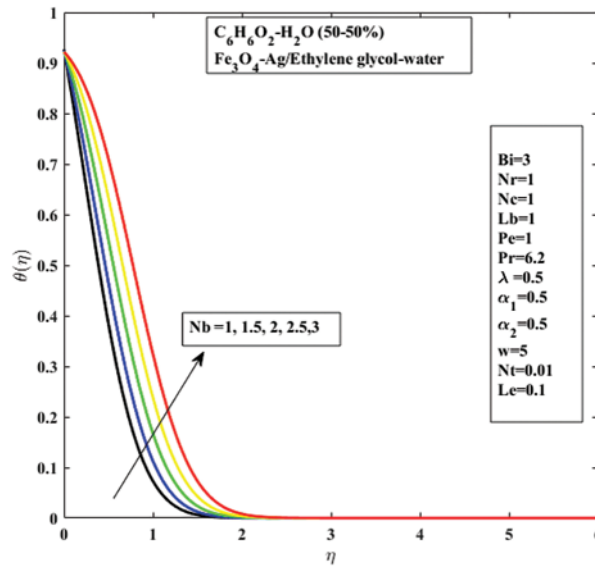


**Figure 5:** Rayleigh bio-number  $Nc$  vs.  $f'$

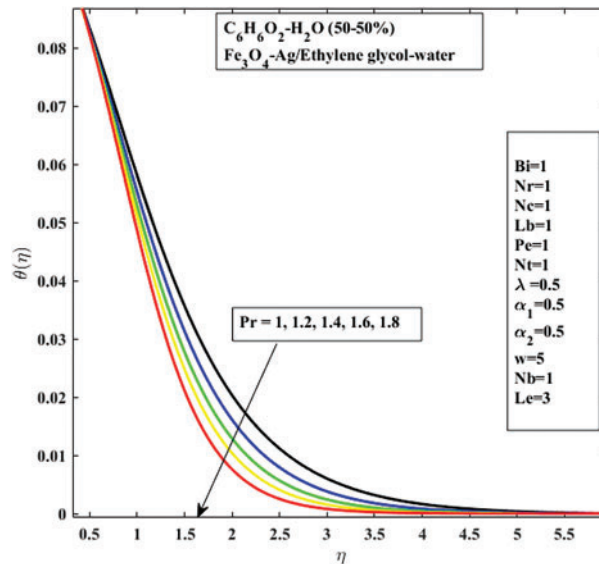


**Figure 6:** Buoyancy ratio  $Nr$  impact on  $f'$

The effects of  $N_b$  and  $Pr$  on temperature are shown in Figs. 7, 8. It has been demonstrated that as the Prandtl number ( $1 \leq Pr \leq 1.8$ ) values increase, the temperature curves of hybrid nanofluids go down, as shown in Fig. 7. The opposite behaviour is shown by the  $N_b$  parameter in Fig. 8. Brownian motion parameter temperature curves grow as parameter values rise. The effects of  $Bi$  and  $N_r$  on temperature are shown in Figs. 9, 10, respectively. Figs. 9, 10 show that the temperature profile of hybrid nanofluids is enhanced when the thermal Biot number and levels of thermophoresis particle diffusion are increased. By extending the energy gradient, moving forward towards the surface results in a drop in the temperature boundary layer thickness, as seen by the thermal Biot number.



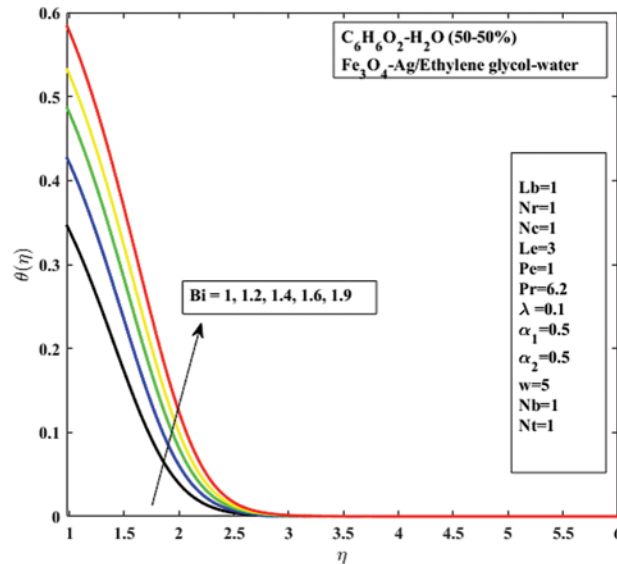
**Figure 7:** Brownian motion  $Nb$  effect on  $\theta$



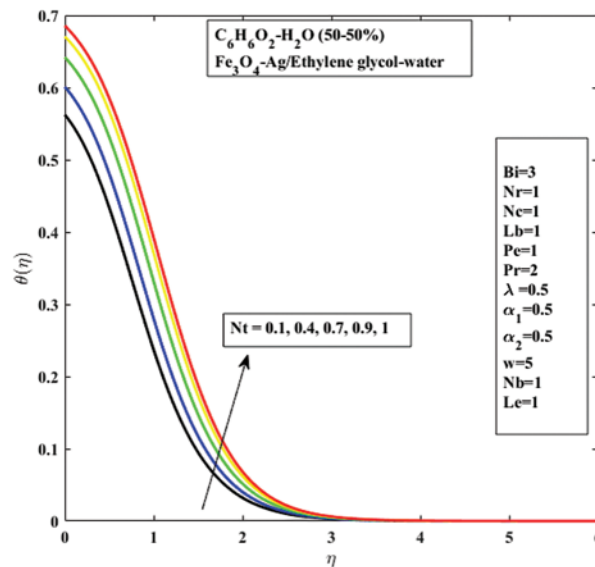
**Figure 8:** Prandtl number  $Pr$  impact on  $\theta$

Fig. 11 demonstrates the effect of the Lewis number ( $3 \leq Le \leq 3.8$ ) on hybrid nanofluid profile over the heated surface. Fig. 11 demonstrates that fluid concentration dropped as  $Le$  values were augmented. The boost in kinetic energy brought through the fluid particles and accelerated velocity causes the boundary layer to enlarge. Theoretically, the particle migrating from peak locations to bottom areas accelerates quickly as ( $Le$ ) parameter values grow. Figs. 12, 13 demonstrate the pattern of the hybrid nanofluid concentration curve for rising values of  $N_b$  and  $N_t$  on the concentration curves. Thermophoresis and Brownian motion have opposite effects on the concentration profile. The increase in Brownian motion declines the concentration profile abruptly, while the increment in thermophoresis

motion enhances the concentration profile. Further, it has been observed that with an increase in thermophoresis particle diffusion, the concentration layer thickness has increased.



**Figure 9:** Thermal Biot number  $Bi$  vs.  $\theta$



**Figure 10:** Thermophoretic  $Nt$  effect on  $\theta$

Figs. 14, 15 demonstrate the hybrid nanofluid behaviour of motile microbes ( $X(\eta)$ ) under the impact of bio-convective and Peclet numbers, respectively. Fig. 14 demonstrates the effect of the Bio-convection Lewis number ( $0.5 \leq Lb \leq 0.9$ ). Based on outcomes, when  $Lb$  levels rise, microbe profiles also decline. Additionally, it is worth mentioning here that with an increase in the bio-convection number, the thickness of the microbe’s boundary layer has expanded rapidly. Microorganism trends are shown in Fig. 15 under the influence of the Peclet number ( $1 \leq Pe \leq 1.16$ ). It has been observed

that with an increase in Peclet number, the microbes profile significantly enhances. In Figs. 16–18, skin friction and Nusselt number have been demonstrated in a graphical manner under the variation of buoyancy and Brownian motion effect.

Table 1 gives the values of the thermophysical properties of the nanoparticles and working base fluid. Table 2 provides thermophysical relations used in this study for different fluid hybrid nanofluid properties. Table 3 demonstrates the comparison between current results and already published results. Table 4 shows the numerical outcome of the skin friction coefficient and Nusselt number under the influence of different study parameters such as Prandtl fluid parameter  $\alpha_1$ , Elastic parameter  $\alpha_2$ , Brownian motion  $N_b$ , thermophoresis diffusion  $N_t$ , mixed convection parameter  $\lambda$ , thermal Biot number  $Bi$ , Prandtl number  $Pr$  and buoyancy ratio  $N_r$ .

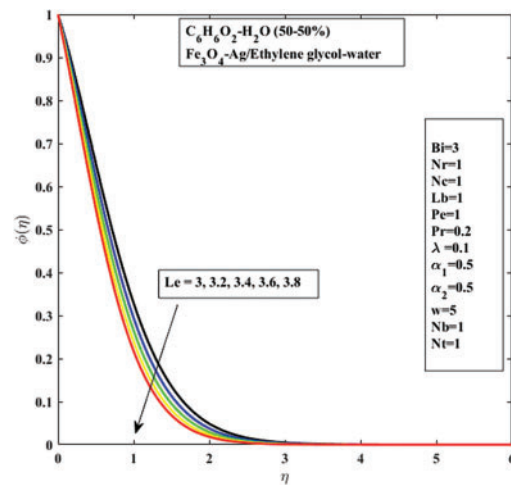


Figure 11: Lewis number  $Le$  impact on  $\varphi$

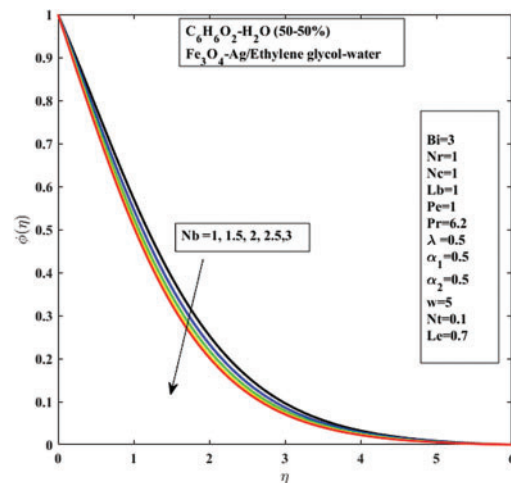
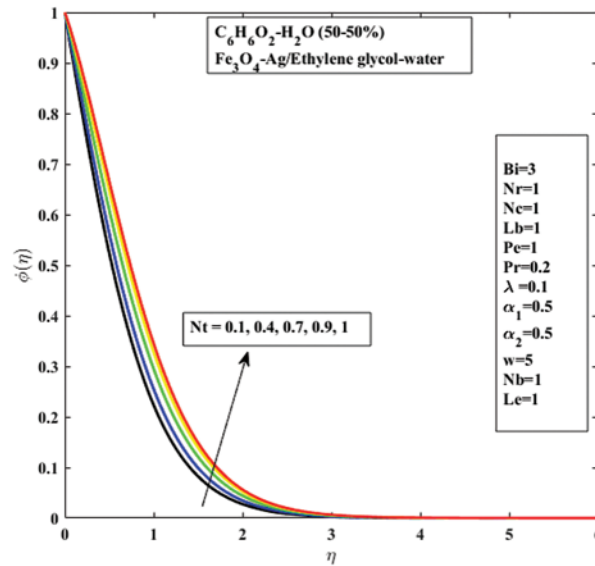
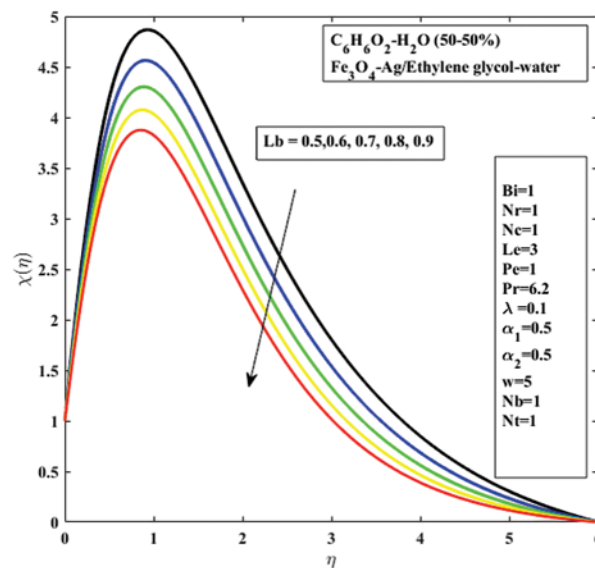


Figure 12: Brownian motion  $Nb$  effect on  $\varphi$



**Figure 13:** Thermophoretic  $Nt$  effect over  $\varphi$



**Figure 14:** Bioconvective Lewis number  $Lb$  effect vs.  $X$

It has been observed that skin friction and Nusselt number both decrease when the Prandtl fluid parameter and elastic parameter values are in the range of  $0.4 < \alpha_1 \leq 0.7$  and  $0.5 < \alpha_2 \leq 0.7$ , respectively. When the Brownian motion values are in the range of  $0.1 < N_b \leq 0.4$ , the Nusselt number is enhanced whereas the skin friction also increases Thermophoretic particle diffusion  $0.1 < N_t \leq 0.4$  significantly decrease the Nusselt number but increase the surface drag. It can be verified from [Table 3](#) that mixed convection  $0.1 < \lambda \leq 0.4$  parameters decrease both Nusselt number and skin friction.

Raising the values of thermal Biot number  $0.5 < Bi \leq 2.0$  decrease the Nusselt number and skin friction. A significant decline has been observed in the Nusselt number by raising the values of Prandtl

number  $1 < Pr \leq 4$ . Whereas buoyancy ratio  $0.5 < Nr \leq 0.8$  decrease the surface drag but it improves the Nusselt number.

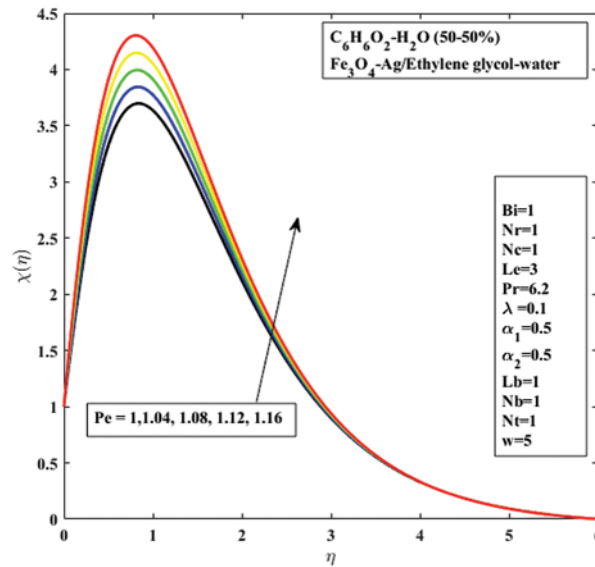


Figure 15: Peclet number  $Pe$  effect on  $X$

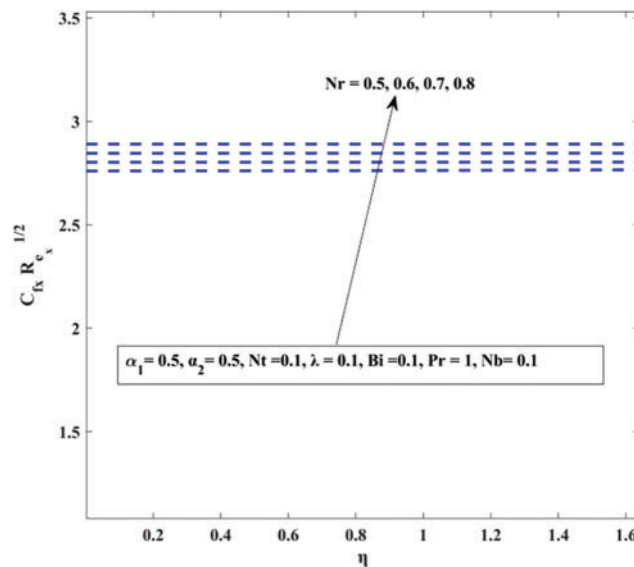
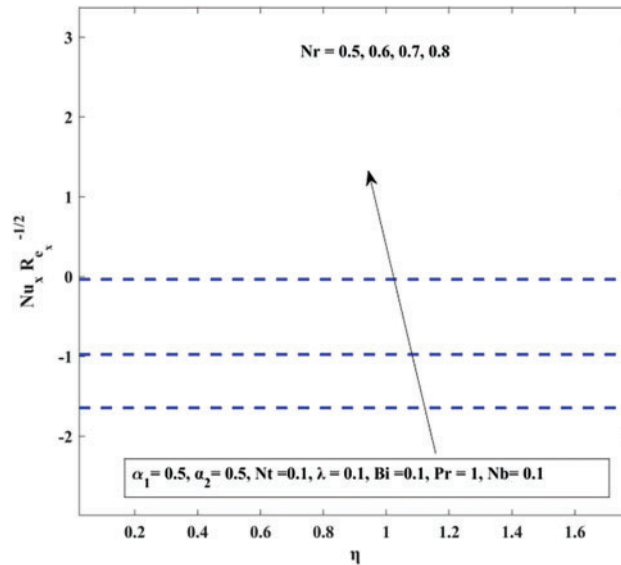


Figure 16: Skin friction under Buoyancy ratio effect

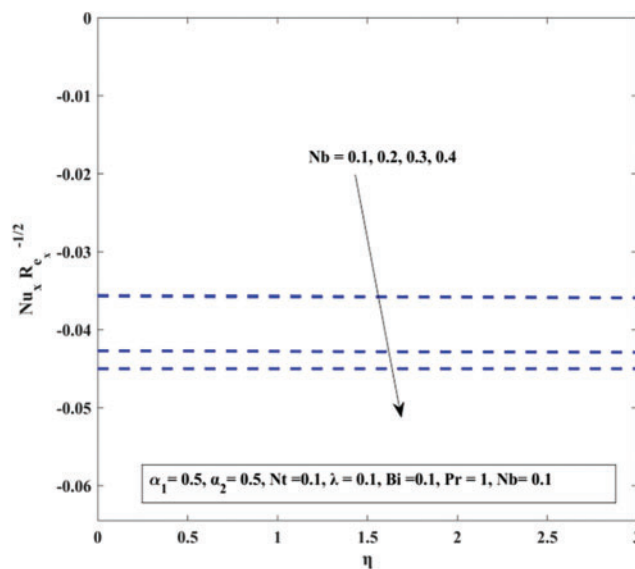
#### 4.1 Models of Thermal Conductivity and Viscosity

It can be easily observed through the above conducted literature review that theoretical viscosity and thermal conductivity models are commonly used by researchers (Srinivasulu et al. [32], Hassan et al. [31] and Saleem et al. [11]). These models have their own pros and cons. In this study, we have utilized thermophysical relations proposed in Yamada Ota (Yamada et al. [33]) model

hybrid nanofluid. This model has been used by numerous researchers such as Hussain et al. [22], Saleem et al. [11], and Srinivasulu et al. [32]. The advantage of these theoretical models is that they are easy to use, time-saving, and easily available, whereas it is difficult to find experimental data and experiments are costly. It is generally more easy to use an already established theory rather than looking for the experimental outcomes. Additionally, with theoretical thermophysical relation optimal outcomes can never be possibly achieved as demonstrated by Xue [34]. Therefore, researchers prefer to use already established results or thermophysical relations.



**Figure 17:** Nusselt number with buoyancy ratio



**Figure 18:** Nusselt number with Brownian motion



**Table 1:** Numeric values of nanoparticles and base fluid properties [31]

Properties	Fe <sub>3</sub> O <sub>4</sub>	Ag	Ethylene glycol-water (50%-50%)
$\rho$ ( $kgm^{-3}$ )	5180	10500	1063.8
$C_p$ ( $Jkg^{-1}K^{-1}$ )	670	235	3630
$k$ ( $Wm^{-1}K^{-1}$ )	9.7	429	0.387

**Table 2:** Thermo-physical properties relations of a hybrid nanofluid [31]

Properties	Hybrid nanofluid
Thermal conductivity	$\frac{k_{hnf}}{k_{bf}} = \frac{(k_{bf-k_{s2}}) \varphi_2 + k_{s2} + (m-1) k_{bf}}{(m-1) k_f + k_{s2} - (m-1) (k_f - k_{s2}) \varphi_2},$ $\frac{k_{bf}}{k_f} = \frac{(k_f - k_{s1}) \varphi_1 + k_{s1} + (m-1) k_f}{(m-1) k_f + k_{s1} - (m-1) (k_f - k_{s1}) \varphi_1}.$
Density	$\rho_{hnf} = (1 - \varphi_2) + \varphi_1 \frac{\rho_{s1}}{\rho_f} (1 - \varphi_1) \rho_f + \varphi_2 \rho_{s2}.$
Dynamic viscosity	$\mu_{hnf} = \frac{\mu_f}{(1 - \varphi_2)^{2.5} (1 - \varphi_1)^{2.5}}.$
Heat capacity	$(\rho C_p)_{hnf} = (\rho C_p)_f (1 - \varphi_2) \left( (1 - \varphi_1) + \varphi_1 \frac{(\rho C_p)_{s1}}{(\rho C_p)_f} \right).$

**Table 3:** Comparison of present outcomes for the heat transfer coefficient (Nusselt number) when  $\alpha_1 = 1$ , and all other parameters are 0

$Pr$	Srinivasulu et al. [32]	Wang [35]	Our outcomes
0.2	0.1723	0.1691	0.172400
0.7	0.4539	0.4539	0.454000
2.0	0.9113	0.9114	0.911370
7.0	1.8954	1.8954	1.895301
70.0	6.4621	6.4622	6.462099

**Table 4:** Heat transfer and drag coefficient outcomes under the impact of varying parametric values of study parameters [22]

$\alpha_1$	$\alpha_2$	$N_b$	$N_t$	$\lambda$	$Bi$	Pr	$N_r$	$Nu_x R_{ex}^{-0.5}$	$C_{fx} R_{ex}^{0.5}$
0.4	0.5	0.1	0.1	0.1	0.5	1.0	0.5	-0.0340	3.3231
0.5								-0.0356	2.7612
0.6								-0.0694	2.4443

(Continued)

**Table 4 (continued)**

$\alpha_1$	$\alpha_2$	$N_b$	$N_t$	$\lambda$	$Bi$	$Pr$	$N_r$	$Nu_x R_{ex}^{-0.5}$	$C_{fx} R_{ex}^{0.5}$
0.7								-0.0992	2.2553
0.5	0.4							-0.0398	2.2463
	0.5							-0.0356	2.7603
	0.7							-0.0315	3.3413
	0.6							-0.0274	3.9945
	0.5	0.1						-0.0356	2.7599
		0.2						-0.0357	3.2535
		0.3						-0.0427	3.4113
		0.4						-0.0450	3.4922
		0.1	0.1					-0.0356	2.7599
			0.2					-0.9767	3.5177
			0.3					-0.9768	4.0750
			0.4					-1.6454	7.5620
			0.1	0.1				-0.0356	2.7599
				0.2				-0.2056	1.8772
				0.3				-0.3157	1.3834
				0.4				-0.4000	1.0706
				0.1	0.5			-0.3558	3.4922
					1.0			-0.3561	3.0871
					1.5			-0.7518	2.8822
					2.0			-0.7519	2.7559
					0.5	0.1		-0.0356	2.7599
						0.2		-0.0556	2.7556
						0.3		-0.0712	2.7513
						0.4		-0.0841	2.7470
						0.1	0.5	-0.0256	2.7599
							0.6	-0.0325	2.8031
							0.7	-0.0295	2.8469
							0.8	-0.0265	2.8911

## 5 Conclusion

In this work, the bio-convection enhanced heat transfer using the hybrid nanofluids with modified Buongiorno's model over a stretched surface has been investigated. The Prandtl fluid model has been utilized and convective boundary condition is also taken into consideration. The flow governing the nonlinear partial differential model is converted into ordinary dimensionless differential equations. The dimensionless flow equations are solved numerically using the bvp4c package in MATLAB. The following are the key findings of the present study:

- The velocity curve of the  $Fe_3O_4$ -Ag/EG-water-base hybrid nanofluid increases by raising the levels of the buoyancy ratio and mixed convection.
- Elastic, thermophoresis parameter and thermal Biot number values enhance temperature profile. The Brownian motion and the Prandtl number, however, exhibit the reverse trend.

- On the concentration profile, the hybrid nanofluid exhibits a decrease in the Lewis number but increases in the elasticity, thermophoresis, and buoyancy ratio.
- Peclet parameters, elastic parameters, and bio-convection numbers all demonstrate an enhancement in the microorganisms' profile in hybrid nanofluids. The Bioconvection Lewis number, however, displays the reverse trend.
- The skin friction coefficient of hybrid nanofluid demonstrates a tendency to decrease for greater inputs of  $Pr$ ,  $\alpha I$ , and  $Bi$ , but an increasing tendency is shown for  $Nb$ ,  $N_t$ ,  $\alpha_2$ , and  $N_r$ .
- The Nusselt number falls for  $Pr$ ,  $Nb$ , and  $Nt$ , whereas it increases for  $\alpha_2$ ,  $Bi$ , and  $N_r$ .

**Acknowledgement:** The authors would like to acknowledge the constructive remarks by worthy reviewers that led to this revised article.

**Funding Statement:** The authors received no specific funding for this study.

**Author Contributions:** The authors confirm their contribution to the paper as follows: study conception and design: Azad Hussain and Ali Hassan; data collection: Saira Raiz; analysis and interpretation of results: Saira Raiz and Ali Hassan; original draft manuscript preparation: Saira Raiz; writing review and editing: Ali Hassan; investigation and validation: Mohamed R. Ali and Abdulkafi Mohammed Saed. All authors reviewed the results and approved the final version of the manuscript.

**Availability of Data and Materials:** Data can be made available from the corresponding author following a special request.

**Conflicts of Interest:** The authors declare that they have no conflicts of interest to report regarding the present study.

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