



**ARTICLE**

# Numerical Simulation of the Mechanical Stirring Process in a Tannin-Based Foaming Precursor Resin

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## ABSTRACT

Tannin foam is a new functional material. It can be widely applied to the automobile industry, construction industry, and packaging industry due to its wide range of raw materials, renewable, easily degraded, low cost and almost no pollution. Preparing tannin foam is a very complex process that includes high temperature, two phases, mechanical agitation, and phase change. To investigate the influence of the stirring velocity and paddle shape, simulation was calculated by making use of the volume of fluid (VOF) method and multiple reference frame (MRF) method in a three-dimensional flow field of tannin-based foaming precursor resin. The gas holdup and velocity magnitude were analysed with various conditions of mechanical velocities and paddle shape in the stirring flow field. The result shows the higher the velocity, the greater the disturbance and paddle shape between the eggbeater and the Rushton turbine, obviously the paddle shape of the eggbeater with a wider range of agitation, which can entrap more air into the tannin-based foaming precursor resin in a short time. Especially when the speed is 1500 rpm, the flow field of the Rushton turbine comes out of a ditch, which decreases the efficiency of mass transfer; there is less air to mix into the tannin-based foaming precursor resin, which causes unevenness. At the same time, the eggbeater shows the marvelous capability of hybrid as it has two vortices and multiple cycles that make a difference from the Rushton turbine, which has only one vortex and two upper and lower loops; the structure makes the flow field more stable allowed evenness of flow field tannin-based foaming precursor resin. The results reveal that it is beneficial for tannin-based foaming precursor resin to use an eggbeater with a speed of 1500 rpm to reduce the consumption of resources while obtaining a uniform flow field.

## KEYWORDS

Tannin-based foaming precursor resin; mechanical agitation; numerical simulation; volume of fluid (VOF); multiple reference frame (MRF)

## 1 Introduction

Foams with excellent properties, such as light weight, thermal insulation, and flame retardancy, have been prepared from tannin material [1–5]. At present, a series of efforts have focused on its preparation and modification, with the aim of designing and fabricating tannin-based foams with outstanding overall



performance. Among those preparation methods, mechanical stirring foaming is considered a promising green way to perform this task, which allows avoidance of foaming agent application; obviously, this approach imposes a lesser burden on the environment.

The process of mechanical stirring is vital to foam structure and performance, but most of the literature on the foaming process is based on metal foam, and there is a lack of research on the foaming process of tannin-based foam. There has been much research investigating the properties of formed tannin-based foams in recent years, but few studies have analysed the foaming process, which presents a large challenge for obtaining relative experimental parameters [6–9]. To investigate the effect of the mechanical stirring process on tannin-based foam, the present work was carried out. During the foaming process, factors such as the mechanical stirring velocity, the shape of the agitator, the physical properties of the solution, the temperature of the environment and the duration of stirring all have crucial impacts on the whole foaming process. In addition, studies of stirred vessel flows are complex, and many factors arise in the process of mechanical stirring [10–13].

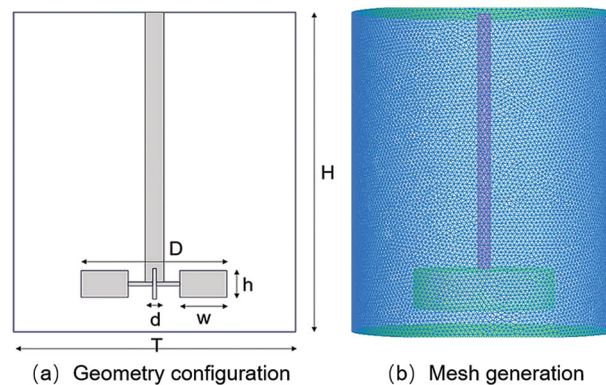
In the past few decades, most of the research on stirred flow has adopted numerical simulation methods. Many studies have focused on the stirring flow field, and most of these studies focused on the geometry, Reynolds number, power number, bubble distribution and rotation speed [14–17]. It seems that the flow field is affected by many factors, so it is vital to investigate the effect on the foaming process to improve the flow field. Mechanical stirring involves not only mixing the contents but also mixing the air, which proves that the foaming process is a part of fluid mechanics. This work explores the foaming process of tannin foam [13], which is prepared without a foaming agent based on mechanical agitation in physical foaming, and through the experience of foam preparation in the laboratory, the mechanical stirring velocity and shape of the agitator have a great influence on the foaming process. Under the working conditions of two-phase flow, the improved mixing and speed increase cause the vortex to occur at the same time and space, leading to the production of more small-scale structures, compression of the other spatial structures of the vortex structure, and an increase in the energy cascade effect; thus, the comprehensive performance under high speed and scope of work is better than that under lower speed, but at the same time, the gas phase can more easily accumulate at a lower rotational speed, and literature studies have also proven that the mechanical rotational speed and impeller type have significant effects on the structure of the flow field [18,19]. On this basis, both the velocity and paddle shape are selected as indexes to investigate the effect on the flow field, which is beneficial for improving the foaming process via mechanical stirring. There are some relationships among the rotation speed, the paddle shape and the stirring effect, and it has been concluded that as the rotation speed increases, the energy cost decreases [20]. Many paddle shapes, including the Rushton turbine, perturbed six-bent-bladed turbine, anchor-type, squire type, disc impeller, etc., were simulated to investigate the effect on the flow field, and the Rushton turbine has received much attention [21–24].

With the development of computational fluid dynamics (CFD) techniques and multiple reference frame (MRF) methods, it has become possible to simulate the foaming process by mechanical stirring. In this study, numerical simulation was used to simulate the mechanical stirring-induced foaming process, and the influence of stirring speed and paddle shape on this process was considered. With respect to the simulation of stirring impellers, there are two classical methods, multiple reference frames and moving meshes; the former fit complex working conditions, but the latter fit the unsteady process of changing the flow field. The MRF method is a mature technology that has been extensively utilized in rotation areas [25]. The interface between the rotation domain and container made use of changing the velocity and mass to mix the entire region. The selection of the interface is usually accomplished in the range of 1/8–1/4 (the value is the interface with the distance of the container). This approach is suitable for the simulation environment for the experimental range. Through numerical simulation of the foaming

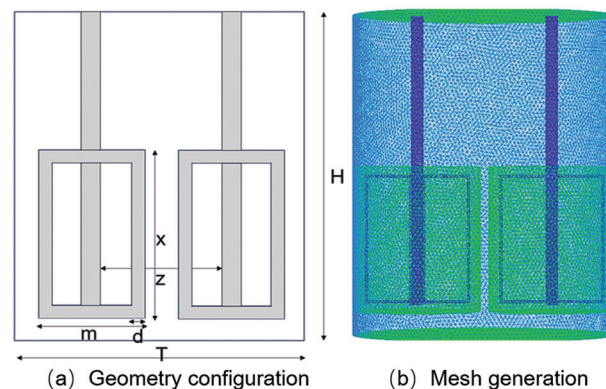
process, the motion characteristics, distribution characteristics and various parameters of bubbles can be more accurately determined to obtain better foaming materials. Therefore, the use of numerical simulations to study the tannin foaming process will be the focus of future research.

## 2 Geometric Description

This study considered the transformation of the flow field caused by changes in the velocity and shape of the paddle in the rotor area of the stirrer beaker, which is also called the rotation domain. To simplify the calculation process, the stirrer beaker is assumed to be a cylinder, which has a standard Rushton impeller of six blades, as shown in Fig. 1a. The diameter of the cylindrical beaker is  $T = 120$  mm, and the height is  $H = 150$  mm. The blade diameter is  $D = 45.3$  mm, the blade width and height are  $w = T/4$  and  $h = T/5$ , respectively, the blade thickness is  $d = 1$  mm, and the impellers are 15 mm from the bottom. And the mesh structures of the working conditions used in this work is shown in Fig. 1b. To investigate the influence of the paddle shape on the flow field, the eggbeater model was established and calculated via the same procedure as for the Rushton turbine, and the size of the eggbeater was the same scale used in the laboratory. The geometry of the eggbeater is illustrated in Fig. 2a. It has two rotation centers, each of which has the same structure as the impellers, and the data were measured in the laboratory by vernier callipers. The impeller width was  $m = 47.4$  mm, the length was  $z = 36$  mm, and the height was  $x = 60$  mm. And the mesh structures of the working conditions used in this work is shown in Fig. 2b.



**Figure 1:** The geometry and mesh of the Rushton turbine



**Figure 2:** The geometry and mesh of the egg beater

### 3 The Flow Pattern in the Beaker

The tannin-based foaming precursor resin was used to simulate the experimental environment, and the pattern of the flow field can be calculated by the stirring Reynolds number according to the equation below:

$$\text{Re} = \rho N D^2 / 60 \mu \quad (1)$$

In the equation, Re is a dimensionless number that can be used to characterize the flow of a fluid,  $\rho$  is the density of the tannin-based foaming precursor resin,  $\rho = 1421.667 \text{ kg/m}^3$ ,  $N$  is the rotational speed, and there are three speeds in the simulation, as demonstrated in the Table 1.  $D$  is the diameter of the impeller,  $D = 45.3 \text{ mm}$ ,  $\mu$  is the dynamic viscosity of the tannin-based foaming precursor resin, and  $\mu = 0.4283 \text{ Pa}\cdot\text{s}$ . In this work, there were three kinds of rotational speeds, and the relative Reynolds number was calculated as given in Table 1, where the results show that the flow pattern is laminar.

**Table 1:** Reynolds number at various rotation speeds

N (rpm)	Re
1000	113.526
1250	141.907
1500	170.289

### 4 Mathematical Model

There are several assumptions that were used to establish the mathematical model. (1) Without regard to the change in energy, (2) the effect of the fluctuation of the tannin resin surface on the initial parameters was ignored. Based on the assumptions above, the laminar flow model is a mathematical model used to simulate the motion of a bubble.

#### 4.1 Governing Dimensional Equations

In the stirring beaker, both the bubble and the tannin resin govern the conservation of mass, momentum, and energy; however, the effects on the flow field of the mechanical velocity and paddle shape were mainly discussed, so the energy conversion was not considered. The governing equations are as follows:

Continuity equations:

$$\frac{\partial(\rho_a \alpha_a)}{\partial t} + \nabla \cdot (\rho_a \alpha_a \mathbf{u}_a) = 0 \quad (2)$$

Momentum equations:

$$\rho \left( \frac{\partial(\rho_a \alpha_a \mathbf{u}_a)}{\partial t} + \nabla \cdot (\rho_a \alpha_a \mathbf{u}_a \mathbf{u}_a) \right) = -\alpha_a \nabla p + \nabla \cdot (\alpha_a \mu_{eff} (\nabla \mathbf{u}_a + (\nabla \mathbf{u}_a)^\top)) + \alpha_a \rho_a \mathbf{g} + \mathbf{F}_a \quad (3)$$

In the equation,  $\mathbf{F}$  is the interfacial force of a two-phase fluid,  $\rho$  is the density of the fluid,  $\mathbf{u}$  is the average velocity,  $p$  is the pressure,  $\mu_{eff}$  is the effective viscosity,  $\mathbf{g}$  is the gravitational acceleration, and  $a$  represents the tannin-based foaming pressure and gas phase.



#### 4.2 VOF Model

The VOF model simulates two or more kinds of immiscible fluids by solving a single momentum equation and tracking the volume fraction of the single-phase fluid in the region, which helps accurately track the movement changes of the gas–liquid phase interface. In this simulation, the phase transition process is not considered, and its governing equation is as follows:

Continuity equations:

$$\frac{\partial(a_q \rho_q)}{\partial t} + \nabla(a_q \rho_q v_q) = 0 \quad (4)$$

In the equation,  $a_q$  is the volume fraction of the  $q$  phase fluid,  $\rho_q$  is the density of the  $q$  phase,  $t$  is time, and  $v_q$  is the velocity of phase  $q$ . The volume fraction of each phase is satisfied in the VOF:

$$\sum_{q=1}^n a_q = 1 \quad (5)$$

#### 4.3 MRF Model

In this study, the multiple reference frame method (MRF) was used to simulate the four-blade impeller. The steady-state approximation is used for each region with different rotation or moving speeds, and the multiple reference frame model (MRF) is more suitable for the boundary of the grid region, assuming that each cell moves at a uniform speed.

The calculation region is divided into several subregions by a multireference model, and the governing equations of each subregion are determined, in which each subregion moves differently. The transfer and exchange of the flow field structure information at the interface of each subdomain is realized by changing the velocity. In the numerical calculation, the rotor of the rotating reference frame is given a fixed angular velocity as the power source. Using the velocity and velocity gradient fluctuations, the transformation from the motion reference frame to the absolute reference frame can be verified. The transformation formula is as follows:

$$\mathbf{r} = \mathbf{x} - \mathbf{x}_0 \quad (6)$$

where  $\mathbf{x}_0$  is the initial position of the rotating shaft and  $\mathbf{x}$  is the guide coordinate of position  $\mathbf{x}_0$ .

The relative velocity in a moving reference frame can be converted to a value in a stationary reference frame via Eq. (7):

$$\mathbf{v} = \mathbf{v}_r + (\boldsymbol{\omega} \times \mathbf{r}) + \mathbf{v}_t \quad (7)$$

where  $\mathbf{v}_r$  is the relative velocity in the noninertial reference frame,  $\mathbf{v}$  is the absolute velocity in the inertial reference frame, and  $\mathbf{v}_t$  is the translation velocity of a noninertial reference frame.

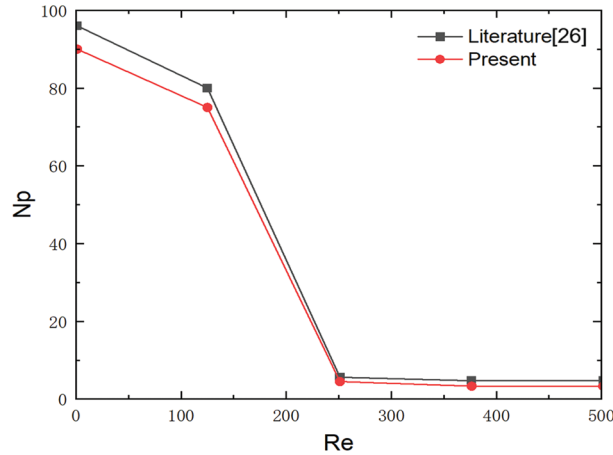
Based on the relative velocity, the vector gradient of the absolute velocity is obtained:

$$\nabla \mathbf{v} = \nabla \mathbf{v}_r + \nabla(\boldsymbol{\omega} \times \mathbf{r}) \quad (8)$$

#### 4.4 Validation of the Mathematical Model

To verify the simulation, the work was completed under the same simulation conditions reported in the literature [26], and the variation in the power number vs. the Reynolds number is displayed in Fig. 3. The qualitative comparison between the simulation results and experimental results in the same calculation

domain showed good agreement with the literature results, indicating that the results of the numerical simulation are reliable.



**Figure 3:** Impeller power vs. the Reynolds number

## 5 Numerical Simulation

The commercial software FLUENT 2022R1 is used to simulate the gas–liquid two-phase flow in tannin resin to determine the influence of the mechanical agitation speed and paddle shape by solving the governing equations of the VOF model and laminar model. The simulation solution method is based on finite volume discretization. Postprocessing was performed with a TECPLOT360. The solver type is pressure based, and time is defined as transient. The coupling of pressure and velocity is implemented by the SIMPLE algorithm. The momentum is second-order upwind, and the discrete format of the volume fraction is first-order implicit. Due to the complex structure of the mixing stirrer beaker, tetrahedral unstructured mesh with strong spatial adaptability was used to mesh the fluid in the stirrer beaker. A preprocessor (ANSYS Meshing) was used to discretize the flow domain with an unstructured tetrahedral mesh. A mesh test was performed to ensure the accuracy of our predicted results. The mesh had a tetrahedral shape, and there were 362,167 cells, 750,693 cells, and 1,159,786 cells. Grid independence was confirmed by comparing results that showed a difference between the last two reference values of less than 5% (Table 2). To reduce the amount of computation, the middle cells were selected for calculation. The MRF technique was used to calculate the flow field by dividing the calculation domain to realize the exchange and replacement of flow field information at the interface. The stirring shaft, impeller and blade wall are set as moving wall boundary conditions: the impeller, blade, and part of the rotating shaft wall in the moving region are static relative to the surrounding moving fluid. The axis wall in the static region is in motion with respect to the surrounding static fluid. The inner wall of the mixing tank is set as the static wall boundary condition. The time step was 0.001 to ensure the convergence of the solution process. The calculation was deemed to be complete when the residuals of all variables were less than 0.001.

They are reported in ANSYS: [report]-[forces], where the moment center inputs the center of gravity of the impeller, and in the wall zones, all the rotating wall surfaces are chosen; the results are printed, and the data are calculated via the following equations:

$$p_a = N_p \rho N^3 D^5 \quad (9)$$

$$p_a = 2\pi N \tau \quad (10)$$

$$N_p = 2\pi \tau / \rho N^2 D^5 \quad (11)$$

**Table 2:** Mesh independence test

Elements	362,167	750,693	1,159,786
Np	0.0488	0.0516	0.0508

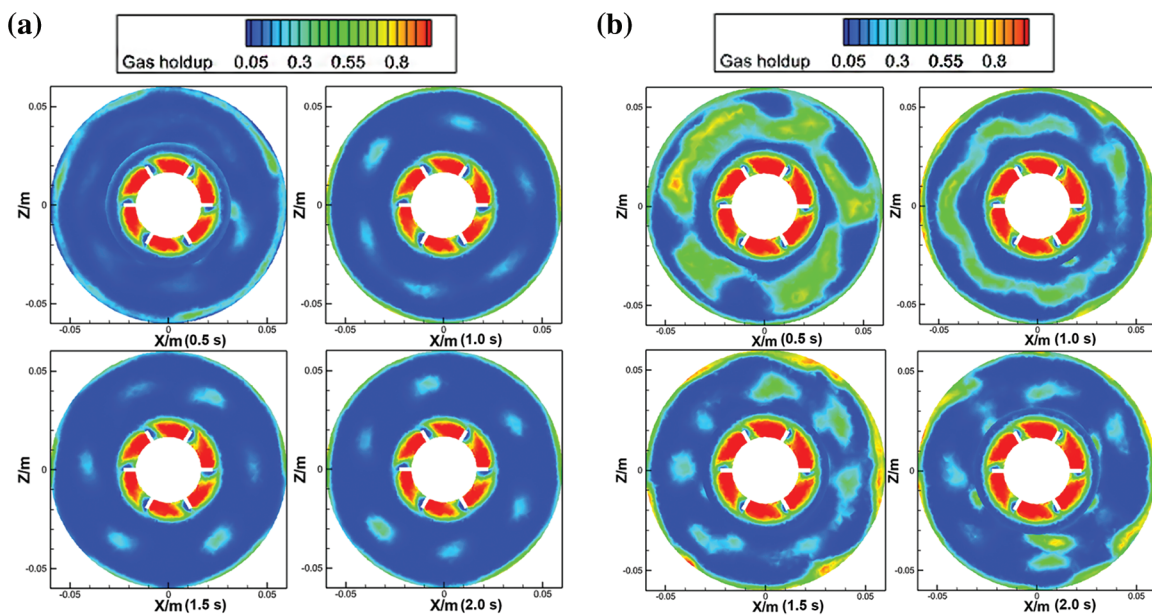
## 6 Results and Discussion

In a mechanical stirring system, there is multiphase flow, mass transfer is closely related to the interface area that affects the gas dispersion behavior, and the gas dispersion in the stirring tank is strongly influenced by the liquid flow pattern [27], stirring speed [28] and shape of the agitator [29]. To investigate the effect of the mechanical velocity and paddle shape on the flow field, a simulation was performed to simulate various conditions, which are discussed as follows.

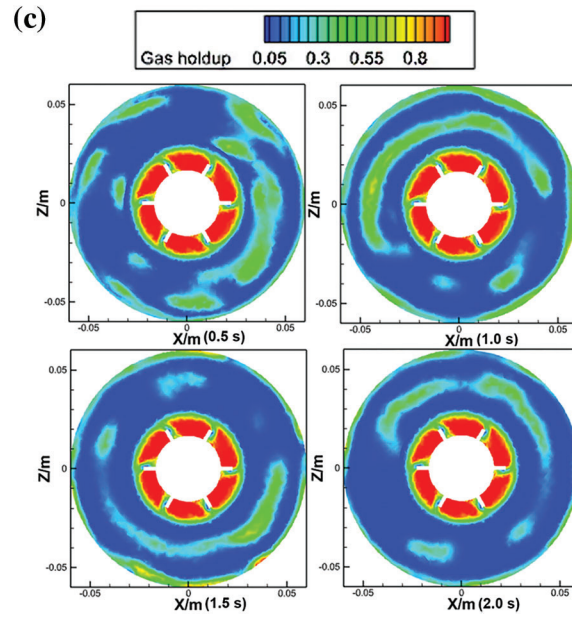
### 6.1 Effect of Mechanical Velocity on Flow Field

The rotating speeds used were 1000, 1250 and 1500 rpm. To observe and analyse the distribution of the flow field, the results were analysed in Tecplot360, and the reference surface was at  $y = -0.055$  and  $z = 0$  m.

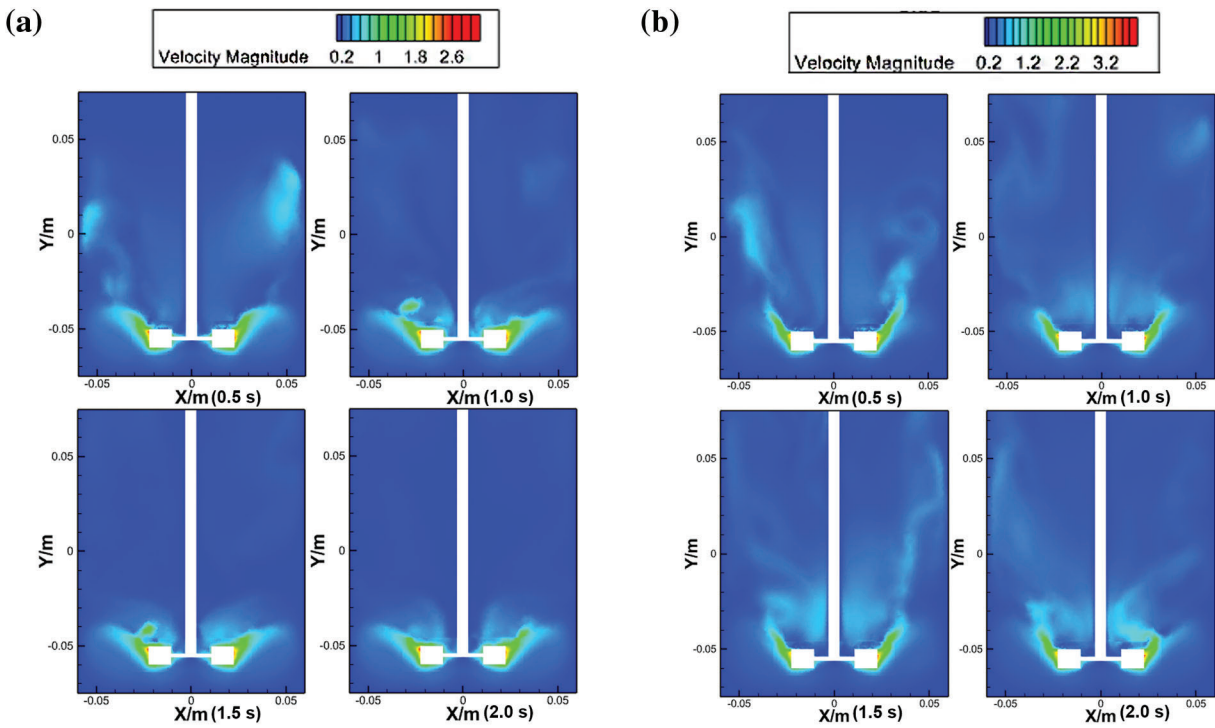
The results show the gas holdup at the cross section of  $y = -0.055$  m and the velocity magnitude at the cross section of  $z = 0$  m. When the velocity is 1000 rpm, as shown in Fig. 4a, the gas holdup tends to increase, and the range is mainly around the paddle as the stirring time increases. It can be easily observed that the gas holdup appears as a strip distribution at the beginning, as shown in Fig. 4a, and R1\_0.5 s but later becomes a circle distribution, as shown in Fig. 4a, and R1\_2.0 s. As the mechanical velocity increases to 1250 rpm, as shown in Fig. 4b, the range of gas holdup becomes wider, and almost the whole cross section becomes nearly the same. When the velocity increases from 1250 to 1500 rpm, the gas holdup presents a different trend from the previous trend, as shown in Fig. 4c, which is likely a ditch caused by too much speed. The results demonstrated that the gas holdup of the flow field is greater and the gas is more evenly distributed at the middle velocity in the cross section, so the velocity is beneficial for preparing tannin-based foam. The prediction of the gas distribution under various rotation speeds is similar to the experimental finding that the higher the rotation speed is, the wider the gas range [30].



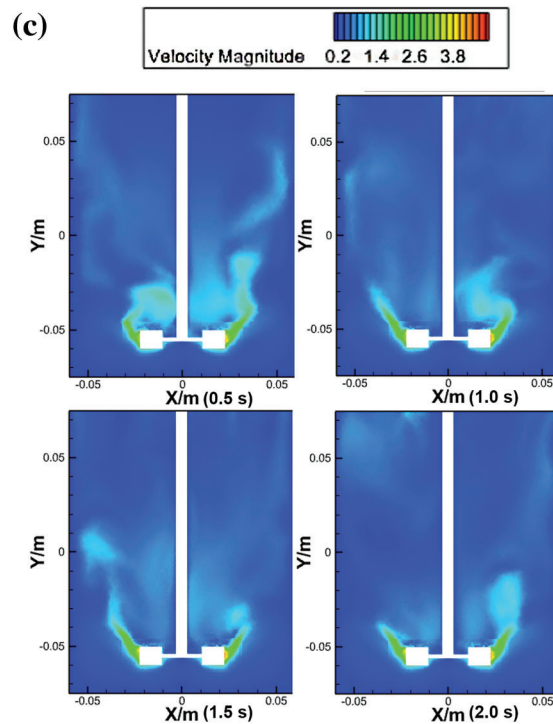
**Figure 4:** (Continued)



**Figure 4:** Gas holdup of the Rushton turbine at  $y = -0.055$  m. (a) R1 is 1000 rpm, (b) R2 is 1250 rpm, and (c) R3 is 1500 rpm



**Figure 5:** (Continued)



**Figure 5:** Velocity of the Rushton turbine at  $z = 0$  m. (a) R1 is 1000 rpm, (b) R2 is 1250 rpm, and (c) R3 is 1500 rpm

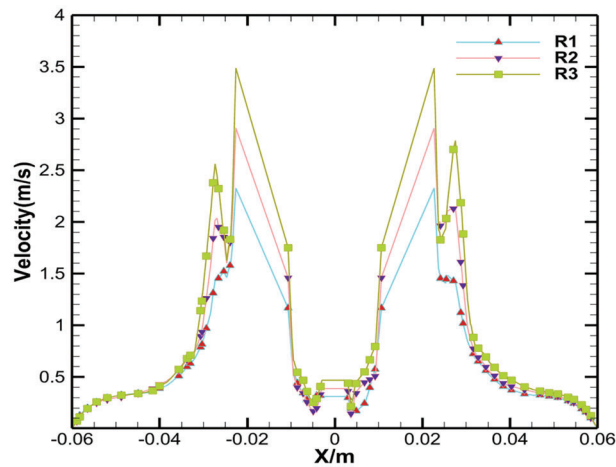
In the foaming process, as much air as possible should be mixed into the resin by mechanical whipping, so the larger the range of disturbance is, the greater the likelihood of mixing into more air and the greater the porosity of the resulting tannin-based foam. The results were analysed at the direction of  $z = 0$  m, as shown in Fig. 5. When the velocity is 1000 rpm, the range of disturbance is just around the blade, the tip of the blade has the greatest speed, and the top of the flow field remains almost constant, as shown in Fig. 5a. When the speed increases to 1250 rpm, the agitation range becomes wider, and the flow field has a larger velocity than ever before. At the same time, the top of the cross section presents disturbance, as exhibited in Fig. 5b. When the velocity increases from 1250 to 1500 rpm, the velocity of the flow field also increases, as shown in Fig. 5c. Furthermore, it can be concluded that the distribution of the velocity was mainly concentrated around the mixing shaft; therefore, the stirring speed and the fluid velocity were positively correlated, as shown in Figs. 6 and 7.

### 6.2 Effect of the Paddle Shape on the Flow Field

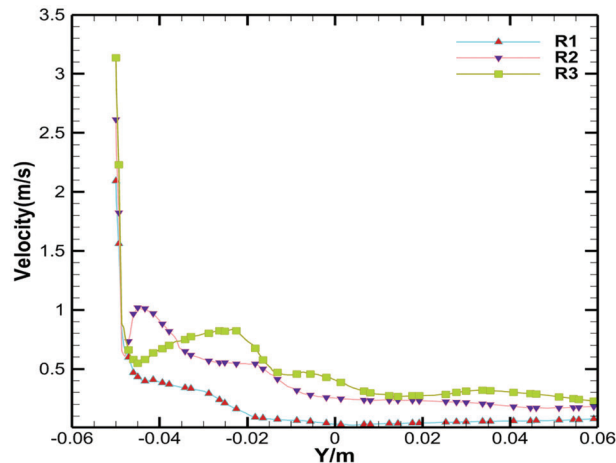
The eggbeater offers stirring in the flow field at three rotation speeds, which is the same as Rushton stirring. The results of eggbeater stirring are presented in Fig. 8. The gas holdup increases with increasing velocity, and the range of disturbance is wider.

Most of the gas is present at the bottom of the eggbeater; as the rotation speed increases, the gas gradually moves closer to the wall of the beaker, and the two rotation centers cause opposite flow fields, which is the reason for the mirror gas content. This can control the gas in the beaker to provide two symmetrical parts that can make the foam construction more stable. However, the symmetrical circles caused the resin surface to sag, which is not conducive to air mixing, under either of the two paddle shapes. Compared with the result of the Rushton turbine, when the rotation speed was 1000 rpm, as

shown in Fig. 4a, the eggbeater showed a narrow disturbance range, and there was no impact on the flow field except for the rotation domain; however, the disturbance of the Rushton turbine was wider, and the gas holdup was just around the top of the impeller outside of the rotation domain, as shown in Fig. 8a. As the stirring velocity increases to 1250 rpm, the Rushton turbine shows a wider range of disturbances in almost the whole cross section, as shown in Fig. 4b. At the same time, the gas appears outside the rotating region in the eggbeater paddle shape, and the position of the gas is on the axis of the tangent plane in the two rotating domains, which also presents a mirror gas distribution, as exhibited in Fig. 8b. When the rotation speed reached 1500 rpm, the distribution of gas content in the shape of an eggbeater paddle was wider than that at a stirring speed of 1250 rpm, but the trend of the gas distribution was the same as that before 1.5 s, rather than the result of the shape of the Rushton turbine paddle, which showed a more stable structure, as shown in Fig. 8c. The gas holdup in the shape of Rushton turbine impellers is larger than that in the eggbeater paddle shape in the whole flow field of the cross section. However, in the case of the two types of lower paddle shapes, the amount of tannin-based foaming precursor resin in the beaker was low inside and high outside due to stirring, as shown in Fig. 9a, but the paddle shape of the eggbeater alternated between high and low, as shown in Fig. 9b. This is because the tannin-based foaming precursor resin is removed by shearing, while the eggbeater has two rotation domains, which causes an alternating phenomenon.

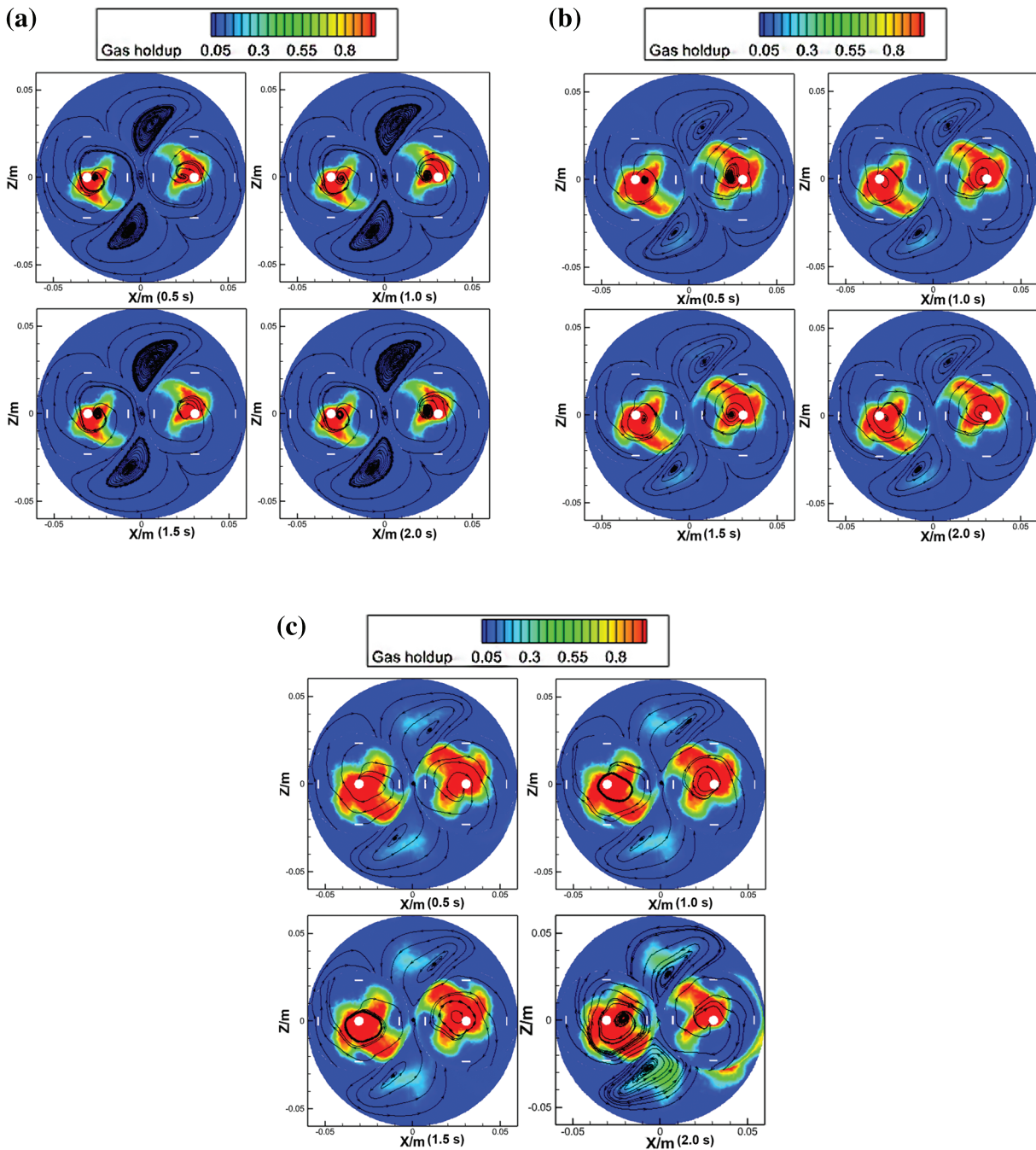


**Figure 6:** Velocity of the Rushton turbine in the X direction



**Figure 7:** Velocity of the Rushton turbine in the Y direction

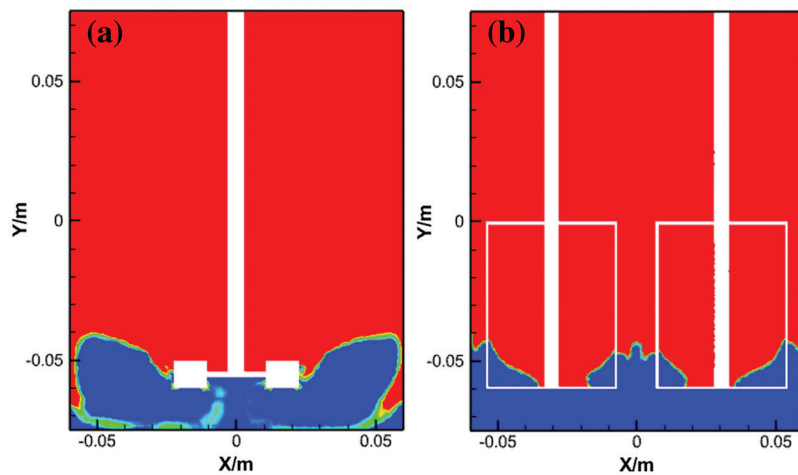




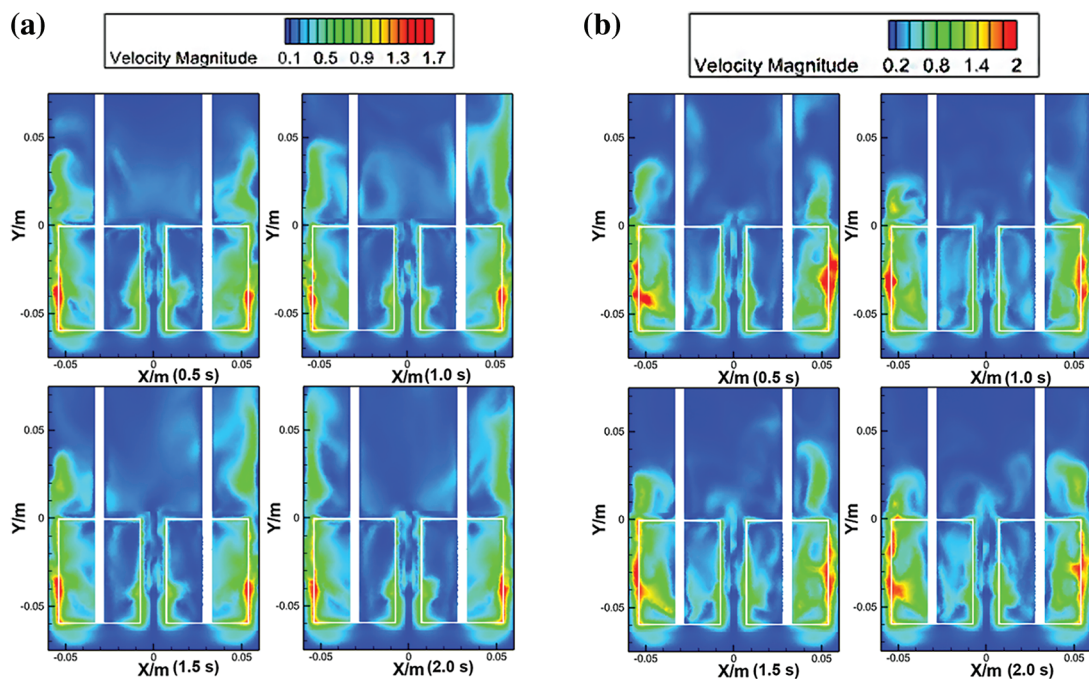
**Figure 8:** Gas holdup of the egg beater at  $y = -0.055$  m. (a) D1 is at 1000 rpm, (b) D2 is at 1250 rpm, (c) D3 is at 1500 rpm

The shape of the agitation has a significant impact on the velocity of the flow field. Fig. 10 shows the velocity of the flow field for the two shapes of impellers in the  $xy$  direction. When the rotation speed was 1000 rpm, as shown in Fig. 10a, the greatest velocity of the eggbeater is outside of the impellers, and the closer impellers exhibit a lower velocity in the flow field during the process of stirring. With the increase of stirring speed, it can be seen that the velocity of the flow field outside the impeller is significantly increased and the range is expanded as shown in Fig. 10b,c. In the flow field of the eggbeater paddle shape, the

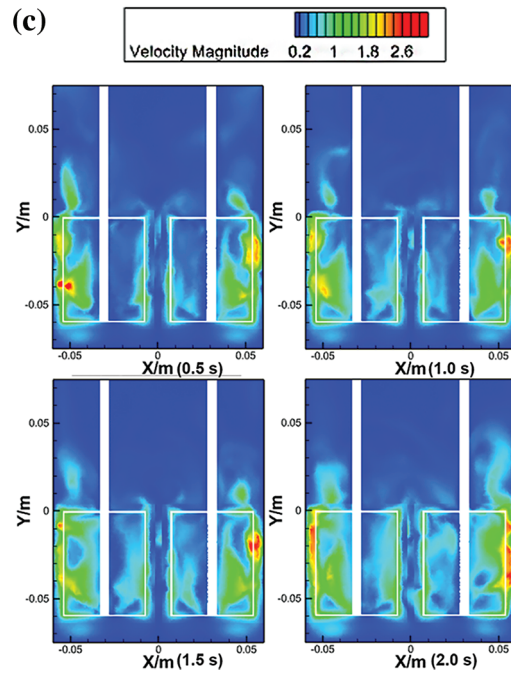
velocity magnitude is lower than that of the Rushton turbine paddle shape, and the distribution of the velocity affects the calculation results, as shown in Figs. 11 and 12. The interaction between the two regions will produce opposite forces, and some will be cancelled out, whereas some will participate in rotation. The shape of the Rushton turbine impellers moves around the axis, as shown in Fig. 4, but the cross section of the eggbeater is partitioned into four parts with the presence of a vortex in every single part, as shown in Fig. 8, which provides more gas to the resin. The function of a mechanical agitator is to mix the material evenly, which is different from the function of an eggbeater, which wraps air into a resin. Therefore, the structure of the paddle shape determines the application. During the foaming process, a large amount of air arrives and stays in the resin by stirring, so it is obvious that the use of an eggbeater is more suitable for the process.



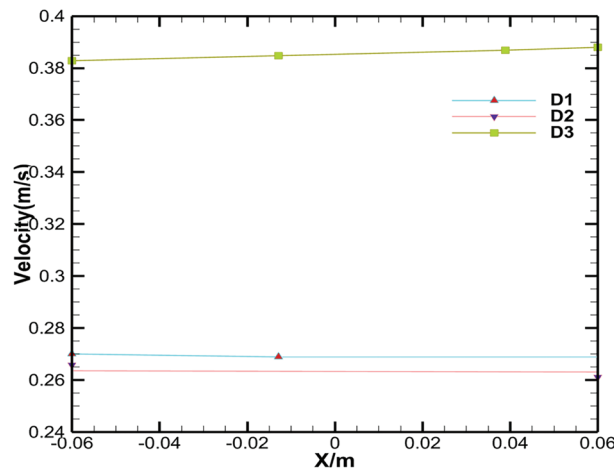
**Figure 9:** Gas holdup of the Rushton turbine and egg beater at  $z = 0$  m. (a) R1 is 1000 rpm, (b) D1 is 1000 rpm



**Figure 10:** (Continued)

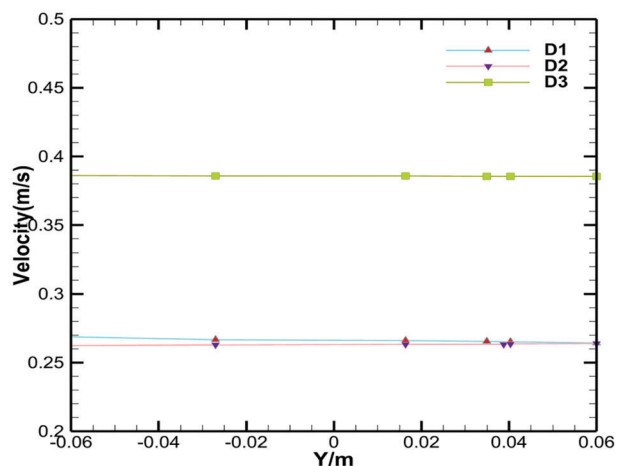


**Figure 10:** Velocity of the egg beater at  $z = 0$  m. (a) D1 is at 1000 rpm, (b) D2 is at 1250 rpm, (c) D3 is at 1500 rpm



**Figure 11:** Velocity of the egg beater in the X direction

The dead zones were generated by two-phase simulation, which is a significant indicator for evaluating the flow field. By comparing the dead zone in the case of two shapes of impellers and three kinds of mechanical stirring speeds, it can be concluded that the dead zone in the Rushton turbine is wider than that in the eggbeater, especially at the bottom of the simulation region, and in the radial direction. With the velocity increasing from 1000 to 1500 rpm, the dead zone in the Rushton turbine gradually decreases, but the egg beater presents the opposite phenomenon; that is, the two centers of rotation interact to cancel out some of the work leading to the dead zone. Therefore, the effect of the paddle shape and mechanical stirring speed is beneficial for the tannin-based foaming process.



**Figure 12:** Velocity of the egg beater in the Y direction

## 7 Conclusions

In this study, two kinds of paddle shapes and three stirring speeds were simulated, and the effects on the flow field were analysed in a tannin-based foaming precursor resin. Some concluding remarks can be outlined:

(1) Under stirring at different gradient speeds, it can be concluded that by increasing the rotation speed from 1000 to 1500 rpm, the range of shear action is extended, the gas phase distribution is more uniform, and there is more entrapped air; at the same time, the mixing of the liquid phase also increases. Due to the presence of eddy currents, the velocity of the outermost layer is slightly greater than that of the inner layer, which proves that the shape of the eggbeater is less likely to cause death far from the impellers in the tannin-based foaming precursor resin process, so the use of an egg whisker is more environmentally friendly and cost effective.

(2) Concerning the effect of the impeller pattern, the eggbeater has four vortexes in the cross section, but the Rushton turbine has only one vortex, which leads to a lower disturbance of the flow field in the paddle shape of the eggbeater. The work performed in the opposite direction was partially offset, which caused the flow field of the tannin-based foaming precursor to be more stable than that of the Rushton turbine impellers. The flow pattern is also different: the Rushton turbine has upper and lower loops, but the eggbeater has multiple cycles, and those cycles can improve the mixing efficiency of tannin-based foaming precursors.

(3) When the stirring speed is set at 1500 rpm, the Rushton turbine produces a ditch, which decreases the mass transfer and leads to unevenness in the structure. At the same time, the eggbeaters showed good stirring behaviour. It is better to work with the eggbeater set at a speed of 1500 rpm to obtain a stable and even flow field of tannin-based foaming precursor resin, meaning that the state of the flow field is vital to the structure of tannin-based foam.

The results have certain reference significance for the design and optimization of tannin-based foam foaming reaction processes. In future work, the shear angle and number of impellers and the axis of the agitator should be considered.

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