# A Sliding Mode Control Algorithm for Solving an Ill-posed Positive Linear System

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For the numerical solution of an ill-posed positive linear system we Abstract: combine the methods from invariant manifold theory and sliding mode control theory, developing an affine nonlinear dynamical system with a positive control force and with the residual vector as being a gain vector. This system is proven asymptotically stable to the zero residual vector by using an argument from the Lyapunov stability theory. We find that the system fast tends to the sliding surface and then moves with a sliding mode, such that the resultant sliding mode control algorithm (SMCA) is robust against large noise and stable to find the numerical solution of an ill-posed linear system. It is interesting that even under a random noise with an intensity  $10^{-5}$  we can obtain a quite accurate solution of the linear Hilbert problem with dimension n = 500. For this highly ill-conditioned problem the number of iterations is still smaller than 100. Numerical tests, including the inverse problems of backward heat conduction problem and Cauchy problems, confirm that the present SMCA has superior computational efficiency and accuracy even for a highly ill-conditioned linear equations system under a large noise.

**Keywords:** Ill-posed linear equations, Invariant manifold, Sliding mode control method, Asymptotically stable

### 1 Introduction

In this paper we propose a robust and easily-implemented algorithm to solve the following linear equations system:

$$\mathbf{B}\mathbf{x} = \mathbf{b},\tag{1}$$

where  $\mathbf{B} \in \mathbb{R}^{n \times n}$  is a given positive-definite matrix, and  $\mathbf{x} \in \mathbb{R}^n$  is an unknown vector to be determined from the input data  $\mathbf{b} \in \mathbb{R}^n$ . When **B** is severely ill-conditioned

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and **b** is perturbed by noise, we may encounter the problem that the numerical solution of Eq. (1) will deviate from the exact one to a great extent. Therefore, an algorithm compromises stability, efficiency and accuracy is desired.

There are several regularization techniques developed after the pioneering work of Tikhonov and Arsenin (1977). Previously, the author and his co-workers have developed several methods to solve the ill-posed linear problems: using the fictitious time integration method as a filter for ill-posed linear system [Liu and Atluri (2009a)], a modified polynomial expansion method [Liu and Atluri (2009b)], the non-standard group-preserving scheme [Liu and Chang (2009)], a vector regularization method [Liu, Hong and Atluri (2010)], the preconditioners and postconditioners generated from a transformation matrix, obtained by Liu, Yeih and Atluri (2009) for solving the Laplace equation with a multiple-scale Trefftz basis function-s, the relaxed steepest descent method [Liu (2011a, 2012a)], the optimal iterative algorithm [Liu and Atluri (2011a)], an optimally scaled vector regularization method [Liu (2012b)], the best vector iterative method [Liu (2012c)], a globally optimal iterative method [Liu (2012e)], the optimal tri-vector iterative methods [Liu (2013a, 2014a)], as well as an adaptive Tikhonov regularization method [Liu (2013b)].

It is known that iterative methods for solving the system of algebraic equations can be derived from the discretization of certain ODEs system [Bhaya and Kaszkurewicz (2006); Chehab and Laminie (2005); Liu and Atluri (2008)]. Particularly, the descent methods can be interpreted as the discretizations of gradient flows [Brown and Bartholomew-Biggs (1989); Helmke and Moore (1994)]. For a large scale system the main choice is using an iterative regularization algorithm, where a regularization parameter is presented by the number of iterations. The iterative method works if an early stopping criterion is used to prevent the reconstruction of noisy components in the approximate solutions.

Liu and Atluri (2008) have developed a very simple ODEs system for iteratively solving algebraic equations. After the work by Liu and Atluri (2008), there were several works applied the fictitious time integration method (FTIM) to solve engineering problems, e.g., [Liu and Atluri (2009a); Liu (2008a, 2009a, 2009b, 2009c, 2010); Chi, Yeih and Liu (2009), Ku, Yeih, Liu and Chi (2009); Chang and Liu (2009)]. In this paper we will develop an affine nonlinear ODEs system with a derived controller from the sliding mode control theory to accelerate the convergence speed in the solution of ill-posed linear problems.

### 2 An invariant manifold

There are several regularization methods to deal with Eq. (1) when **B** is ill-conditioned, of which the most prominent and best well understood ones are the Tikhonov method, the Landweber iteration method, and truncated singular value decomposition, all being linear and all being strongly convergent with an appropriate a priori parameter choice [Eicke, Louis and Plato (1990)]. In this paper we consider an iterative regularization method for Eq. (1) by investigating

$$\mathbf{r}(\mathbf{x}) = \mathbf{B}\mathbf{x} - \mathbf{b}.\tag{2}$$

We start from a continuous manifold [Absil, Baker and Gallivan (2007); Adler, Dedieu, Margulies, Martens and Shub (2002); Baker, Absil and Gallivan (2008); Smith (1994); Yang (2007)], defined in terms of the square-residual-norm of **r** and a function Q(t):

$$h(\mathbf{x},t) := \frac{1}{2}Q(t)\|\mathbf{r}(\mathbf{x})\|^2 = C.$$
(3)

Here, we let **x** be a function of a fictitious time-like variable *t*, and expect that in our algorithm Q(t) > 0 is an increasing function of *t*, and the residual error can be decreased with time. We let Q(0) = 1, and *C* is determined by the initial condition  $\mathbf{x}(0) = \mathbf{x}_0$  given by

$$C = \frac{1}{2} \|\mathbf{r}(\mathbf{x}_0)\|^2.$$
 (4)

Usually, C > 0, and C = 0 when the initial value  $\mathbf{x}_0$  is just the solution of Eq. (1). When C > 0 and Q > 0, the manifold defined by Eq. (3) is continuous and differentiable, and thus the following differential operation carried out on the manifold makes sense. For the requirement of "consistency condition", i.e.,  $\mathbf{x}(t)$  always on the manifold in time, we have

$$\frac{1}{2}\dot{Q}(t)\|\mathbf{r}(\mathbf{x})\|^2 + Q(t)(\mathbf{B}^{\mathrm{T}}\mathbf{r})\cdot\dot{\mathbf{x}} = 0,$$
(5)

which is obtained by taking the time differential of Eq. (3) with respect to *t* and considering  $\mathbf{x} = \mathbf{x}(t)$ .

The governing equation of  $\mathbf{x}$  cannot be uniquely determined by Eq. (5); however, we suppose that  $\mathbf{x}$  is governed by a gradient-flow, like that for the steepest-descent method (SDM):

$$\dot{\mathbf{x}} = -\lambda \frac{\partial h}{\partial \mathbf{x}} = -\lambda Q(t) \mathbf{B}^{\mathrm{T}} \mathbf{r},\tag{6}$$

where  $\lambda$  is to be determined. Inserting Eq. (6) into Eq. (5) we can solve

$$\lambda = \frac{\dot{Q}(t) \|\mathbf{r}\|^2}{2Q^2(t) \|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2}.$$
(7)

Thus by inserting the above  $\lambda$  into Eq. (6) we can obtain an ODEs system for **x**:

$$\dot{\mathbf{x}} = -q(t) \frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \mathbf{B}^{\mathrm{T}}\mathbf{r},\tag{8}$$

where

$$q(t) := \frac{Q(t)}{2Q(t)}.$$
(9)

If Q(t) can be guaranteed to be an increasing function of t, we have an absolutely convergent numerical method to solve Eq. (1) shown as follows:

$$\|\mathbf{r}(\mathbf{x})\|^{2} = \frac{2C}{Q(t)}.$$
(10)

When *t* is large enough the above equation will enforce the residual error  $||\mathbf{r}(\mathbf{x})||$  tending to zero, and meanwhile the solution of Eq. (1) is obtained approximately.

### **3** The sliding mode control method

Equation (8) is an ODEs system defined on the invariant manifold in Eq. (10). Based on this equation, Liu (2011a) has derived the relaxed steepest descent method (RSDM) for solving linear system (1). Although the original design of the numerical algorithm is for the purpose of keeping the orbit of  $\mathbf{x}$  on the invariant manifold in Eq. (10), but in practice the RSDM cannot satisfy this requirement. For this purpose we can insert a *compensated controller*  $\tilde{u}$  in Eq. (8):

$$\dot{\mathbf{x}} = -\frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \mathbf{B}^{\mathrm{T}}\mathbf{r} + \tilde{u}\mathbf{r},\tag{11}$$

such that a suitable design of  $\tilde{u}$  can help us to achieve this purpose. In above we let q = 1 and **r** is taken as a gain vector.

The theory of sliding mode control for system (11) is designed an invariant hypersurface [Utkin (1978, 1992); Liu (2014b)]:

$$s(\mathbf{x}) = 0 \tag{12}$$

in the state space, such that by keeping the orbit of  $\mathbf{x}$  on the surface we require

$$\dot{s}(\mathbf{x}) = \frac{\partial s}{\partial \mathbf{x}} \cdot \dot{\mathbf{x}} = 0.$$
<sup>(13)</sup>

Here we can set up the most simple sliding surface by

$$s = \mathbf{a} \cdot \mathbf{r} = 0, \tag{14}$$

where **a** is a constant vector. Hence, we have

$$\dot{s} = \mathbf{a} \cdot \dot{\mathbf{r}} = \mathbf{a} \cdot (\mathbf{B} \dot{\mathbf{x}}) = 0. \tag{15}$$

Inserting Eq. (11) for  $\dot{\mathbf{x}}$  we can derive

$$-\frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2}\mathbf{a}\cdot(\mathbf{A}\mathbf{r}) + \tilde{u}\mathbf{a}\cdot(\mathbf{B}\mathbf{r}) = 0,$$
(16)

where

$$\mathbf{A} := \mathbf{B}\mathbf{B}^{\mathrm{T}}.\tag{17}$$

From Eq. (16) we can solve

$$\tilde{u} = \frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \frac{\mathbf{a} \cdot (\mathbf{A}\mathbf{r})}{\mathbf{a} \cdot (\mathbf{B}\mathbf{r})},\tag{18}$$

and then by Eq. (11) we have

$$\dot{\mathbf{x}} = -\frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \mathbf{B}^{\mathrm{T}}\mathbf{r} + \frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \frac{\mathbf{a} \cdot (\mathbf{A}\mathbf{r})}{\mathbf{a} \cdot (\mathbf{B}\mathbf{r})} \mathbf{r};$$
(19)

however, this ODEs system may be unstable. Thus we recast Eq. (19) to be

$$\dot{\mathbf{x}} = -\frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \mathbf{B}^{\mathrm{T}}\mathbf{r} - u\mathbf{r},\tag{20}$$

where

$$u = |\tilde{u}| = \frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}\|^2} \left| \frac{\mathbf{a} \cdot (\mathbf{A}\mathbf{r})}{\mathbf{a} \cdot (\mathbf{B}\mathbf{r})} \right| > 0.$$
(21)

In the control theory, Eq. (20) is a form of an *affine nonlinear system*, which is linear with the control force u(t), while the residual vector **r** is acting as a gain vector.

Now we prove that the above dynamical system is asymptotically stable, tending to the zero dynamics of the residual vector  $\mathbf{r}$ . Consider the following Lyapunov function:

$$V = \frac{1}{2} \|\mathbf{r}\|^2 \ge 0.$$
(22)

Taking the derivative of V and inserting  $\dot{\mathbf{r}} = \mathbf{B}\dot{\mathbf{x}}$  and Eq. (20) we can derive

$$\dot{V} = \mathbf{r} \cdot \dot{\mathbf{r}} = \mathbf{r} \cdot (\mathbf{B}\dot{\mathbf{x}})$$
  
=  $-\frac{\|\mathbf{r}\|^2}{\|\mathbf{B}^T\mathbf{r}\|^2} \mathbf{r} \cdot (\mathbf{A}\mathbf{r}) - u\mathbf{r} \cdot (\mathbf{B}\mathbf{r}) < 0,$  (23)

where we have used the positiveness of **A**, **B** and *u*. As compared with Eq. (8), the new ODEs system (20) possesses an extra term  $u\mathbf{r}$  to enforce the dynamics to the zero point of  $\mathbf{r} = \mathbf{0}$ , which is gained by a sliding effect with the aid of a gain vector  $\mathbf{r}$  and a positive control force *u*.

Given an initial point  $\mathbf{r}_0 = \mathbf{B}\mathbf{x}_0 - \mathbf{b}$  we can construct the constant vector  $\mathbf{a}$  by an arbitrarily given vector  $\mathbf{a}_0$ , say  $\mathbf{a}_0 = (1, ..., 1)^T$ :

$$\mathbf{a} = \mathbf{a}_0 - \frac{\mathbf{a}_0 \cdot \mathbf{r}_0}{\|\mathbf{r}_0\|^2} \mathbf{r}_0; \tag{24}$$

it is obvious that  $\mathbf{a} \cdot \mathbf{r}_0 = 0$ , such that  $s_0 = \mathbf{a} \cdot \mathbf{r}_0 = 0$ .

We simply use the Euler scheme to integrate Eq. (20), and thus one can derive the following sliding mode control algorithm (SMCA):

(i) Give a stepsize *h*, an initial  $\mathbf{x}_0$ ,  $\mathbf{r}_0 = \mathbf{B}\mathbf{x}_0 - \mathbf{b}$  and  $\mathbf{a}_0 = (1, \dots, 1)^T$ , and compute  $\mathbf{a} = \mathbf{a}_0 - (\mathbf{a}_0 \cdot \mathbf{r}_0)\mathbf{r}_0 / \|\mathbf{r}_0\|^2$ .

(ii) For k = 0, 1, 2, ..., we repeat the following computations:

$$\mathbf{r}_{k} = \mathbf{B}\mathbf{x}_{k} - \mathbf{b},$$

$$u_{k} = \frac{\|\mathbf{r}_{k}\|^{2}}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}_{k}\|^{2}} \left| \frac{\mathbf{a} \cdot (\mathbf{A}\mathbf{r}_{k})}{\mathbf{a} \cdot (\mathbf{B}\mathbf{r}_{k})} \right|,$$

$$\mathbf{x}_{k+1} = \mathbf{x}_{k} - \frac{h\|\mathbf{r}_{k}\|^{2}}{\|\mathbf{B}^{\mathrm{T}}\mathbf{r}_{k}\|^{2}} \mathbf{B}^{\mathrm{T}}\mathbf{r}_{k} - hu_{k}\mathbf{r}_{k}.$$
(25)

If  $\mathbf{x}_{k+1}$  converges according to a given stopping criterion  $\|\mathbf{r}_{k+1}\| < \varepsilon$ , then stop; otherwise, go to step (ii).

#### 4 Numerical examples

In order to assess the performance of the newly developed method of the sliding mode control algorithm (SMCA), let us investigate the following examples. Some

results are compared with that obtained by the relaxed steepest descent method (RS-DM), conjugate gradient method (CGM), the optimal iterative algorithm driven by an optimal descent vector (OIA/ODV), and the globally optimal iterative algorithm (GOIA).

# 4.1 Hilbert problems

Finding an *n*-degree polynomial function  $p(x) = a_0 + a_1x + ... + a_nx^n$  to best match a continuous function f(x) in the interval of  $x \in [0, 1]$ :

$$\min_{\deg(p) \le n} \int_0^1 [f(x) - p(x)]^2 dx,$$
(26)

leads to a problem governed by Eq. (1), where **B** is the  $(n+1) \times (n+1)$  Hilbert matrix, defined by

$$B_{ij} = \frac{1}{i+j-1},$$
(27)

**x** is composed of the n + 1 coefficients  $a_0, a_1, \ldots, a_n$  appeared in p(x), and

$$\mathbf{b} = \begin{bmatrix} \int_0^1 f(x)dx \\ \int_0^1 xf(x)dx \\ \vdots \\ \int_0^1 x^n f(x)dx \end{bmatrix}$$
(28)

is uniquely determined by the function f(x).

The Hilbert matrix is a famous example of highly ill-conditioned matrices. Eq. (1) with the coefficient matrix **B** having a large condition number usually displays that an arbitrarily small perturbation of data on the right-hand side may lead to an arbitrarily large perturbation to the solution on the left-hand side.

In this example we consider a highly ill-conditioned linear equation (1) with **B** given by Eq. (27). The ill-posedness of Eq. (1) with the above **B** increases very fast with an exponential growth with n.

# $4.1.1 \quad n = 9$

In order to compare the numerical solutions with exact solutions we suppose that  $x_1 = \ldots = x_n = 1$  to be the exact one, and then by Eq. (27) we have

$$b_i = \sum_{j=1}^n \frac{1}{i+j-1} + \sigma R(i),$$
(29)

where we consider noise being imposed on the data with random numbers  $R(i) \in [-1, 1]$ .

We first calculate this problem for the case with n = 9 and  $\sigma = 0$ . The resulting linear equation is highly ill-conditioned, since the condition number is quite large, up to  $4.93 \times 10^{11}$ .

In the computation by the SMCA we have fixed h = 0.3. Starting from  $x_1 = ... = x_9 = 0.5$ , and with a stopping criterion  $\varepsilon = 10^{-8}$  we find that SMCA converges with 1166 iterations, where the maximum error is  $1.03 \times 10^{-3}$ .

The RSDM with  $\gamma = 0.06$  does not converge within 5000 iterations; however the numerical solution is close to the exact solution as shown in Table 1 with the maximum error being  $1.44 \times 10^{-3}$ ; wherein, for the purpose of comparison, the values obtained by the conjugate gradient method (CGM) are also listed. The maximum error of CGM is  $7.85 \times 10^{-3}$ .

 $4.1.2 \quad n = 50$ 

Let us increase the ill-posedness of the linear Hilbert problem to n = 50. For this problem the condition number is about  $1.1748 \times 10^{19}$ . We first consider a constant solution  $x_1 = \ldots = x_{50} = 1$ . The convergence criterion  $\varepsilon$  is fixed to be  $\varepsilon = 10^{-6}$  for case 1 with noise-free  $\sigma = 0$ , and case 2 with  $\varepsilon = 10^{-3}$  and with a noise  $\sigma = 10^{-4}$ . Starting from the initial condition  $x_i = 0.5$ ,  $i = 1, \ldots, 50$ , for case 1 the SMCA with h = 0.1 converges with 1184 iterations as shown in Fig. 1(a) for the residual error, Fig. 1(b) for the control force and Fig. 1(c) for the numerical error, of which the maximum error is 0.00392. Very accurate result is obtained. From Fig. 2(a) we can observe that the iterative dynamics approaches to the sliding surface with s = 0 very quickly. For case 2 the SMCA with h = 0.1 converges with 116 iterations as shown in Fig. 2(b) for the sliding quantity *s*, Fig. 2(c) for the residual error and the control force. The maximum error is 0.037 as shown in Fig. 1(c) by the dashed line.

4.1.3 n = 200

Let us consider n = 200 and with a noise  $\sigma = 10^{-5}$ . In the computation of this noisy problem by the SMCA, we fix h = 0.35 and  $\varepsilon = 10^{-3}$ , where we show the residual error in Fig. 3(a) and the SMCA converges with 82 iterations. From Fig. 3(b) we can see that the sliding mode is obtained very fast. The numerical solution obtained by the SMCA converges to the exact solution very accurately as shown in Fig. 3(c) with the maximum error being 0.0258. We also apply the CGM to this problem under the same noise. Under the above convergence criterion the CGM converges very fast with 9 steps, and the maximum error as shown in Fig. 3(c) is 0.0993. It can be seen that the accuracy of SMCA is better than CGM.





Figure 1: For example 1 of Hilbert problem with n = 50, (a) the residual and (b) the control force for noise free case, and (c) comparing the numerical errors.



Figure 2: For example 1 of Hilbert problem with n = 50, comparing the sliding quantities for (a) noise free and (b) under a noise, and (c) the control force and residual for the noised case.



Figure 3: For example 1 of Hilbert problem with n = 200, (a) the residual and (b) the sliding quantity for noised case, and (c) comparing the numerical errors.



Figure 4: For example 1 of Hilbert problem with n = 500, (a) the residual and (b) the sliding quantity for noised case, and (c) comparing the numerical errors.

# $4.1.4 \quad n = 500$

Let us consider n = 500 and with a noise  $\sigma = 10^{-5}$ . In the computation by the SMCA, we fix h = 0.4 and  $\varepsilon = 10^{-3}$ , where we show the residual error in Fig. 4(a) and the SMCA can converge with 96 iterations. From Fig. 4(b) we can see that the sliding mode is obtained very fast. The numerical solution of SMCA is very accurate as shown in Fig. 4(c) with the maximum error being 0.024. The maximum error of CGM as shown in Fig. 4(c) is large up to 0.1973. It can be seen that the accuracy of SMCA is much better than CGM. From the above cases for this highly ill-posed Hilbert problem we can observe that the present SMCA is robust and effective, which is insensitive to the degree of the ill-posedness. The accuracy of the SMCA is much better than other methods.

#### 4.2 Example 2

When the backward heat conduction problem (BHCP) is considered in a spatial interval of  $0 < x < \ell$  by subjecting to the boundary conditions at two ends of a slab:

$$u_t(x,t) = \alpha u_{xx}(x,t), \ \ 0 < t < T, \ \ 0 < x < \ell,$$
(30)

$$u(0,t) = u_0(t), \ u(\ell,t) = u_\ell(t), \tag{31}$$

we solve *u* under a final time condition:

$$u(x,T) = u^T(x).$$
(32)

The fundamental solution to Eq. (30) is given as follows:

$$K(x,t) = \frac{H(t)}{2\sqrt{\alpha\pi t}} \exp\left(\frac{-x^2}{4\alpha t}\right),\tag{33}$$

where H(t) is the Heaviside function.

The method of fundamental solutions (MFS) has a broad application in engineering computations. However, the MFS has a serious drawback in that the resulting system of linear equations is always highly ill-conditioned, when the number of source points is increased [Golberg and Chen (1996)], or when the distances of source points are increased [Chen, Cho and Golberg (2006)].

In the MFS the solution of *u* at the field point  $\mathbf{z} = (x, t)$  can be expressed as a linear combination of the fundamental solutions  $U(\mathbf{z}, \mathbf{s}_i)$ :

$$u(\mathbf{z}) = \sum_{j=1}^{N} c_j U(\mathbf{z}, \mathbf{s}_j), \ \mathbf{s}_j = (\eta_j, \tau_j) \in \Omega^c,$$
(34)

|           |         | F       | C       |          |          | -       |            |          |          |
|-----------|---------|---------|---------|----------|----------|---------|------------|----------|----------|
| Solutions | $x_1$   | $x_2$   | $x_3$   | $x_4$    | $\chi_5$ | $9\chi$ | <i>X</i> 7 | $8\chi$  | $6\chi$  |
| Exact     | 1.0     | 1.0     | 1.0     | 1.0      | 1.0      | 1.0     | 1.0        | 1.0      | 1.0      |
| CGM       | 0.99966 | 1.00396 | 0.99251 | 0.99769  | 1.00371  | 1.00606 | 1.00439    | 0.999457 | 0.992143 |
| RSDM      | 1.00000 | 0.99986 | 1.00087 | 88866'0  | 0.99928  | 1.00058 | 1.00128    | 1.00066  | 95866.0  |
| SMCA      | 1.00000 | 0.99975 | 1.00099 | 0.999912 | 0.99923  | 1.00029 | 1.00099    | 1.00063  | 16866.0  |
|           |         |         |         |          |          |         |            |          |          |

Table 1: Comparing the numerical results for Example 1 with different methods

where *N* is the number of source points,  $c_j$  are unknown coefficients, and  $s_j$  are source points located in the complement  $\Omega^c$  of  $\Omega = [0, \ell] \times [0, T]$ . For the heat conduction equation we have the basis functions:

$$U(\mathbf{z},\mathbf{s}_j) = K(x - \eta_j, t - \tau_j).$$
(35)

It is known that the location of source points in the MFS has a great influence on the accuracy and stability. In a practical application of MFS to solve the BHCP, the source points are uniformly located on two straight lines parallel to the *t*-axis and not over the final time, which was adopted by Hon and Li (2009) and Liu (2011b), showing a large improvement than the line location of source points below the initial time. After imposing the boundary conditions and the final time condition on Eq. (34) we can obtain a linear equations system:

$$\mathbf{V}\mathbf{x} = \mathbf{b}_1,\tag{36}$$

where

$$V_{ij} = U(z_i, s_j), \ \mathbf{x} = (c_1, \cdots, c_N)^{\mathrm{T}}, \mathbf{b}_1 = (u_\ell(t_i), \ i = 1, \dots, m_1; u^{\mathrm{T}}(x_j), \ j = 1, \dots, m_2; u_0(t_k), \ k = m_1, \dots, 1)^{\mathrm{T}}.$$
(37)

The number  $n = 2m_1 + m_2$  of collocation points does not necessarily equal to the number *N* of source points. By applying the SMCA to solve the above problem we solve a normal equation:

$$\mathbf{V}^{\mathrm{T}}\mathbf{V}\mathbf{x} = \mathbf{V}^{\mathrm{T}}\mathbf{b}_{1}.$$
(38)

Since the BHCP is highly ill-posed, the ill-conditioning of the matrix  $\mathbf{B} = \mathbf{V}^{\mathrm{T}}\mathbf{V}$  in Eq. (38) is serious. To overcome the ill-posedness of Eq. (38) we can employ the SMCA to solve this problem. Here we compare the numerical solution with an exact solution:

$$u(x,t) = \cos(\pi x) \exp(-\pi^2 t).$$

For the case with T = 1 the value of final data is in the order of  $10^{-4}$ , which is small in a comparison with the value of the initial temperature  $u_0(x) = \cos(\pi x)$  to be retrieved, which is O(1).

We apply the SMCA to solve Eq. (38) with h = 0.05 and  $\varepsilon = 10^{-4}$ , and a relative random noise with an intensity  $\sigma = 10\%$  is added on the final time data. The SMCA converges with 899 iterations as shown in Fig. 5(a), while the sliding quantity *s* is shown in Fig. 5(b). The numerical error as shown in Fig. 5(c) by the solid line is smaller than 0.0438. For the purpose of comparison we also apply the RSDM with  $\gamma = 0.05$  to solve this problem; however, over 10000 iterations it does not converge as shown in Fig. 5(a) by the dashed line for the residual error. However, the result as shown in Fig. 5(c) by the dashed line for numerical error is smaller than 0.05.



Figure 5: For example 2, (a) comparing the residual errors obtained by RSDM and SMCA, (b) the sliding quantities o SMCA, and (c) comparing the numerical errors.

#### 4.3 Example 3

We solve the Cauchy problem of the Laplace equation under overspecified boundary conditions:

$$\Delta u = u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta} = 0, \ r < \rho, \ 0 \le \theta \le 2\pi,$$
(39)

$$u(\rho, \theta) = h(\theta), \ 0 \le \theta \le \pi, \tag{40}$$

$$u_n(\rho,\theta) = g(\theta), \ 0 \le \theta \le \pi, \tag{41}$$

where  $h(\theta)$  and  $g(\theta)$  are given functions, and  $\rho = \rho(\theta)$  is a given contour to describe the boundary shape. The contour in the polar coordinates is specified by  $\Gamma = \{(r, \theta) | r = \rho(\theta), 0 \le \theta \le 2\pi\}$ , which is the boundary of the problem domain  $\Omega$ , and *n* denotes the normal direction.

In the potential theory, it is well known that the method of fundamental solutions (MFS) can be used to solve the Laplace equation when a fundamental solution is known. In the MFS the trial solution of *u* at the field point  $\mathbf{z} = (r \cos \theta, r \sin \theta)$  can be expressed as a linear combination of the fundamental solutions  $U(\mathbf{z}, \mathbf{s}_i)$ :

$$u(\mathbf{z}) = \sum_{j=1}^{n} c_j U(\mathbf{z}, \mathbf{s}_j), \ \mathbf{s}_j \in \Omega^c,$$
(42)

where *n* is the number of source points,  $c_j$  are the unknown coefficients,  $\mathbf{s}_j$  are the source points, and  $\Omega^c$  is the complementary set of  $\Omega$ . For the Laplace equation (39) we have the fundamental solutions:

$$U(\mathbf{z},\mathbf{s}_j) = \ln r_j, \ r_j = \|\mathbf{z} - \mathbf{s}_j\|.$$
(43)

In the practical application of MFS, frequently the source points are uniformly located on a circle with a radius R, such that after imposing the boundary conditions (40) and (41) on Eq. (42) we can obtain a linear equations system:

$$\mathbf{V}\mathbf{x} = \mathbf{b}_1,\tag{44}$$

where

$$\mathbf{z}_{i} = (z_{i}^{1}, z_{i}^{2}) = (\boldsymbol{\rho}(\theta_{i}) \cos \theta_{i}, \boldsymbol{\rho}(\theta_{i}) \sin \theta_{i}),$$
  

$$\mathbf{s}_{j} = (s_{j}^{1}, s_{j}^{2}) = (R \cos \theta_{j}, R \sin \theta_{j}),$$
  

$$V_{ij} = \ln \|\mathbf{z}_{i} - \mathbf{s}_{j}\|, \text{ if } i \text{ is odd},$$
  

$$V_{ij} = \frac{\boldsymbol{\eta}(\theta_{i})}{\|\mathbf{z}_{i} - \mathbf{s}_{j}\|^{2}} \left(\boldsymbol{\rho}(\theta_{i}) - s_{j}^{1} \cos \theta_{i} - s_{j}^{2} \sin \theta_{i} - \frac{\boldsymbol{\rho}'(\theta_{i})}{\boldsymbol{\rho}(\theta_{i})} \left[s_{j}^{1} \sin \theta_{i} - s_{j}^{2} \cos \theta_{i}\right]\right), \text{ if } i \text{ is even},$$
  

$$\mathbf{x} = (c_{1}, \dots, c_{n})^{\mathrm{T}}, \ \mathbf{b}_{1} = (h(\theta_{1}), g(\theta_{1}), \dots, h(\theta_{m}), g(\theta_{m}))^{\mathrm{T}},$$
(45)

in which n = 2m, and

$$\eta(\theta) = \frac{\rho(\theta)}{\sqrt{\rho^2(\theta) + [\rho'(\theta)]^2}}.$$
(46)

A noise with an intensity  $\sigma = 10\%$  is imposed on the given data. We fix n = 60and use a circle with a constant radius R = 14 to distribute the source points. By applying the SMCA to solve the normal equation (38) obtained from Eq. (44) we fix h = 0.3. Under the convergence criterion  $\varepsilon = 10^{-2}$ , the SMCA converges with 7308 iterations and the residual error is shown in Fig. 6(a), while the control force is shown in Fig. 6(b) and the sliding quantity *s* is plotted in Fig. 7(a). Along the lower half contour  $\rho(\theta) = \sqrt{10 - 6\cos(2\theta)}, \pi \le \theta < 2\pi$ , in Fig. 6(c) we compare the numerical solution with the exact data given by  $u = \rho^2 \cos(2\theta), \pi \le \theta < 2\pi$ , of which the maximum error is found to be 0.121.

For the purpose of comparison we also apply the OIA/ODV developed by Liu and Atluri (2011a) and the GOIA developed by Liu (2012d) to solve this problem under the same conditions. In Fig. 7(b) we compare the relative residuals obtained by these three methods, while in Fig. 7(c) we compare the numerical errors, where the maximum error obtained by the SMCA is 0.121, the maximum error by the GOIA is 0.141, and the maximum error by the OIA/ODV is 0.148. Both the GOIA and OIA/ODV do not converge within 10000 iterations under the above convergence criterion  $\varepsilon = 10^{-2}$ . We can say that the performance of SMCA is better than GOIA and OIA/ODV.

### 4.4 Example 4

In this example we consider an inverse Cauchy problem of the following biharmonic equation:

$$\Delta^2 u = 0, \quad (x, y) \in \Omega, \tag{47}$$

where  $\Omega$  is an interior domain in the plane. This inverse Cauchy problem of biharmonic equation is under an incomplete set of data given by

$$u(\rho, \theta) = h(\theta), \ u_n(\rho, \theta) = g(\theta), \ 0 \le \theta \le 2\beta\pi.$$
(48)

When  $\beta = 1$  we recover to the direct problem. Here we let  $\beta < 1$  and do not use the overspecified data, such that the present problem is a Cauchy problem with an incomplete set of given data.

For the purpose of comparison we suppose that the exact solution is

$$u(x,y) = x^3 + y^3,$$



Figure 6: For example 3, (a) the residual error obtained by SMCA, (b) the control force, and (c) comparing the numerical solution with the exact one.



Figure 7: For example 3, (a) the sliding quantity obtained by SMCA, (b) comparing the relative residuals, and (c) comparing the numerical errors.



Figure 8: For example 4 of an inverse Cauchy problem of biharmonic equation, comparing the numerical solutions obtained by the SMCA with exact ones.

and the domain is defined by

$$\rho(\theta) = \sqrt{26 - 10\cos(4\theta)}.\tag{49}$$

For this example we apply the following Trefftz method [Liu (2008b)] to solve the inverse Cauchy problem of biharmonic equation:

$$u(r,\theta) = a_0 + \sum_{k=1}^{m} \left[ a_k \left( \frac{r}{R_0} \right)^k \cos k\theta + b_k \left( \frac{r}{R_0} \right)^k \sin k\theta \right] + c_0 r^2 + \sum_{k=1}^{m} \left[ c_k \left( \frac{r}{R_0} \right)^{k+2} \cos k\theta + d_k \left( \frac{r}{R_0} \right)^{k+2} \sin k\theta \right].$$
(50)

Under the following parameters  $R_0 = 200$ , m = 20,  $\beta = 0.8$  and  $\sigma = 0.01$  we apply the SMCA to solve this Cauchy problem, of which the results are compared with the exact solutions in Fig. 8, where  $v = \Delta u$ . The results obtained by the SMCA are quite accurate.

### 5 Conclusions

In order to tackle the problem of the discrepancy of iteration orbit from the invariant manifold, which causes a slow convergence, in this new invariant-manifold based theory we introduce a sliding mode control method by adding an extra controller with the residual vector as being a gain vector into the evolution equation. The new dynamics is proven to be asymptotically stable towards the zero point of the residual vector, which means that we can quickly find the real solution of a positive linear system by the sliding mode control algorithm (SMCA). In several numerical examples we have observed the sliding behavior with a fast sliding phase, and by comparing with exact solutions, the SMCA can work very effectively for the highly ill-conditioned linear equations system under a large noisy perturbation. We have obtained very accurate solution of the linear Hilbert problem with dimension n = 500, and the number of iterations is still smaller than 100. For the inverse problems the SMCA is also effective, although the noise being imposed on the input data is large up to 10%.

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