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Absolute Stability of Chaotic Asynchronous Multi-Interactions Schemes for Solving ODE

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Abstract: Multi Interaction Systems, used in the context of Virtual Reality, are dedicated to real-time interactive simulations. They open the way to the *in virtuo experimentation*, especially useful in the domain of biochemical kinetics. To this purpose, chaotic and asynchronous scheduling of autonomous processes is based upon desynchronization of phenomena involved in the system. It permits interactivity, especially the capability to add or remove phenomena in the course of a simulation. It provides methods of resolution of ordinary differential systems and partial derivative equations. Proofs of convergence for these methods have been established, but the problem of absolute stability, although it is crucial when considering multiscale or stiff problems, has not yet been treated. The aim of this article is to present absolute stability conditions for chaotic and asynchronous schemes. We give criteria so as to predict instability thresholds, and study in details the significant example of a damped spring-mass system. Our results, which make use of random matrices products theory, stress the point that the desynchronization of phenomena, and a random scheduling of their activations, can lead to instability.

Keywords: Chaotic asynchronous scheduling, Multi-interaction systems, Ordinary differential systems, Absolute stability, Random matrices products.

1 Introduction

Multi Interaction Systems (MIS) (Desmeulles, Bonneaud, Redou, Rodin, and Tisseau, 2009) were introduced in the context of Virtual Reality. The purpose was precisely to provide medical researchers with a simulator dedicated to virtual experimentation and satisfying four essential points :

[1] UEB, Enib, LISYC- EA 3883, CERV.
[2] UEB, LISYC- EA 3883, CervVal.
[3] UEB, UBO, LISYC- EA 3883.

- the ability to interact with the simulated system in the course of the simulation, without stopping it, by adding or removing interactions or constituents in the system, so as to be as close as possible from the *in vitro* experimentation.
- the possibility for a computer-illiterates to achieve this interactive simulation of biochemical reactions without knowing anything about programmation or numerical methods for solving differential systems.
- the capacity to take into account widely different time and space scales for simulated phenomena.
- the property of the simulator to give as precise results as possible, particularly when solving differential or partial derivative systems.

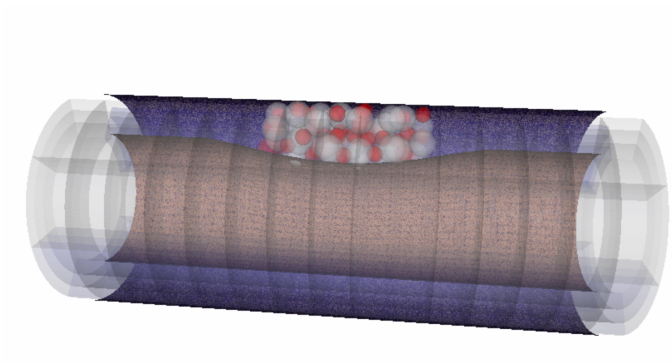

Figure 1: In Virtuo simulation of atherosclerosis using a MIS.

Such a challenge implies to consider different levels of description (Fitzgerald, Goldbeck-Wood, Kung, Petersen, Subramanian, Wescott, and Source, 2008) : *quantum scale, molecular scale, mesoscale...*, and interfacing these different levels is an additional and not negligible difficulty (Chirputkar, Qian, and Source, 2008). Thereby, the challenge was to create a virtual reality simulator for what was called *in virtuo* experimentation (Tisseau, 2001), that is, to summarize, *in silico* computations in the conditions of *in vitro* experiments. For this purpose, the MIS paradigm proposed to reify interactions into the system instead of constituents, with the main and basical advantage to provide modularity, i.e. adding or removing interactions in course of simulation. Thus, a MIS can be seen as a collection of autonomous processes-interactions-, each acting on a collection of variables-constituents, and carrying its own time step. This radical change in perspective has made feasible the constraints, outlined above, of *in virtuo* experimentation. It also led to the choice of a new kind of simulation algorithms, based upon random scheduling of interactions inside the system: chaotic asynchronous scheduling (Desmeulles, Bonneaud, Redou, Rodin, and Tisseau, 2009; Redou, Kerdélo, Le Gal, Querrec, Rodin, Abgrall, and Tisseau, 2005; Redou, Desmeulles, Abgrall, Rodin, and Tisseau, 2007). The principle is to consider each (physical, biological) phenomenon acting on the system -i.e. an interaction between constituents- as autonomous. The simulation engine evolves interactions asynchronously (one after the other, into cycles) and chaotically (the order of interventions changes randomly from one cycle to the other). This scheduling was chosen in order to avoid the typical inflexibility of synchronous systems, as well as bias in numerical results.

From a formal point of view, chaotic asynchronous scheduling provides methods of resolution of ordinary differential equations or systems (say, to simplify, ODE) (Redou, Kerdélo, Le Gal, Querrec, Rodin, Abgrall, and Tisseau, 2005), as well as methods for partial derivative equations (Redou, Desmeulles, Abgrall, Rodin, and Tisseau, 2007). The present work deals with the case of numerical resolution of ODE. Let us give the principle of chaotic asynchronous scheduling in this context: if one wants to solve the cauchy problem

$$Y'(t) = (f_1 + \cdots + f_p)(t, Y(t)), \quad Y(t_0) = Y_0 \tag{1}$$

the principle is to consider functions f_i as autonomous agents, what is necessary when desynchronizing the different phenomena represented by each of these functions. Considering a numerical method for solving (1), the matching chaotic asynchronous method will be given by successive applications of the chosen method, one for each function. These resolutions take place during the same time step, and the order of resolutions, that is, the order of interventions of functions/phenomena f_i, changes randomly at each time step. Details about this process are given in section 3.

This desynchronization eases a modular and incremental building of the numerical model. This is especially usefull when building biochemical models, since the modeller usually selects, subjectively, the reactions which are most likeky involved, and runs the model. If results are not correct enough, the model is incremented with other reactions, etc., until a satisfying model is obtained. Modularity makes this process natural and doesn't require to stop the simulation to modify the code of equations.

Furthermore, chaotic asynchronous simulation provides a means to bear with non-determinism, which occurs most of the time in chaotic systems because of causality between phenomena at the beginning of the experiment, at a very small scale

(Devaney, 2003). Introducing random causality inside a computation time step facilitates the construction of simulators able to report a non-determinist behavior.

Many applications have been achieved in different domains, though, as said above, biochemical kinetics is a natural application context for chaotic and asynchronous scheduling: a classical example is given by cancer, since chromosomic instability (Hanahan and Weinberg, 2000) implies on a regular basis modifications or creations of new reactions (Bos, 1989). In this context, an application of this scheduling to computer simulation of multiple myeloma was recently achieved (Rodin, Querrec, Ballet, Bataille, Desmeulles, Abgrall, and Tisseau, 2009). Notice that it is also used for simulation of MAPK pathway (Querrec, Rodin, Abgrall, Kerdelo, and Tisseau, 2003), and simulation of the extrinsic pathway of blood coagulation (Lu, Broze, and Krishnaswamy, 2004). In an other context, chaotic asynchronous scheduling is used for simulation of sea states, which is typically multi-model and multi-scale (Le Gal, Parenthoen, Béal, and Tisseau, 2007).

Proofs of convergence for these methods have been established (Redou, Kerdélo, Le Gal, Querrec, Rodin, Abgrall, and Tisseau, 2005), (Redou, Desmeulles, Abgrall, Rodin, and Tisseau, 2007), but the problem of absolute stability (Ascher and Petzold, 1998) has not yet been treated, despite its importance: indeed, the region of absolute stability can be seen as the set of values of the time step outside which the distance between the exact solution and the approximate gets out of control. Thus, when simulating multiscale problems, one has to find a compromise between precision and a realistic time simulation, and this choice can not be made without knowing the region of absolute stability. Another important case where this knowledge is crucial is given by stiff problems (Hairer, Norsett, and Wanner, 1996), with brutal variations of the solution of an ODE. The aim of this article is to present absolute stability conditions for chaotic and asynchronous schemes. We give general results, based upon the theory of products of random matrices, and stress the point that in certain circumstances, these schemes may impose strong conditions on the time step, mainly when opposing forces are at work in the system. A significative illustration is the case of a damped spring-mass system where the different physical phenomena are desynchronized.

In section 2, we remind the reader of the problem of absolute stability of methods for solving ODEs, so as properties of classical explicit and implicit schemes. In section 3, we describe how desynchronization of phenomena leads to define asynchronous and chaotic asynchronous schemes. We also recall results of convergence for these methods. Sections 4 and 5 expose the main results of this paper : we study absolute stability for asynchronous and chaotic asychronous schemes, in a general context, providing conditions on integration steps. Finally, section 6 exposes the practical example of a damped spring-mass system, where the three phenomena involved are not considered as synchronous. In this case, we show that absolute stability conditions can be drastical.

2 Absolute stability issues

In this section, we simply remind the reader of absolute stability issues Hairer, Norsett, and Wanner (1983, 1996), so as classical cases.

2.1 Definitions

We consider the following differential system (classically named *test equation*)

$$X'(t) = A \cdot X(t) \tag{2}$$

where A is a square matrix with *distinct eigenvalues all lying in the negative half-plane* $\mathrm{Re}(z) < 0$. Its general solution is

$$X(t) = \exp(tA) \cdot X(0)$$

One has, under these conditions,

$$\lim_{t \to \infty} X(t) = \vec{0}$$

that is, the solution tends towards its stationary point.

Consider the one dimensional case $y' = \lambda y$, $\mathrm{Re}(\lambda) < 0$, and assume that, with the given method, y_n approximates the exact solution $y(t_n)$ at time t_n. The *region of absolute stability for a method* is the set of values of the time step h and of λ for which

$$\lim_{n \to \infty} y_n = \vec{0}$$

is verified. This definition is illustrated in the following section, but one can consider absolute stability as the capability of a method to bare brutal variations of the solution, even with large time steps. This is of great interest when dealing with real-time multiscale simulations, what induces the choice of optimal time steps.

In the multidimensional case given by equation (2), it is easy to see that a necessary condition for the absolute stability of a method is that $h\lambda$ be in the stability region of this method for each eigenvalue λ of A and h the largest time step.

2.2 *Examples of classical Euler methods*

Let us recall classical results about simple methods, before exposing what regards asynchronous schemes. The most simple method for the resolution of (2) is the Euler algorithm. It is given by

$$X_n = X_{n-1} + hA \cdot X_{n-1} \tag{3}$$

Considering the one dimensional case, one easily gets the absolute stability region : this is the open disk defined by $\{z = hA \in \mathbb{C} \ : \ |1+z| < 1\}$.
Let us consider the simple example $A = -\lambda$, $\lambda \in \mathbb{R}_+$. Equation (2) is simply $X'(t) = -\lambda X(t)$, and its solution is $X(t) = e^{-\lambda t} X(0)$. Applying explicit Euler scheme, one obtains $X_n = (1-\lambda h)^n X_0$, and the absolute stability condition is $|1-\lambda h| < 1$, that is, $h < 2/\lambda$. Figure 2 shows different approximations of the solution with $\lambda = 6$, i.e. $2/\lambda = 1/3$.

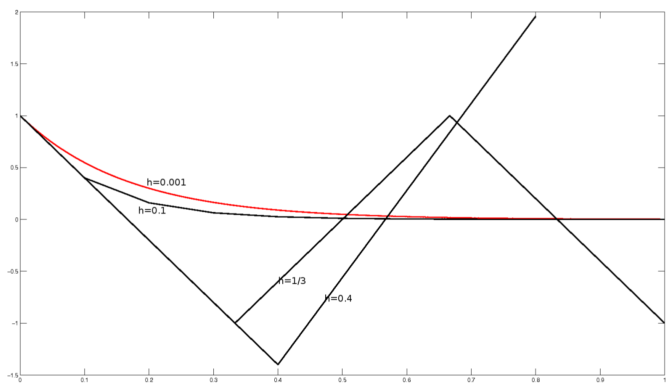

Figure 2: An application of Euler scheme for the test-equation $y' = -6y$ with different values of the time step. The absolute stability threshold $h = 1/3$ is highlighted.

In the general multidimensional case, equation (3) gives

$$X_n = (I + hA)^n \cdot X_0 \tag{4}$$

where I is the identity matrix. Therefore (see section 4), the absolute stability condition is here

$$\rho(I + hA) < 1$$

with $\rho(M)$ the spectral radius of M.

A more efficient algorithm, regarding absolute stability, is given by the Implicit Euler method

$$X_n = X_{n-1} + hA \cdot X_n \tag{5}$$

Here, one easily gets the fact that the absolute stability region is the whole complex plane. Indeed, for the one dimensional test-equation
$X'(t) = -\lambda X(t)$, one gets with implicit Euler method $X_n = \frac{1}{(1+\lambda h)^n} X_0$, so that $\lim_{n\to\infty} X_n = 0\ \forall h$, and absolute stability is guaranteed. Figure 3 shows different approximations, for the same example and the same values of the time step as in figure 2.

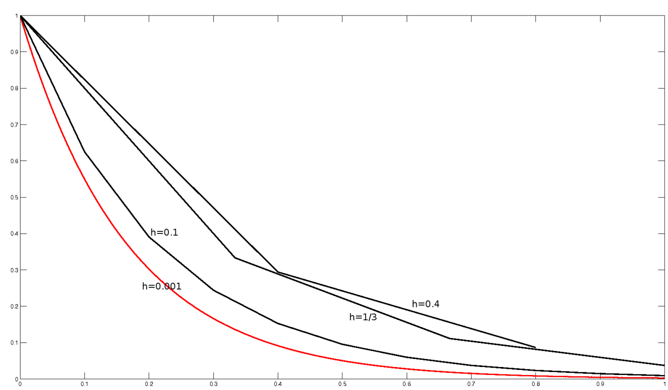

Figure 3: An application of Implicit Euler scheme for the test-equation $y' = -6y$ with the same values of the time step as in figure 2.

In the general multidimensional case, equation (3) gives

$$X_n = (I - hA)^{-n} \cdot X_0$$

and the absolute stability condition is here

$$\rho((I - hA)^{-1}) < 1$$

3 Asynchronous and chaotic asynchronous schemes

Chaotic asynchronous schemes were presented in Redou, Kerdélo, Le Gal, Querrec, Rodin, Abgrall, and Tisseau (2005) and Redou, Desmeulles, Abgrall, Rodin, and Tisseau (2007), where their general definition and convergence properties were detailed. For the sake of simplicity, and because we deal with absolute stability, we will simply remind the reader of the principle of asynchronous and chaotic asynchronous scheduling, when applied to test-equation (2). The example of explicit Euler scheme, though simple, will enable us to stress the difference between asynchronous and chaotic asynchronous schemes, so as problems posed by “poor" properties of the spectral radius.

Here is the principle: we consider equation (2) and assume that matrix A is written

$$A = \sum_{i=1}^{m} A_i \tag{6}$$

As regards applications in the domain of interactive real-time simulations, each A_i is the matricial representation of a distinct phenomenon. Each of these phenomena will be activated at specific moments. In the asynchronous case one defines a scheduling that will be repeated all along the simulation. In the chaotic asynchronous case, this order of phenomena activations changes randomly at each cycle.

The next sections describe in details these simulation methods.

3.1 Asynchronous Euler schemes

Consider a fixed permutation $\sigma \in S_m$, where S_m is the symmetric group of permutations of m elements. This permutation is used at each time step, and characterizes the scheduling of A_i's interventions in cycles. We recall that this “desynchronization” mainly makes it easy to add or remove phenomena in the course of a running simulation, without stopping it.

The principle is to execute the same algorithm (here explicit Euler) successively with each phenomenon involved, according to the order of interventions fixed by the permutation σ. On one time step, the execution of asynchronous explicit Euler algorithm gives :

$$\begin{aligned} X^{*1} &= X_{n-1} + hA_{\sigma(1)} \cdot X_{n-1} \\ X^{*2} &= X^{*1} + hA_{\sigma(2)} \cdot X^{*1} \\ &\vdots \\ X_n &= X^{*(m-1)} + hA_{\sigma(m)} \cdot X^{*(m-1)} \end{aligned}$$

Thus, one gets

$$X_n = \prod_{i=1}^{m} (I + hA_{\sigma(i)}) \cdot X_{n-1} \iff X_n = \left(\prod_{i=1}^{m} (I + hA_{\sigma(i)}) \right)^n \cdot X_0 \tag{7}$$

We stress again the point that the same permutation σ is used here on each time step.
In a similar way, asynchronous scheme applied to implicit Euler algorithm leads to :

$$X_n = \left(\prod_{i=1}^{m} \left(I - hA_{\sigma(m-i+1)} \right)^{-1} \right)^n \cdot X_0 \tag{8}$$

3.2 Chaotic asynchronous explicit Euler scheme

The fundamental difference between asynchronous and chaotic asynchronous schemes is that a new permutation is chosen at each time step for the scheduling of phenomena. During time step n, the order of interventions of phenomena involved makes matrices intervene the following way : $A_{\sigma_n(1)}, A_{\sigma_n(2)}, \ldots, A_{\sigma_n(m)}$, where σ_n is the permutation of m operators A_i which is involved at time n.
For this time step, chaotic asynchronous Euler algorithm gives :

$$\begin{aligned} X^{*1} &= X_{n-1} + hA_{\sigma_n(1)} \cdot X_{n-1} \\ X^{*2} &= X^{*1} + hA_{\sigma_n(2)} \cdot X^{*1} \\ &\vdots \\ X_n &= X^{*(m-1)} + hA_{\sigma_n(m)} \cdot X^{*(m-1)} \end{aligned}$$

Thus, one gets

$$X_n = \prod_{i=1}^{m} (I + hA_{\sigma_n(i)}) \cdot X_{n-1}$$

that is

$$X_n = \prod_{k=1}^{n} \prod_{i=1}^{m} (I + hA_{\sigma_k(i)}) \cdot X_0$$

Here again, this chaotic asynchronous scheme may be applied to implicit Euler algorithm and leads to

$$X_n = \prod_{k=1}^{n} \prod_{i=1}^{m} (I - hA_{\sigma_k(m-i+1)})^{-1} \cdot X_0$$

As an introduction to the kind of problems that arise when using these methods, the next part deals exclusively with the asynchronous case. The chaotic case will be even more difficult to handle, because it involves stochastic processes.

4 Issues and results about asynchronous Euler schemes

In the following, we denote by $\rho(M)$ the spectral radius of a matrix M. We will make use of the following fundamental property:

Theorem 4.1. *Quarteroni, Sacco, and Saleri (2000) Let M a matrix in* $\mathbb{C}^{n\times n}$.

$\lim_{n\to\infty} M^n = 0 \iff \rho(M) < 1$

4.1 Stability regions

Considering equation (4), theorem 4.1 implies that the absolute stability condition for explicit Euler scheme is given by

$$\rho(I + hA) < 1 \tag{9}$$

The same way, considering equation (7), the absolute stability condition for asynchronous explicit Euler scheme is given by

$$\rho\left(\prod_{i=1}^{m}(I + hA_{\sigma(i)})\right) < 1$$

with σ the fixed permutation chosen at the beginning of the execution. An obvious remark is that this condition is not as easy to check as (9), and may induce complex computations (our damped mass-spring example will exhibit this complexity). This is the reason why it is important to provide absolute stability conditions for these asynchronous schemes. This is what we present in the following.

Moreover, since any permutation may be initially chosen and then used during the whole simulation, we get the trivial following criteria for explicit and implicit asynchronous Euler schemes:

Proposition 4.2. *1. The absolute stability domain for asynchronous explicit Euler scheme, when resolving* $X' = A\cdot X = (\sum_{i=1}^{m} A_i)\cdot X$, *is given by the set*

$$S_A = \left\{ h \in \mathbb{R}_+ :\ \forall \sigma \in S_m,\ \rho\left(\prod_{i=1}^{m}(I + hA_{\sigma(i)})\right) < 1 \right\} \tag{10}$$

2. The absolute stability domain for asynchronous implicit Euler scheme, when resolving
$X' = A\cdot X = (\sum_{i=1}^{m} A_i)\cdot X$, *is given by the set*

$$S_A = \left\{ h \in \mathbb{R}_+ :\ \forall \sigma \in S_m,\ \rho\left(\prod_{i=1}^{m}(I - hA_{\sigma(m-i+1)})^{-1}\right) < 1 \right\} \tag{11}$$

In section 6, a detailed example will show that these criteria may induce complex conditions on time steps, when applied to concrete cases. But even the most simple case of a one dimensional equation leads to non trivial conditions, the following example may be instructive.

4.2 Examples of stability regions in one dimension

In this section we illustrate the non triviality of absolute stability conditions for asynchronous schemes, even in elementary cases. We want to show that the conditions of absolute stability for asynchronous schemes, in both cases of explicit and implicit Euler, are uneasy to handle in general. Even in the simple case of one differential equation, where all A_i are real numbers and commute, conditions (10) and (11) are $|\prod_{i=1}^m (1+hA_i)| < 1$ and $|\prod_{i=1}^m \frac{1}{1-hA_i}| < 1$, so that a general condition on h is not easy to extract. For instance, one can consider the special case where $m=2$ and A_1, A_2 are real numbers, here denoted $-\lambda_1$ and $-\lambda_2$: we assume in the following $\lambda_1+\lambda_2>0$, so that the problem

$$x'(t) = -(\lambda_1+\lambda_2)x(t), \quad \lambda_1+\lambda_2>0 \tag{12}$$

remains stiff.

In the case of the explicit Euler scheme, the absolute stability condition for (12) is $|(1-\lambda_1 h)(1-\lambda_2 h)| < 1$. A direct study leads to the following alternative:

Proposition 4.3.
- *If $\lambda_1\lambda_2 > 0$, the absolute stability condition for (12) is*

$$h \in \left]0; \frac{\lambda_1+\lambda_2-\sqrt{(\lambda_1+\lambda_2)^2-8\lambda_1\lambda_2}}{2\lambda_1\lambda_2}\right[\cup \left]\frac{\lambda_1+\lambda_2+\sqrt{(\lambda_1+\lambda_2)^2-8\lambda_1\lambda_2}}{2\lambda_1\lambda_2}; \frac{\lambda_1+\lambda_2}{\lambda_1\lambda_2}\right[$$

- *If $\lambda_1\lambda_2 < 0$, the absolute stability condition for (12) is*

$$h \in \left]0; \frac{\lambda_1+\lambda_2-\sqrt{(\lambda_1+\lambda_2)^2-8\lambda_1\lambda_2}}{2\lambda_1\lambda_2}\right[$$

On the other hand, the absolute stability condition for the implicit scheme (5) is $\left|\frac{1}{(1+h\lambda_1)(1+h\lambda_2)}\right| < 1$ what leads to another alternative:

Proposition 4.4.
- *If $\lambda_1\lambda_2 > 0$, the absolute stability condition for (12) is trivial, so that the method is absolutely stable.*

- *If $\lambda_1\lambda_2 < 0$, the absolute stability condition for (12) is*

$$h \in \left]0; \frac{\lambda_1+\lambda_2}{-\lambda_1\lambda_2}\right[\cup \left]\frac{\lambda_1+\lambda_2+\sqrt{(\lambda_1+\lambda_2)^2-8\lambda_1\lambda_2}}{-2\lambda_1\lambda_2}; +\infty\right[$$

For example, let us consider the case $\lambda_1 = -3,\ \lambda_2 = 11$, so that $\lambda = \lambda_1+\lambda_2 = 8$. Therefore our problem is the stiff one :

$$x' = -8x = -(-3+11)x$$

Absolute stability conditions are in this case :

- classical synchronous Euler : $h < \frac{2}{8} = 0.25$
- explicit asynchronous Euler: $h < 0.1531$
- implicit asynchronous Euler :

$$h \in]0; 0.2424[\cup]0.3956; +\infty[\tag{13}$$

One can check these results with different simulations.

This simple example suggests that the exact absolute stability region of a general asynchronous scheme may be really complex. Nevertheless, we can prove an easier-to-apply (but less precise) criterion for the explicit case.

4.3 Criterion of absolute stability (explicit scheme)

Proposition 4.5. *Consider the decomposition $A = \sum_{i=1}^m A_i$ where $A \in \mathscr{L}(\mathbb{C}^n)$. Let P be the passage matrix into a base where A is triangular, and the norm defined by*

$$\|v\|_A = \|P^{-1}v\|_1$$

Let $M = \max_i \|A_i\|_A$, and $P_m(X)$ the polynomial defined by
$P_m(X) = (X+1)^m - 1 - mX$. Then, the absolute stability region for Euler chaotic asynchronous scheme, with the desynchronization considered, contains the set

$$S_A = \{h > 0 :\ 0 < \rho(I_n + hA) < 1 - P_m(hM)\}.$$

Proof. If we denote $M_{\|.\|} = \max_i \|A_i\|$ for a norm $\|.\|$, we get

$$\left\|\prod_{i=1}^m (I_n + hA_\sigma(i))\right\| \le \|I_n + hA\| + P_m(hM_{\|.\|})$$

Considering the norm $\|v\|_A = \|P^{-1}v\|_1$ it is easy to see that $\rho(I_n + hA) = \|I_n + hA\|_A$. Therefore, one has

$$\left\| \prod_{i=1}^{m} (I_n + hA_\sigma(i)) \right\| \leq \rho(I_n + hA) + P_m(hM_A)$$

Finally, it suffices to extract the condition on hM from the relation

$$0 < \rho(I_n + hA) < 1 - P_m(hM_A)$$

□

We will see an illustration of this criterion in section 6. But for now, in the next section, we study the chaotic asynchronous case.

5 Chaotic asynchronous Euler schemes

This section presents our main results about absolute stability of chaotic asynchronous schemes. We recall that the fundamental difference between asynchronous and chaotic asynchronous schemes is the fact that, in the latter case, a new permutation is chosen at each time step for the scheduling of phenomena. This leads to radically different properties of stability, as detailed below.

Let us first recall that the execution of chaotic asynchronous Euler schemes, when solving equation (2) with the decomposition (6), leads to the following formulas, where σ_k is the permutation used at step k, $k \leq n$:

- For chaotic asynchronous explicit Euler scheme,

$$X_n = \prod_{k=1}^{n} \prod_{i=1}^{m} (I + hA_{\sigma_k(i)}) \cdot X_0$$

- For chaotic asynchronous implicit Euler scheme,

$$X_n = \prod_{k=1}^{n} \prod_{i=1}^{m} (I - hA_{\sigma_k(m-i+1)})^{-1} \cdot X_0$$

In this section, we prove a general criterion which ensures that the upper Lyapunov exponent associated with a distribution on $GL(d, \mathbb{R})$ is negative. Then we apply this criterion to the absolute stability of chaotic asynchronous methods.

5.1 Negative upper Lyapunov exponent

Let us start with some common notations and definitions (see Bougerol and Lacroix (1985) for a detailed theory about the products of random matrices).

Definition 5.1. *If $(B_i)_{i\geq 1}$ is a sequence of i.i.d random matrices, we write β_n the product $B_n.\cdots.B_1$. If $\ln^+\|B_1\|$ is integrable, then the following limit exists and is called upper Lyapunov exponent of the sequence (or equivalentely of the distribution associated to the sequence):*

$$\lim_n \frac{1}{n}\mathbb{E}\left[\ln\|\beta_n\|\right] = \gamma$$

Definition 5.2. *If μ is a probability measure on $GL(d,\mathbb{R})$, G_μ is the smallest closed subgroup of $GL(d,\mathbb{R})$ that contains the support of μ.*

Definition 5.3. *A subset S of $GL(d,\mathbb{R})$ is said to be irreducible if there is no proper subspace $V\subset\mathbb{R}^d$ such that $M(V)=V$ for all $M\in S$.*

We will use the following lemma:

Lemma 5.4. *Let $\{B_n, n\geq 1\}$ be a sequence of independent random matrices of $GL(d,\mathbb{R})$ with common distribution μ, and $\beta_n = B_n.\cdots.B_1$. We suppose that:*

1. *G_μ is irreducible.*
2. *$\ln^+\|B_1\| + \ln^+\left\|B_1^{-1}\right\|$ is integrable*

Then

$$\lim_n \frac{1}{n}\sup_{\|x\|=1}\mathbb{E}\left[\ln\|\beta_n\cdot x\|\right] = \gamma$$

Proof. First of all, let us check that the sequence

$$a_n = \sup_{\|x\|=1}\mathbb{E}\left[\ln\|\beta_n\cdot x\|\right]$$

is subadditive. For any integer n and m one has

$$\begin{aligned}\mathbb{E}\left[\ln\|\beta_{n+m}.x\|\right] &= \mathbb{E}\left[\ln\|B_{n+m}\cdots B_{n+1}B_n\cdots B_1\cdot x\|\right]\\ &= \mathbb{E}\left[\ln\left\|B_{n+m}\cdots B_{n+1}\frac{B_n\cdots B.x}{\|B_n\cdots B_1\cdot x\|}\right\|\right] + \mathbb{E}\left[\ln\|B_n\ldots B_1\cdot x\|\right]\end{aligned}$$

As $\frac{B_n \cdots B_1 \cdot x}{\|B_n \cdots B_1 \cdot x\|}$ is unitary, considering the upper bounds on $\|x\| = 1$ leads to

$$a_{n+m} \leq a_n + a_m$$

Thereby, the sequence $\frac{a_n}{n}$ converges: we denote γ' its limit.
Since G_μ is irreducible and $\mathbb{E}\left[\ln^+ \|B_1\|\right]$ is finite, we know that, for any $x \neq 0$ (see Bougerol and Lacroix (1985) p.72):

$$\lim_n \frac{1}{n} \ln \|\beta_n \cdot x\| = \gamma \quad \text{almost surely.}$$

Now, an easy computation shows that

$$\frac{1}{n} |\ln \|\beta_n \cdot x\|| \leq \frac{1}{n} \sum_{i=1}^{n} \left(\ln^+ \|B_i\| + \ln^+ \left\|B_i^{-1}\right\|\right)$$

From the law of large numbers, the right hand side converges in L^1, thereby, it is uniformly integrable. Thus, the left hand side is also uniformly integrable, and as it converges almost surely to γ, it converges as well in L^1:

$$\lim_n \frac{1}{n} \mathbb{E}\left[\ln \|\beta_n . x\|\right] = \gamma \tag{14}$$

And since $\frac{1}{n} \mathbb{E}\left[\ln \|\beta_n \cdot x\|\right] \leq \frac{a_n}{n}$, one has

$$\gamma \leq \gamma'$$

On the other hand, one has

$$\sup_{\|x\|=1} \mathbb{E}\left[\ln \|\beta_n \cdot x\|\right] \leq \mathbb{E}\left[\sup_{\|x\|=1} \ln \|\beta_n \cdot x\|\right] = \mathbb{E}\left[\ln \|\beta_n\|\right]$$

The right hand side, by definition, converges to γ, so that $\gamma' \leq \gamma$, which ends the proof. □

From this lemma, one can deduce the following result based on the negativity of the upper Lyapunov exponent:

Proposition 5.5. *Let $\{B_n, n \geq 1\}$ be a sequence of independent random matrices of $GL(d,\mathbb{R})$ with common distribution μ that satisfies*

1. *G_μ is irreducible.*

2. $\ln^+ \|B_1\| + \ln^+ \left\|B_1^{-1}\right\|$ *is integrable*

If there exists an integer m such that

$$\sup_{\|x\|=1} \mathbb{E}\left[\ln \|\beta_m \cdot x\|\right] < 0$$

Then, for any x

$$\lim_n \beta_n \cdot x = 0 \quad \textit{almost surely.}$$

Proof. From lemma 5.4, we know that

$$\lim_n \frac{1}{n} \sup_{\|x\|=1} \mathbb{E}\left[\ln \|\beta_n \cdot x\|\right] = \gamma$$

But, as $a_n = \sup_{\|x\|=1} \mathbb{E}\left[\ln \|\beta_n \cdot x\|\right]$ is a subadditive sequence, we know that

$$\lim_n \frac{a_n}{n} = \inf_m \frac{a_m}{m}$$

Our hypothesis ensures that $\inf_m \frac{a_m}{m} < 0$, so that $\gamma < 0$. But, as in lemma 5.4, we know that for any $x \neq 0$

$$\lim_n \frac{1}{n} \ln \|\beta_n \cdot x\| = \gamma \quad \text{almost surely.}$$

This suffices to deduce the result. □

5.2 Absolute stability of chaotic asynchronous schemes

In this section we will simply apply proposition 5.5 to chaotic asynchronous schemes. In this particular context, assumptions of this proposition are generally satisfied, so that the following criterion is relevant.

Definition 5.6. *In the following proposition, a matrix is said to be associated with a chaotic asynchronous method if it is a random product of matrices intervening at each time step: for instance, matrices associated with chaotic asynchronous explicit euler scheme for the resolution of* $X' = (\sum_{i=1}^m A_i) \cdot X$ *will be the following products:*

$$B_k = \prod_{i=1}^{m} (I + hA_{\sigma_k(i)}), \; \sigma_k \in S_m$$

Of course, our problem regards the limit of products of such associated matrices.

Proposition 5.7. *Let $B = \{B_1, \ldots, B_N\} \subset GL(d, \mathbb{R})$ be the matrices associated with a chaotic asynchronous method applied to a linear equation. We suppose that B is irreducible, then, if there exists an integer m such that*

$$\sup_{\|x\|=1} \prod_{1\leq, i_1, \cdots, i_m \leq N} \|B_{i_1} \cdots B_{i_m} \cdot x\| < 1$$

Then the method is almost surely absolutely stable.

Proof. First, it is easy to check that if B is irreducible, then G_μ is also irreducible (where μ is the uniform distribution on B).

Since the matrices are equidistributed, one has:

$$\begin{aligned}
\sup_{\|x\|=1} \mathbb{E}[\ln \|\beta_m \cdot x\|] &= \sup_{\|x\|=1} \frac{1}{N^m} \sum_{1\leq, i_1, \ldots, i_m \leq N} \ln \|B_{i_1} \cdots B_{i_m} \cdot x\| \\
&= \sup_{\|x\|=1} \frac{1}{N^m} \ln \left(\prod_{1\leq, i_1, \ldots, i_m \leq N} \|B_{i_1} \cdots B_{i_m} \cdot x\| \right) \\
&= \frac{1}{N^m} \ln \left(\sup_{\|x\|=1} \prod_{1\leq, i_1, \ldots, i_m \leq N} \|B_{i_1} \cdots B_{i_m} \cdot x\| \right)
\end{aligned}$$

Our hypothesis insures that

$$\sup_{\|x\|=1} \mathbb{E}[\ln \|\beta_m \cdot x\|] < 0$$

Since we have only a finite number of matrices, the condition of integrability of $\ln^+ \|B_1\| + \ln^+ \left\|B_1^{-1}\right\|$ is satisfied. Thus we may apply proposition 5.5 and conclude. □

With quite simple calculus this criterion can indicate, depending on the value of h, that a chaotic asynchronous method is stable. Nevertheless in some cases, the criteria is not applicable because the sequence of functions $\prod_{1\leq, i_1, \cdots, i_m \leq N} \|B_{i_1} \cdots B_{i_m} \cdot x\|$ converges only almost everywhere. We could have improved the criterion to handle this fact, and produce a result like the following one :

Proposition 5.8. *Let $S^d = \{x \in \mathbb{R}^d, \|x\| = 1\}$. If there exists an integer m and a subset $N \subset S^d$ of null measure such that*

$$\sup_{x \in S-N} \prod_{1\leq, i_1, \cdots, i_m \leq N} \|B_{i_1} \cdots B_{i_m} \cdot x\| < 1$$

Then the method is almost surely absolutely stable.

But such a proposition would be useless in practice. Even in the case of a negative Liapunov exponent, the quantity $\prod_{1\leq,i_1,\cdots,i_m\leq N}\|B_{i_1}\cdots B_{i_m}\cdot x\|$ may grow to infinity on a set of null measure (this is precisely the case of system S_1 in section 6.5). In these cases, the estimation of Lyapunov exponent may become the only way to compute stability conditions.

The next section is devoted to examples and illustrations of all the previous results and observations.

6 Applications and illustrations

The damped mass-spring system is a particular case of desynchronisation of one single differential equation, this is why we first describe this general case.

6.1 Desynchronization of one single differential equation

The case of one single linear differential equation with order m is given by:

$$x^{(m)}(t)-\sum_{i=0}^{m-1}a_ix^{(i)}(t)=0 \tag{15}$$

This equation can be written as a linear differential system: with the notations $z_i=x^{(i)}$, $0\leq i\leq m-1$, one gets the system

$$\begin{pmatrix} z_0\\ z_1\\ \vdots\\ z_{m-1}\end{pmatrix}'=\begin{pmatrix} 0 & 1 & 0 & \dots\\ 0 & 0 & \ddots & 0\\ 0 & \dots & 0 & 1\\ a_0 & a_1 & \dots & a_{m-1}\end{pmatrix}\cdot\begin{pmatrix} z_0\\ z_1\\ \vdots\\ z_{m-1}\end{pmatrix}$$

what can be denoted, with obvious notations,

$$Z'=A\cdot Z.$$

Consider the elementary matrices E_{ij}, $1\leq i,j\leq m$. Assume that coefficients a_i each characterize a distinct phenomenon: we can associate to a_i the matrix

$$P_i=a_iE_{m,i+1}$$

and introduce an “integration phenomenon" given by the matrix

$$\mathrm{Int}=\sum_{i=1}^{m-1}E_{i,i+1}$$

With these notations, one easyly gets

$$A = \text{Int} + \sum_{i=0}^{m-1} P_i$$

Therefore, one can apply an asynchronous scheme (chaotic or not), where the P_is and the integration phenomenon are desynchronized. Our main example of a damped mass-spring system will illustrate this process.

6.2 Damped mass-spring system

In sections 4 and 5, we have exposed absolute stability conditions for asynchronous and chaotic asynchronous schemes. In the following, we propose an illustration of these results in the case of a second order linear differential equation, with drastic absolute stability conditions when physical phenomena involved are desynchronized. We volunteerly consider a typical case of antagonist phenomena leading to a more significative unstability when they are desynchronized. Indeed, we consider the case of a damped spring-mass system, that can be represented by the following equation

$$x'' = -g - \frac{k}{m}x - \frac{\gamma}{m}x' \tag{16}$$

where :

- g is the gravity field
- m is the mass of the object
- k is the elasticity constant of the spring
- γ is the damp coefficient

All along this section, we will carry simple computations in order to illustrate our problems. We will consider two cases of such systems, defined by the following parameters :

$$(\gamma, k, m) = (1, 4, 1) \text{ referred as } (S_1) \tag{17}$$

and

$$(\gamma, k, m) = (8, 1, 1) \text{ referred as } (S_2) \tag{18}$$

But for now, we will try to explore our system in the general case. According to our theoretical study, we will first deal with the asynchronous case, before dealing

with the chaotic asynchronous one. This example will clearly expose how chaotic schemes, though they are a bit more precise than non chaotic ones, may suffer from great instability.
Using the notations

$$x_1 = x, \quad x_2 = x'$$

equation (16) can be written as the system

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix}' = \begin{bmatrix} 0 & 1 \\ -k/m & -\gamma/m \end{bmatrix} \cdot \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ -g \end{bmatrix}$$

The simple change of variables

$$X = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 0 & 1 \\ -k/m & -\gamma/m \end{bmatrix}^{-1} \cdot \begin{bmatrix} 0 \\ -g \end{bmatrix}$$

leads to the equivalent system

$$X' = \begin{bmatrix} 0 & 1 \\ -k/m & -\gamma/m \end{bmatrix} \cdot X$$

In the following, we use the notations:

$$A = \begin{bmatrix} 0 & 1 \\ -k/m & -\gamma/m \end{bmatrix}, A_1 = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

$$A_2 = \begin{bmatrix} 0 & 0 \\ -k/m & 0 \end{bmatrix}, \quad A_3 = \begin{bmatrix} 0 & 0 \\ 0 & -\gamma/m \end{bmatrix}$$

so that $A = A_1 + A_2 + A_3$.
Our study of absolute stability implies that the eigenvalues of A both be in the negative half-plane. Since these eigenvalues are

$$\lambda_{\pm} = \frac{-\gamma \pm \sqrt{\gamma^2 - 4mk}}{2m}$$

a direct computation shows that $Re(\lambda_{\pm}) < 0 \iff (k, \gamma) \in (\mathbb{R}_+^*)^2$
Before exposing results for asynchronous schemes, we first recall classical results as regards equation (16) in the case of classical Euler schemes.

6.3 Classical Euler schemes

First, in the explicit Euler case, the absolute stability condition is here $\rho(I+hA) \in [0;1[$, and is equivalent to $|1+h\lambda_{\pm}| < 1$. Therefore, we get the conditions :

- If $\gamma^2 - 4mk \geq 0$,
$$h < \frac{4m}{\gamma + \sqrt{\gamma^2 - 4mk}}.$$
- If $\gamma^2 - 4mk \leq 0$,
$$h < \frac{\gamma}{k}.$$
In the case of system (S_1) the condition is $h < 0.25$ and for (S_2) one gets $h < 8 - 2\sqrt{15} \sim 0.254$

In the implicit case, the absolute stability condition is $\rho((I-hA)^{-1}) < 1$. Nevertheless, this condition is trivial, since we have seen that implicit Euler method is absolutely stable, with no condition on the time step. This can obviously be verified by considering eigenvalues of $(I-hA)^{-1}$.

6.4 Absolute stability conditions for asynchronous Euler schemes

Now we turn to asynchronous methods and we will show how conditions (10) and (11), though simple, can lead to difficult computations, even on our elementary example.

6.4.1 Asynchronous explicit Euler

We prove the following result :

Proposition 6.1. *Conditions of absolute stability for asynchronous explicit Euler scheme for the damped mass-spring system are the following :*

- *If* $2\gamma^2 - k < 0$,
$$h < -\frac{\gamma}{m} + \sqrt{\frac{\gamma^2}{m^2} + 4\frac{m}{k}} \tag{19}$$
- *If* $2\gamma^2 - k \geq 0$,
$$h < \frac{m}{\gamma} \tag{20}$$

Proof. Absolute stability condition is here

$$\forall \sigma \in S_3,\ \rho\left((I+hA_{\sigma(1)})(I+hA_{\sigma(2)})(I+hA_{\sigma(3)})\right) < 1.$$

Straightforward computations lead to the following results:

- The matrices $(I+hA_1)(I+hA_2)(I+hA_3)$, $(I+hA_2)(I+hA_3)(I+hA_1)$, $(I+hA_3)(I+hA_1)(I+hA_2)$ have the same eigenvalues, that is,

  $$\zeta_\pm = \frac{1}{2m}\left(2m - h\gamma - h^2k \pm \sqrt{-4mh^2k + h^2\gamma^2 + 2h^3\gamma k + h^4k^2}\right)$$

  Denote $P(h) = -4mh^2k + h^2\gamma^2 + 2h^3\gamma k + h^4k^2$.
  If $P(h) < 0$, that is, if $h \in]0; -\frac{\gamma}{k} + 2\sqrt{\frac{m}{k}}[$, one gets $|\zeta_\pm|^2 = 1 - h\frac{\gamma}{m} < 1$.
  If $P(h) > 0$, that is, if $h > -\frac{\gamma}{k} + 2\sqrt{\frac{m}{k}}$, the condition $\zeta_+ < 1$ is always satisfied and the condition $\zeta_- > -1$ implies $h < -\frac{\gamma}{m} + \sqrt{\frac{\gamma^2}{m^2} + 4\frac{m}{k}}$.

- The matrices $(I+hA_1)(I+hA_3)(I+hA_2)$, $(I+hA_2)(I+hA_1)(I+hA_3)$, $(I+hA_3)(I+hA_2)(I+hA_1)$ have the same eigenvalues, that is,

  $$\lambda_\pm = \frac{1}{2m^2}\left(2m^2 - mh\gamma - mh^2k + h^3\gamma k \pm \sqrt{\phi(h)}\right)$$

  with

  $$\phi(h) = -4m^3h^2k + 6m^2h^3\gamma k + m^2h^2\gamma^2 - 2mh^4\gamma^2k + m^2h^4k^2 - 2mh^5k^2\gamma + h^6\gamma^2k^2$$

  We first replace $\frac{\gamma}{m}$ by γ and $\frac{k}{m}$ by k, then denote $X = h\sqrt{\frac{k}{m}}$ and $\alpha = \frac{\gamma}{\sqrt{k}}$
  Therefore, we get

  $$\lambda_\pm = 1 - \frac{\alpha}{2}X - \frac{X^2}{2} + \frac{\alpha}{2}X^3 \pm \frac{X}{2}\sqrt{Q(X)}$$

  with

  $$Q(X) = \alpha^2 - 4 + 6\alpha X + (1-2\alpha^2)X^2 - 2\alpha X^3 + \alpha^2X^4$$

  One has $Q(X) = \alpha^2\prod_{j=1}^4 (X - x_j)$, with

  $$x_1 = \frac{1-2\alpha+\sqrt{1+4\alpha}}{2\alpha} \qquad x_2 = \frac{1-2\alpha-\sqrt{1+4\alpha}}{2\alpha}$$
  $$x_3 = \frac{1+2\alpha+\sqrt{1-4\alpha}}{2\alpha} \qquad x_4 = \frac{1+2\alpha-\sqrt{1-4\alpha}}{2\alpha}$$

Remark that x_2 and x_1 are always real and $(x_3, x_4) \in \mathbb{R} \iff \alpha < 1/4$. In this case, one has $x_2 < 0 < x_4 < x_3 < x_1$.

It suffices now to consider the sign of $Q(X)$ according to the values of X, and compute that :

-If $Q(X) < 0$, one has $|\lambda_\pm|^2 = 1 - \alpha X < 1$.

-If $Q(X) > 0$, the strongest condition is given by $\lambda_+ < 1$, which is satisfied if and only if $X < \frac{1}{\alpha}$. This value is to be taken in account only in the case

$$\frac{1}{\alpha} < -\frac{\gamma}{m} + \sqrt{\frac{\gamma^2}{m^2} + 4\frac{m}{k}} \iff \alpha > \frac{1}{\sqrt{2}} \iff 2\gamma^2 - k > 0$$

□

Once again we will illustrate these results with our two systems (S_1) and (S_2). In the first case, the condition is $h < \sqrt{2-1} \sim 0.4142$, and in the second case $h < 0.125$. Moreover, we computed the criterion given in section 4.5. The following table summarizes all these results :

System	Classic	Asynchronous	Criterion
S_1	0.25	0.414	0.084
S_2	0.254	0.125	0.061

Table 1: Comparison of stability conditions on h for explicit methods. The first column shows the stability conditions for the classic explicit method, the second column shows the exact conditions in the asynchronous case and the third one shows the condition based on the proposition 4.5.

One can notice that in the case of system (S_1), asynchronous explicit Euler scheme gives better results than the classical scheme. Moreover, the criterion given in proposition 4.5 is quite easy to use, but gives strong majorations.

6.4.2 Asynchronous implicit Euler

We prove the following result :

Proposition 6.2. *A sufficient condition for absolute stability of asynchronous implicit euler scheme, in the case of a damped spring-mass system, is*

$$h < \frac{\tau}{\sqrt{k}} \tag{21}$$

with τ the biggest real positive root of the polynomial

$$\Delta(X) = \alpha X^3 + X^2 - 2\alpha X - 4, \quad \alpha = \frac{\gamma}{\sqrt{mk}}$$

Remark : A good approximation of $h < \tau$ is given by

$$h < \frac{m}{2\gamma}\left(-1 + 2\frac{\gamma}{\sqrt{mk}} + \sqrt{1 + 4\frac{\gamma}{\sqrt{mk}}}\right)$$

Proof. A sufficient condition for absolute stability is here

$$\forall \sigma \in S_3,\ \rho((I - hA_{\sigma(1)})^{-1}(I - hA_{\sigma(2)})^{-1}(I - hA_{\sigma(3)})^{-1}) < 1.$$

- Once again, we compute that three of the six matrices $\prod_{i=1}^{3}(I - hA_{\sigma(i)})^{-1}$ have the same eigenvalues, which are

  $$\nu_\pm = 1/2\,\frac{2m + h\gamma - h^2k \pm \sqrt{-4mh^2k + h^2\gamma^2 - 2h^3\gamma k + h^4k^2}}{m + h\gamma}$$

  For these matrices, direct computations give

  $$|\nu_\pm| < 1 \iff h \in [0, \chi[, \quad \chi = \frac{\gamma}{k} + \sqrt{\frac{\gamma^2}{k^2} + 4\frac{m}{k}}$$

- The three other matrices all have the two eigenvalues

  $$\eta_\pm = \frac{2m^2 + mh\gamma - mh^2k - h^3\gamma k \pm \sqrt{\psi(h)}}{2(m + h\gamma)m}$$

  with

  $$\begin{aligned}\psi(h) = & -4m^3h^2k - 6m^2h^3\gamma k + m^2h^2\gamma^2 - 2mh^4\gamma^2k + m^2h^4k^2 \\ & +2mh^5k^2\gamma + h^6\gamma^2k^2\end{aligned}$$

  Once again, we first replace $\frac{\gamma}{m}$ by γ and $\frac{k}{m}$ by k, then denote $X = h\sqrt{\frac{k}{m}}$ and $\alpha = \frac{\gamma}{\sqrt{k}}$. Remark that with these notations one has

  $$h < \chi \iff X < \alpha + \sqrt{4 + \alpha^2}$$

  and

  $$\eta_\pm = \frac{2 + \alpha x - X^2 - \alpha X^3 \pm X\sqrt{R(X)}}{2(1 + \alpha X)}$$

with

$$R(X) = \alpha^2 X^4 + 2\alpha X^3 + (1-2\alpha^2)X^2 - 6\alpha X + \alpha^2 - 4$$

Consider the following real roots of $R(X)$:

$$x_1 = \frac{-1+2\alpha+\sqrt{1+4\alpha}}{2\alpha}$$
$$x_2 = \frac{-1+2\alpha-\sqrt{1+4\alpha}}{2\alpha}$$

Thus :

- If $X \in [x_2, x_1]$ one has $R(X) < 0$ and $|\eta_\pm| < 1$ without condition.
- If $X \in \mathbb{R}_+ \setminus [x_2, x_1]$, the condition $\eta_+ < 1$ is always satisfied, and

  $$\eta_- > -1 \iff X\sqrt{R(X)} < 4 + 3\alpha X - X^2 - \alpha X^3$$

  This leads to two conditions :

  1. $\delta_1(X) < 0$, with

     $$\delta_1(X) = -4 - 3\alpha X + X^2 + \alpha X^3$$

  2. $\delta_2(X) < 0$, with

     $$\delta_2(X) = \alpha(X + \tfrac{1}{\alpha})(\alpha X^3 + X^2 - 2\alpha X - 4)$$

     The condition on δ_1 is weaker than the one on δ_2, so that these conditions lead to $h \in [0, x_2] \cup [x_1, \tau[$, where $\tau < \chi$ and τ is the biggest real positive root of δ_2.

□

There is no need to explore our systems (S_1) and (S_2) according to implicit asynchronous method. Indeed, implicit Euler scheme is absolutely stable, but from the previous result, we know that asynchronous implicit Euler scheme is not stable (for any value of h). This illustrates clearly the loss of performance of this method.

6.5 Absolute stability of chaotic asynchronous Euler schemes

We finally illustrate our theoretical results for the damped spring-mass system with our systems (S_1) and (S_2). First, we illustrate the complex behavior of the Lyapunov exponent, and the fact that it need not be better from explicit scheme to implicit ones. In each of the cases exposed on figures 4 and 5, we compute numerically (using approximations of invariant measures) the Lyapunov exponents in function of the time step.

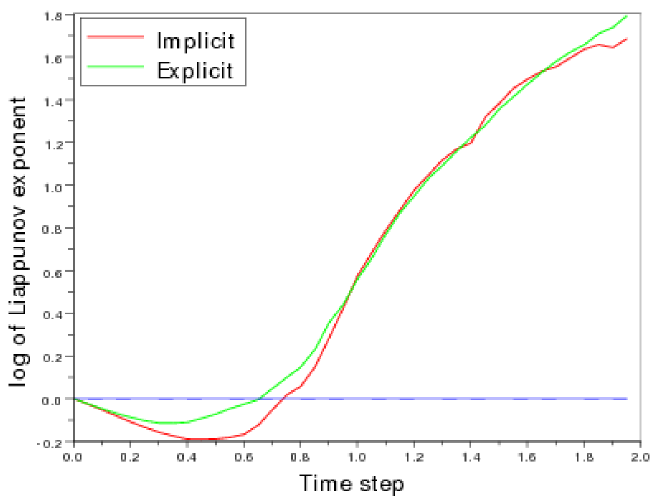

Figure 4: comparison of chaotic asynchronous explicit and implicit Euler, system S_1

These two figures show that the behavior of the upper Lyapunov exponent does not make implicit chaotic schemes more stable than the explicit ones, unlike in the classical Euler schemes. Any case may occur.

To end with, we computed the values of the different criteria from proposition 5.7, for different values of m. Table 2 summarizes the calculus.

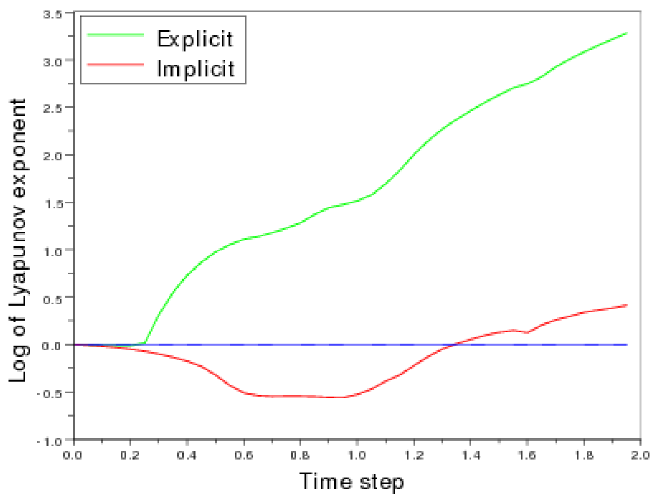

Figure 5: comparison of chaotic asynchronous explicit and implicit Euler, system S_2

System, scheme	Lyapunov exp.	Crit. with $m = 2$	Crit. with $m = 3$
S_1 explicit	0.728	—	—
S_1 implicit	0.652	—	—
S_2 explicit	0.227	0.208	0.217
S_2 implicit	1.341	1.257	1.285

Table 2: Stability conditions. The first column correspond to conditions on h computed from the estimation of Lyapunov exponent, the second and third ones give conditions from the application of the proposition 5.7 with $m = 2$ and $m = 3$.

This table shows that system S_1 is an example of a situation where proposition 5.7 does not apply, as the sequence $\sup_{\|x\|=1} \prod_{1 \leq, i_1, \cdots, i_m \leq N} \|B_{i_1} \cdots B_{i_m} \cdot x\|$ does not converge quickly enough. On the other hand, with system S_2, one can easily compute conditions on h without estimating the Lyapunov exponent.

These simple examples exhibit the fact that a systematic application of chaotic asynchronous methods leads to quite unpredictable systems, as regards their absolute stability.

7 Conclusion and proposition

Chaotic asynchronous schemes for resolving ordinary differential systems have shown their interest in the context of real time interactive simulation of multi interaction systems, especially when dealing with biochemical kinetics. Their main advantage is the capability that is given to the user to add or remove interactions, e.g. chemical reactions or forces, in the course of a simulation. Nevertheless, eventhough proofs of convergence for such schemes have been established, the present work highlights the fact that absolute stability conditions may be difficult to satisfy, when antagonist phenomena are desynchronized: antagonist forces can lead to force the choice of tiny time steps, making impossible the aim of real-time simulation. An illustration is given by the case of a mass-spring system. Therefore, a compromise has to be found between a total desynchronization of phenomena, which leads to instability, and synchronization, which prevents from in virtuo experimentation.

We propose, in this perspective, to adapt splitting methods Mclachlan and Quispel (2002) in order to keep the capacity of interacting by adding or removing phenomena. Indeed, splitting methods seem to be relevant when the phenomena involved in a simulated system have to be considered as autonomous: as in the case of chaotic asynchronous schemes, the resolution of a system $y' = (A+B)y$ is replaced by successive resolutions of systems $y' = Ay$ and $y' = By$. The use of different time steps for each of the subsystems permits to simulate multiscale systems. This is also possible with chaotic asynchronous schemes, but splitting methods have the advantage of absolute stability, by the use of particular scheduling of integrations of each subsystem, each of which being solved by an absolutely stable method. Nevertheless, the choice of splitting methods makes it impossible to add or remove phenomena in the course of a simulation, without stopping the simulation and rewriting algorithms with the new set of phenomena involved.

We have recently developed algorithms that can be seen as an hybridation between chaotic asynchronous schemes and splitting methods: a future work will expose these methods and achieve their theoretical study.

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